

Zoektocht naar de elementaire bouwstenen van de natuur



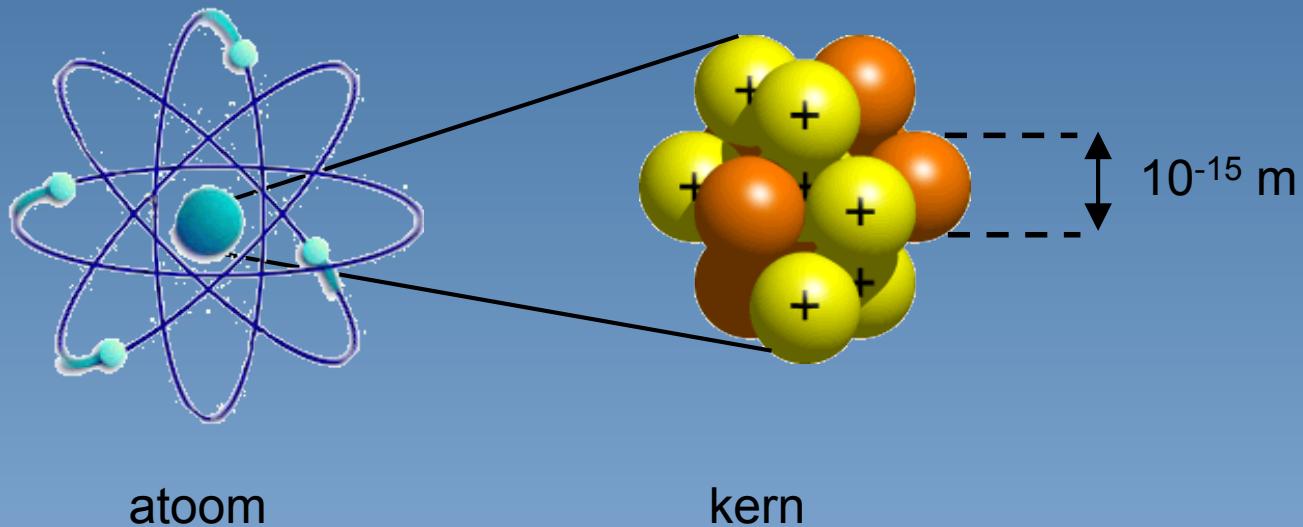
Week-end van de wetenschap, Groningen, 6 oktober 2013

Ivo van Vulpel

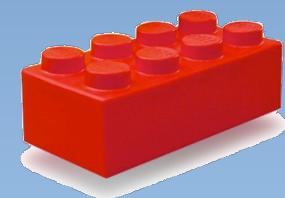
CERN in Genève, Zwitserland



Deeltjesfysica



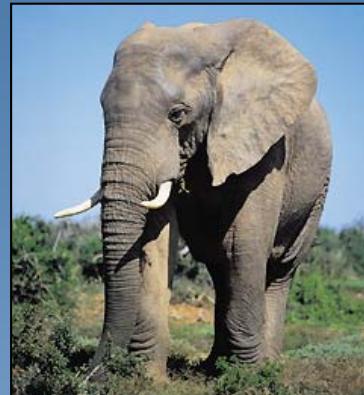
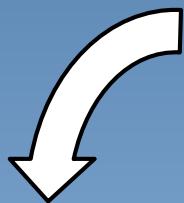
Wat zijn de bouwstenen van de dingen om ons heen ?



Het grootste en het kleinste

... volgens mijn dochter van 3

... volgens haar vader

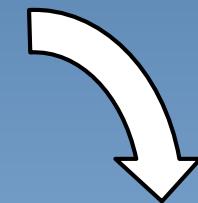


Olifant

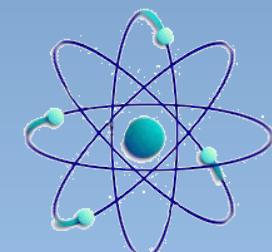
Grootst



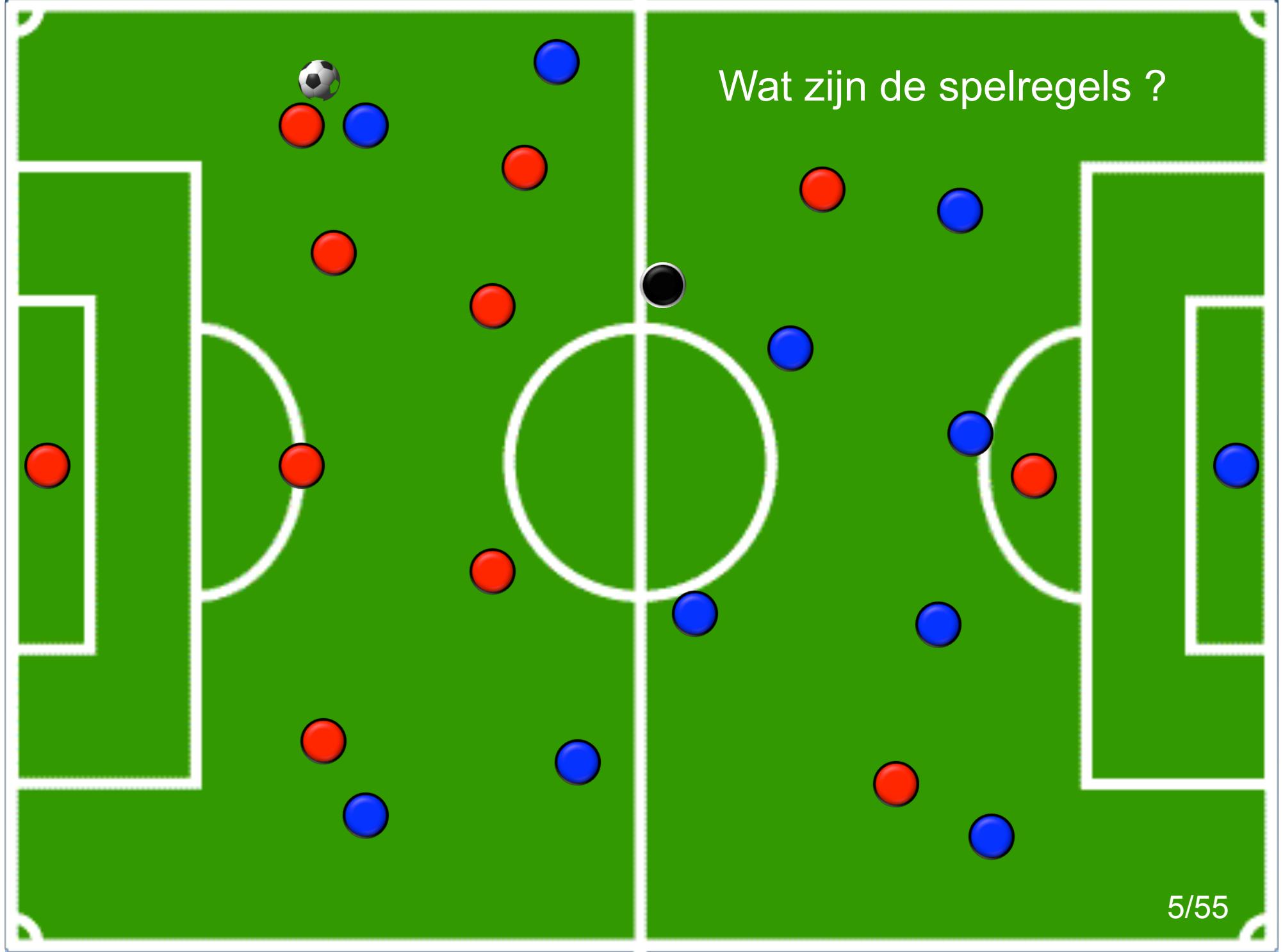
Klein muisje



10^{+22} m



10^{-9} m

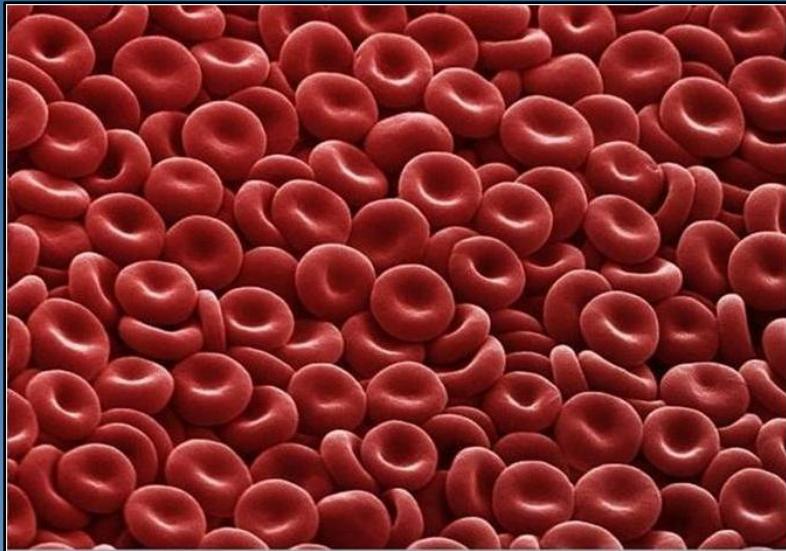


Röntgen foto

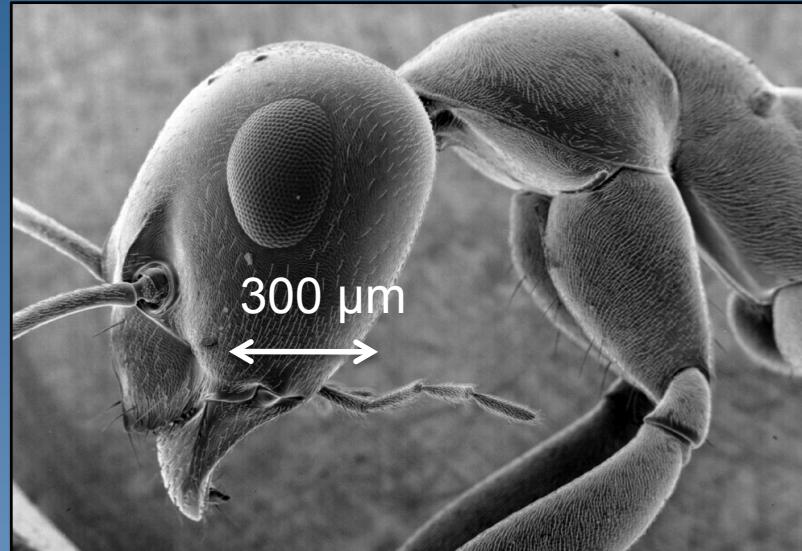


Het zichtbaar maken van
(niet-)afgeketste deeltjes

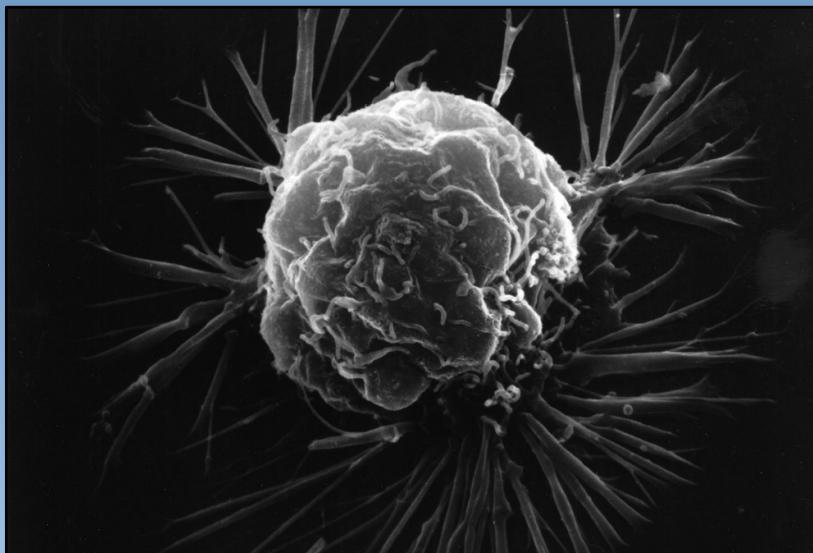
Rode bloedlichaampjes



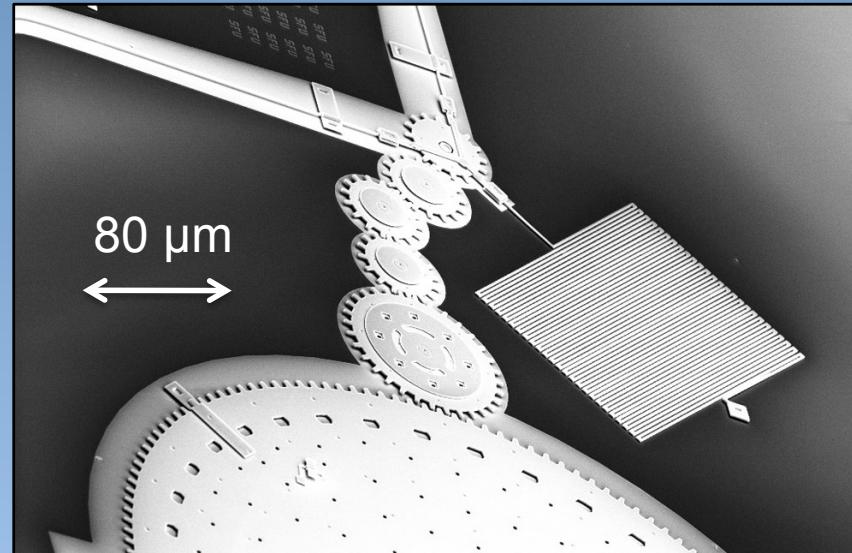
Mier



Borstkanker cel



Microelectronisch systeem



Van *wat* naar *hoe*

oogkleur ouders en kind

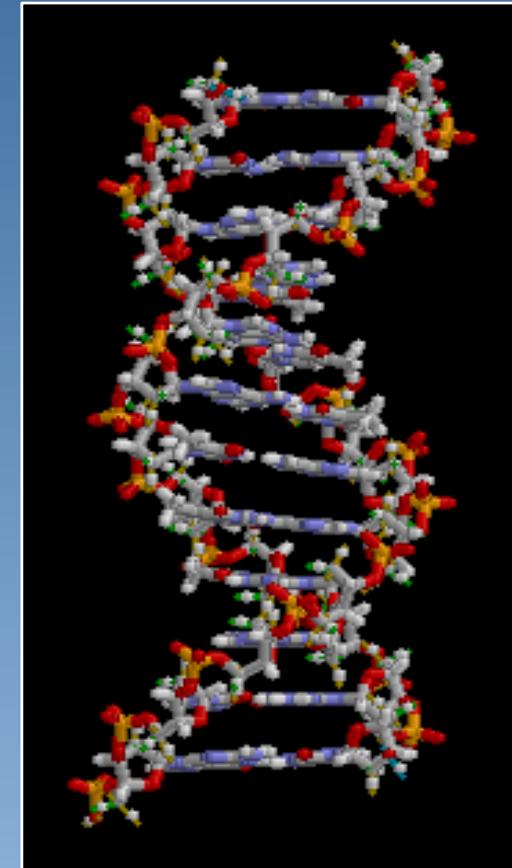
pa

ma

kans op oogkleur baby



= ?

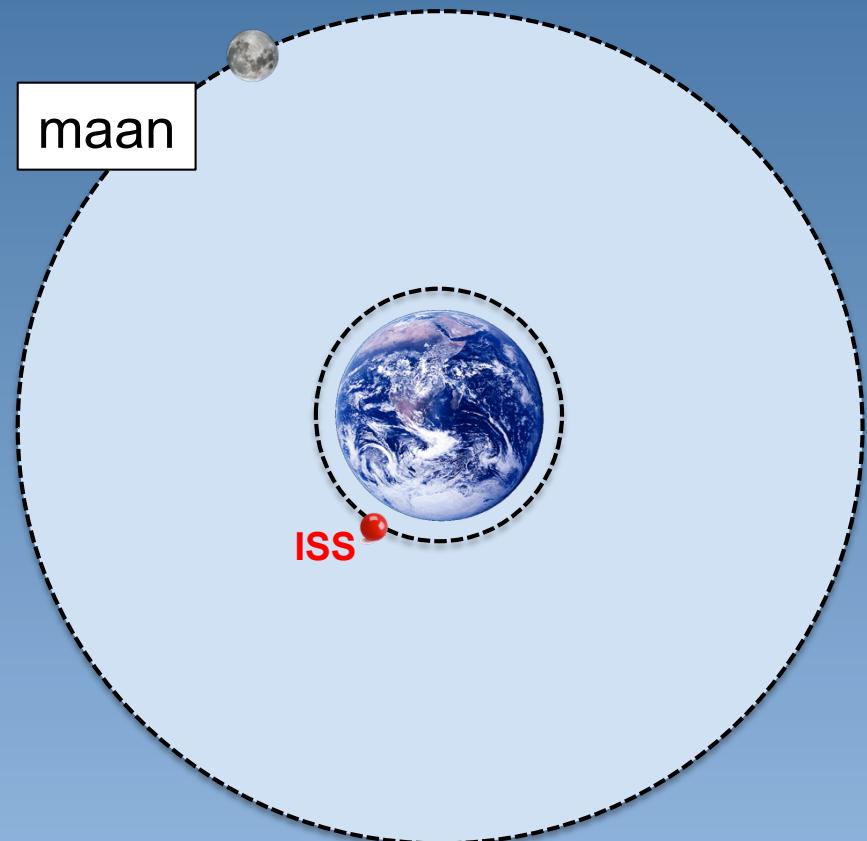


DNA: GTCA-nucleotiden

De atoomwereld: nieuwe fenomenen



Klassieke mechanica



We begrijpen waarom de maan om de aarde draait

$$v \propto \frac{1}{\sqrt{r}}$$

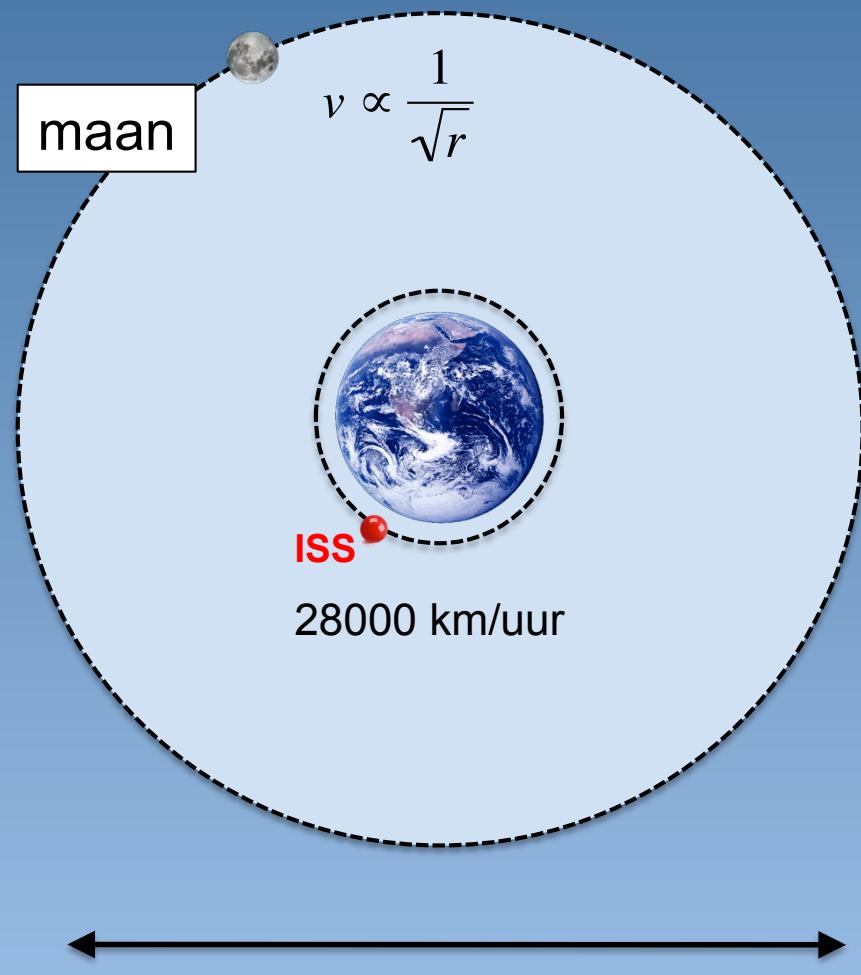
ISS: 400 km = 28.000 km/uur



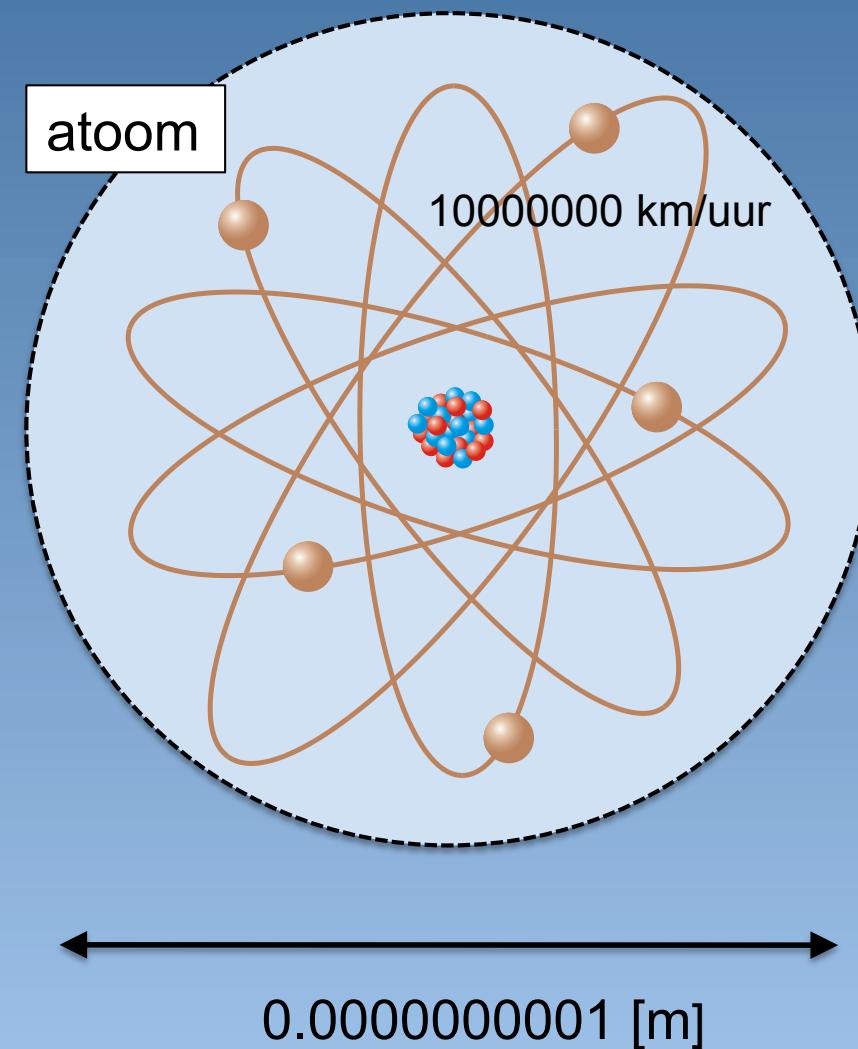
satelliet (pioneer/ISS)

$$\begin{aligned}\ddot{\mathbf{r}}_i = & \sum_{j \neq i} \frac{\mu_j (\mathbf{r}_j - \mathbf{r}_i)}{r_{ij}^3} \left(1 - \frac{2(\beta + \gamma)}{c^2} \sum_{k \neq i} \frac{\mu_k}{r_{ik}} - \frac{2\beta - 1}{c^2} \sum_{k \neq j} \frac{\mu_k}{r_{jk}} - \frac{3}{2c^2} \left[\frac{(\mathbf{r}_j - \mathbf{r}_i) \dot{\mathbf{r}}_j}{r_{ij}} \right]^2 + \frac{1}{2c^2} (\mathbf{r}_j - \mathbf{r}_i) \ddot{\mathbf{r}}_j - \frac{2(1 + \gamma)}{c^2} \dot{\mathbf{r}}_i \dot{\mathbf{r}}_j + \right. \\ & \left. + \gamma \left(\frac{v_i}{c} \right)^2 + (1 + \gamma) \left(\frac{v_j}{c} \right)^2 \right) + \frac{1}{c^2} \sum_{j \neq i} \frac{\mu_j}{r_{ij}^3} \left([\mathbf{r}_i - \mathbf{r}_j] \cdot [(2 + 2\gamma) \dot{\mathbf{r}}_i - (1 + 2\gamma) \dot{\mathbf{r}}_j] \right) (\dot{\mathbf{r}}_i - \dot{\mathbf{r}}_j) + \frac{3 + 4\gamma}{2c^2} \sum_{j \neq i} \frac{\mu_j \ddot{\mathbf{r}}_j}{r_{ij}} \quad (3)\end{aligned}$$

Klassieke mechanica



Quantummechanica



De atoomwereld

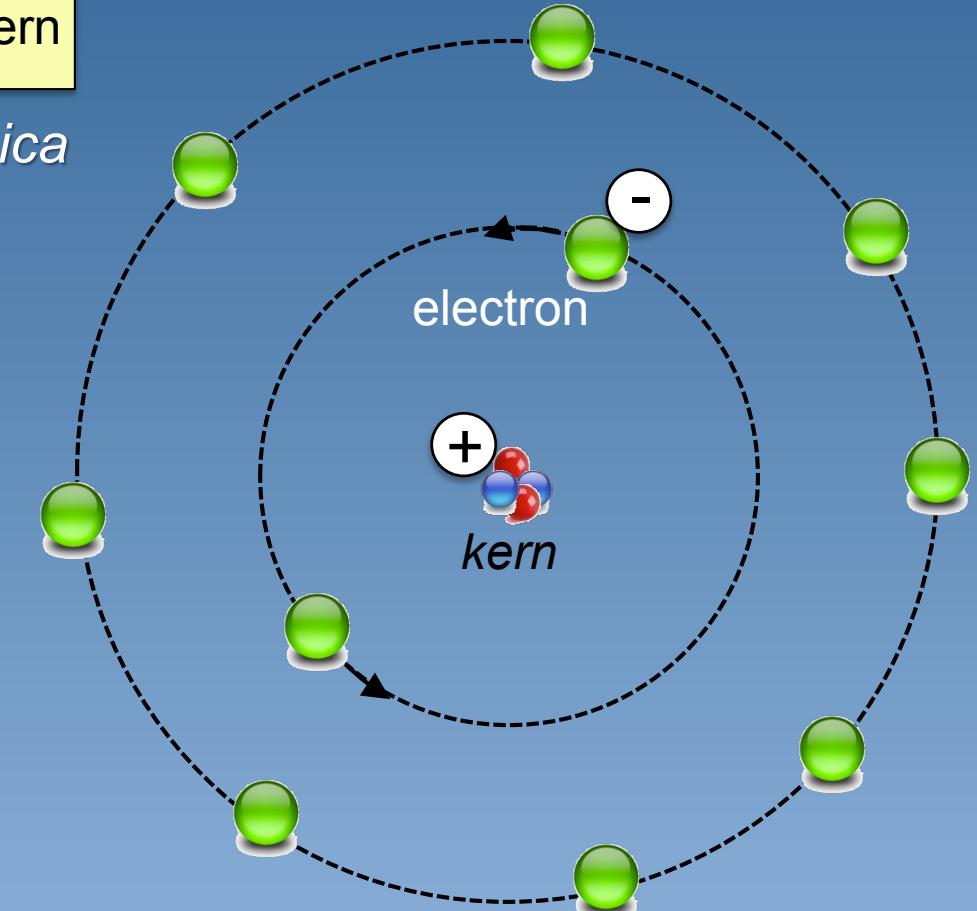
1) De elektronen vallen niet op de kern

Natuurwet: quantummechanica

2) De kerndeeltjes kleven aan elkaar

Natuurkracht: de kernkrachten

3) Maar 2 elektronen in de 1^e baan



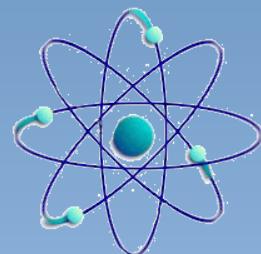
Fenomenen: spin en Pauli's uitsluitingsprincipe

Quantummechanica ... nuttig ?

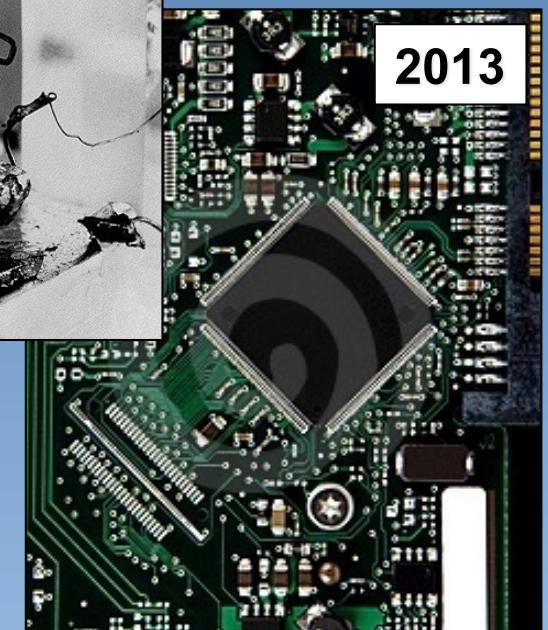
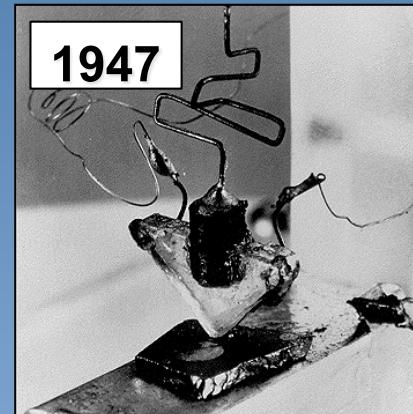
Quantummechanica

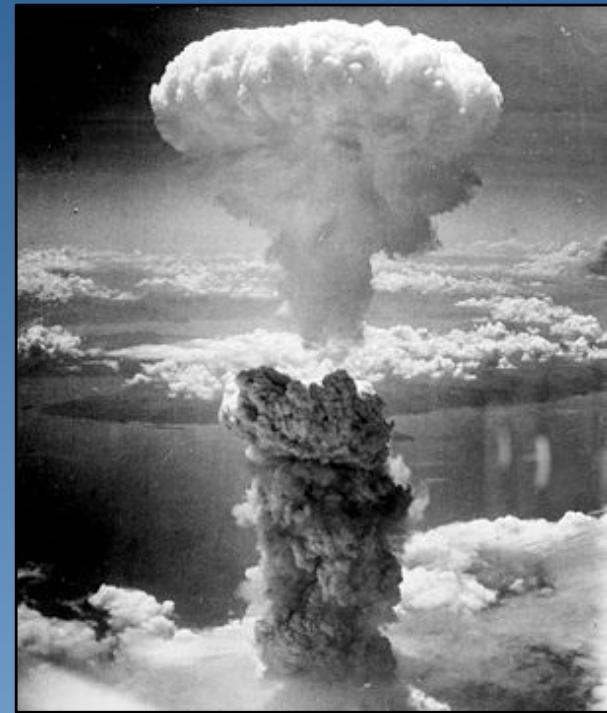


Niels Bohr

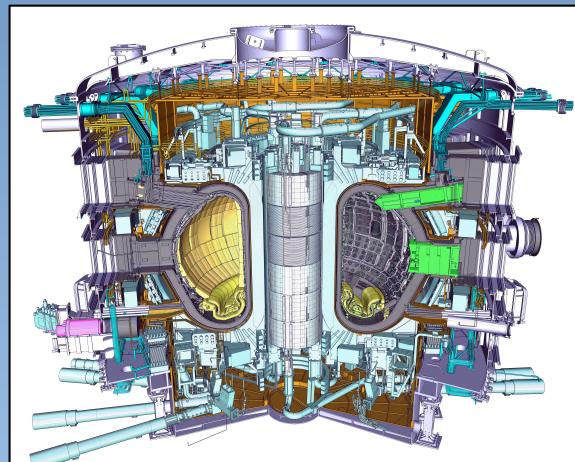


“Quantum mechanica en begrip van atomen cruciaal voor transistors.”





Kernkrachten nuttig ?



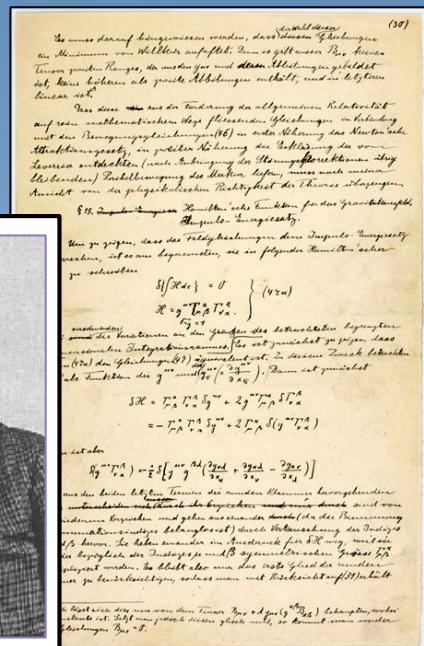
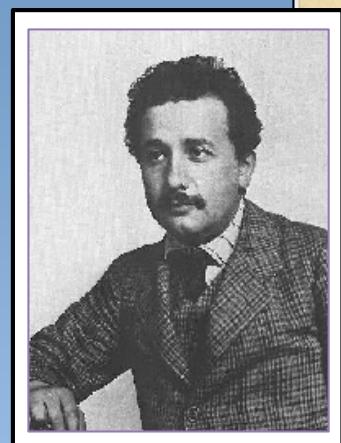
ITER (fusieractor)



Relativiteitstheorie ?

Net zo raar en onlogisch als de quantummechanica

- *tijd gaat langzamer als je hard beweegt*
 - *tjd gaat sneller als je verder weg bent van de aarde*



“Zonder relativiteitstheorie zit de GPS er 10 km per dag naast!”

De stand van zaken op 6 oktober 2013



[http:// pdg.lbl.gov](http://pdg.lbl.gov)

De elementaire deeltjes



up

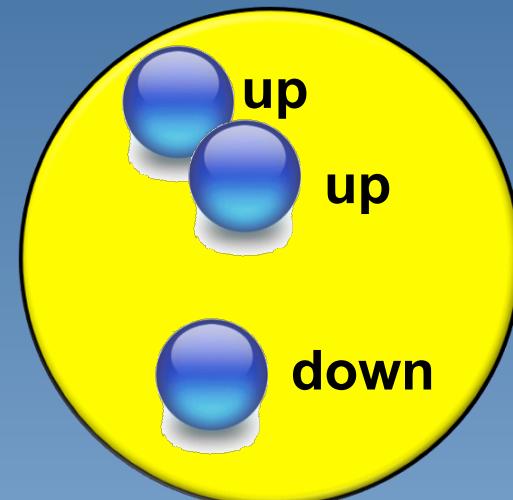


down

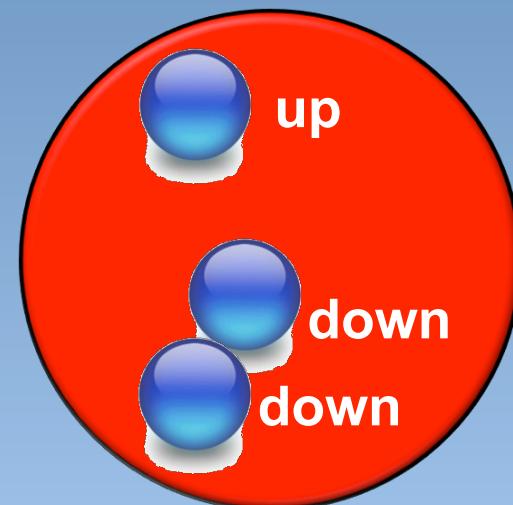


elektron

+ 4 natuurkrachten

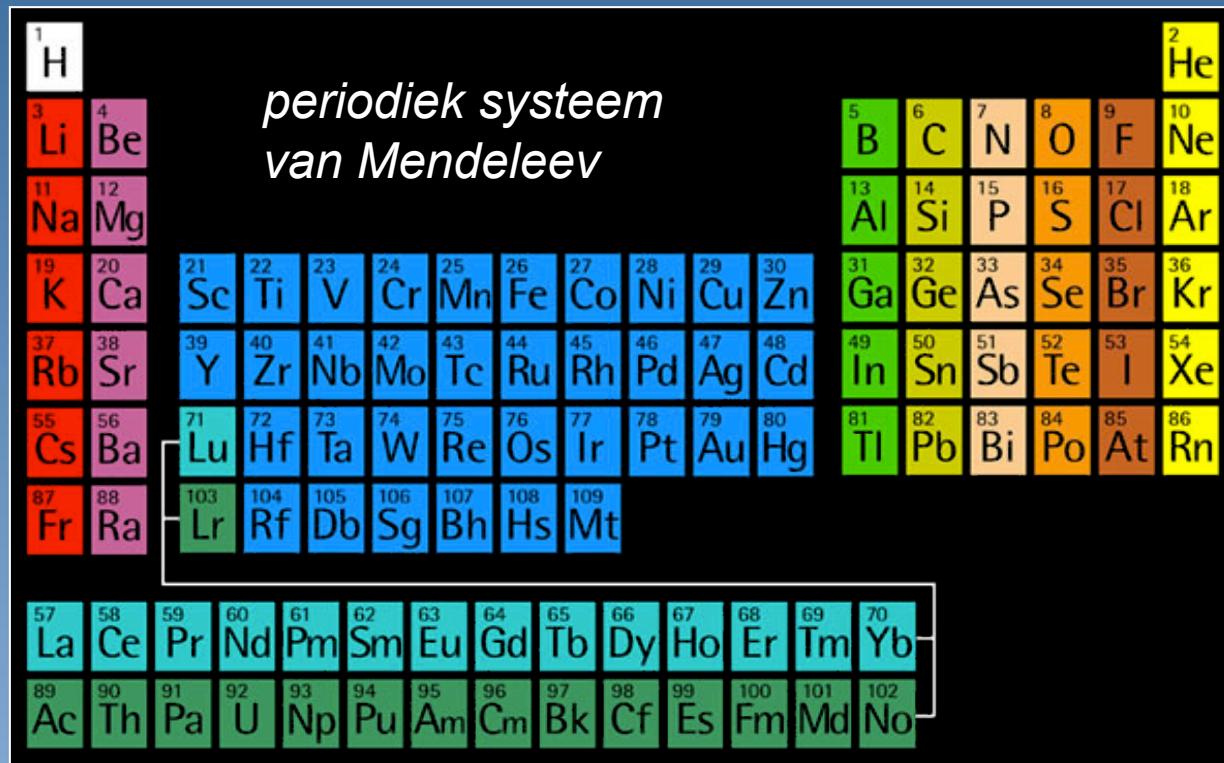


Proton



Neutron

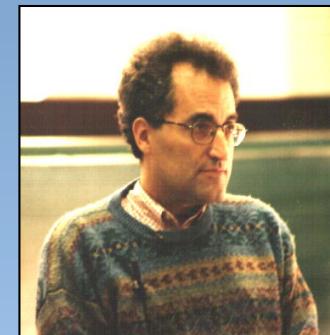
Wat kan je maken met deze 3 bouwstenen ?



Alles!



Alles wat we nodig hebben !

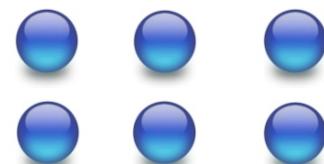


Edward Witten

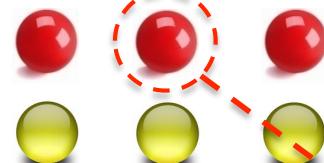
Deeltjes

Krachten

Quarks



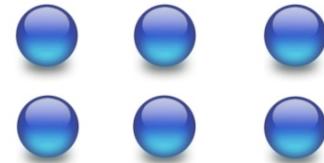
Leptonen



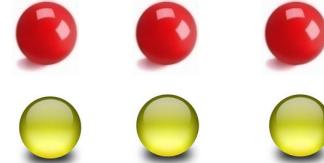
- 1) Electromagnetisme
- 2) Zwakke kernkracht
- 3) Sterke kernkracht
- 4) Zwaartekracht

Anti-deeltjes

Anti-
Quarks



Anti-
Leptonen



muon

Abstractie en patronen vinden

Standaard Model Lagrangiaan

$$SU(2)_L \otimes U(1)_Y \otimes SU(3)_C$$

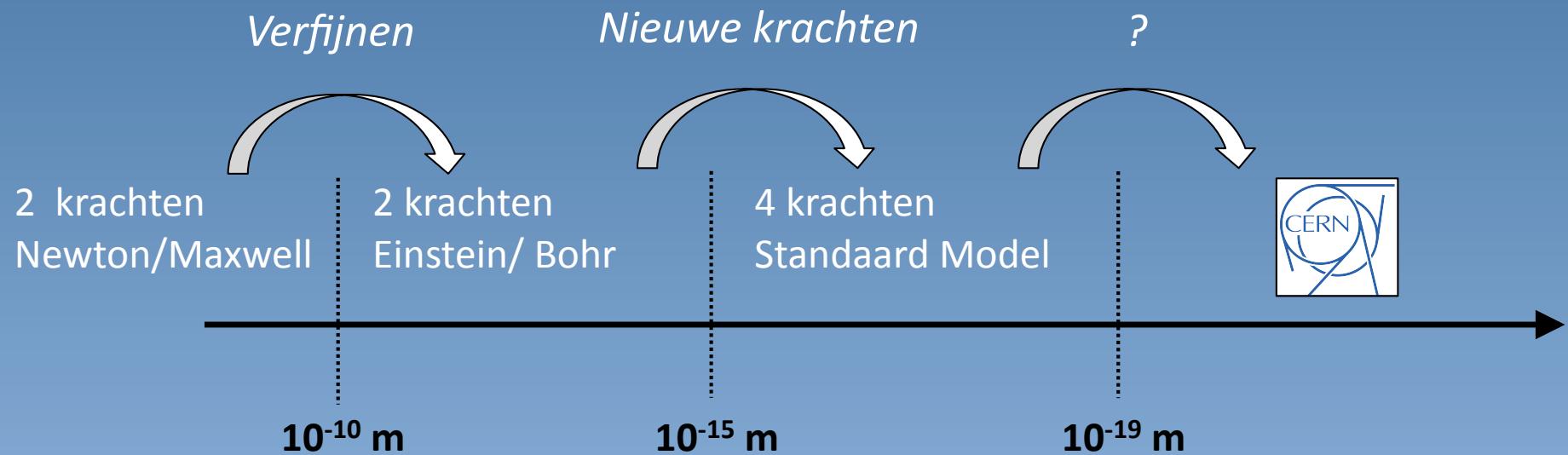
$$\begin{aligned}
& -\frac{1}{2}\partial_\nu g_\mu^a \partial_\nu g_\mu^a - g_s f^{abc} \partial_\mu g_\mu^a g_\mu^b g_\mu^c - \frac{1}{4}g_s^2 f^{abc} f^{ade} g_\mu^b g_\nu^c g_\mu^d g_\nu^e + \frac{1}{2}g_s^2 (q^\sigma \gamma^\mu q^\sigma) g_\mu^a + \\
& \bar{G}^a \partial^\mu G^a + g_s f^{abc} \partial_\mu \bar{G}^a g_\mu^b g_\mu^c - \partial_\mu W_\mu^+ \partial_\nu W_\mu^- - M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\nu Z_\mu^0 \partial_\mu Z_\mu^0 - \frac{1}{2c_w^2} M^2 Z_\mu^0 Z_\mu^0 - \\
& \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - \frac{1}{2}\partial_\mu H \partial_\mu H - \frac{1}{2}m_H^2 H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - M^2 \phi^+ \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \\
& \frac{1}{2c_w^2} M \phi^0 \phi^0 - \beta_h [\frac{2M^2}{g^2} + \frac{2M}{g} H + \frac{1}{2}(H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-)] + \frac{2M^4}{g^2} \alpha_h - ig c_w [\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\
& W_\mu^+ W_\nu^-) - Z_\mu^0 (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+) + Z_\mu^0 (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+)] - ig s_w [\partial_\nu A_\mu (W_\mu^+ W_\nu^- - \\
& W_\mu^- W_\nu^+) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+) + A_\mu (W_\nu^+ \partial_\nu W_\mu^- - W_\nu^- \partial_\nu W_\mu^+)] - \frac{1}{2}g^2 W_\mu^0 W_\mu^- W_\nu^+ W_\nu^- + \\
& \frac{1}{2}g^2 W_\mu^+ W_\mu^- W_\nu^+ W_\nu^- + g^2 w^2 (Z_\mu^0 W_\mu^+ Z_\nu^0 W_\nu^- - Z_\mu^0 Z_\nu^0 W_\mu^+ W_\nu^-) + g^2 s^2 (A_\mu W_\mu^+ A_\nu W_\nu^- - \\
& A_\mu A_\nu W_\mu^+ W_\nu^-) + g^2 s_w c_w [A_\mu Z_\mu^0 (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - 2A_\mu Z_\mu^0 W_\mu^+ W_\nu^-] - g\alpha [H^3 + \\
& H\phi^0 \phi^0 + 2H\phi^+ \phi^-] - \frac{1}{8}g^2 \alpha_h H^4 + (\phi^0)^4 + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + \\
& 2(\phi^0)^2 H^2] - gMW^+ W_\mu^- H - \frac{1}{2}g \frac{M}{c_w^2} Z_\mu^0 Z_\mu^0 H - \frac{1}{2}ig [W_\mu^+ (\phi^0 \partial_\mu \phi^- - \phi^+ \partial_\mu \phi^0) - W_\mu^- (\phi^0 \partial_\mu \phi^+ - \\
& \phi^+ \partial_\mu \phi^0)] + \frac{1}{2}g [W_\mu^+ (H \partial_\mu \phi^- - \phi^+ \partial_\mu H) - W_\mu^- (H \partial_\mu \phi^+ - \phi^+ \partial_\mu H)] + \frac{1}{2}g \frac{1}{c_w} (Z_\mu^0 (H \partial_\mu \phi^0 - \\
& \phi^0 \partial_\mu H) - ig \frac{s_w}{c_w} M Z_\mu^0 (W_\mu^+ \phi^- - W_\mu^- \phi^+) + ig s_w A_\mu (W_\mu^+ \phi^- - W_\mu^- \phi^+) - ig \frac{1-2c_w^2}{2c_w^2} Z_\mu^0 (\phi^+ \partial_\mu \phi^- - \\
& \phi^- \partial_\mu \phi^+) + ig s_w A_\mu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) - \frac{1}{4}g^2 W_\mu^+ W_\mu^- [H^2 + (\phi^0)^2 + 2\phi^+ \phi^-] - \\
& \frac{1}{4}g^2 \frac{1}{c_w^2} Z_\mu^0 Z_\mu^0 [H^2 + (\phi^0)^2 + 2(2s_w^2 - 1)^2 \phi^+ \phi^-] - \frac{1}{2}g^2 \frac{s_w}{c_w} Z_\mu^0 g^0 (W_\mu^+ \phi^- + W_\mu^- \phi^+) - \\
& \frac{1}{2}ig \frac{s_w}{c_w} Z_\mu^0 H (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \frac{1}{2}g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + W_\mu^- \phi^+) + \frac{1}{2}ig^2 s_w A_\mu H (W_\mu^+ \phi^- - \\
& W_\mu^- \phi^+) - g^2 \frac{s_w}{c_w} (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+ \phi^- - g^1 s_w^2 A_\mu A_\nu \phi^+ \phi^- - e^1 (\gamma \partial + m_\lambda^1) e^\lambda - \\
& \bar{\nu}^\lambda \gamma \partial \nu^\lambda - \bar{u}_j^\lambda (\gamma \partial + m_u^\lambda) u_j^\lambda - \bar{d}_j^\lambda (\gamma \partial + m_d^\lambda) d_j^\lambda + ig s_w A_\mu [-(\bar{e}^\lambda \gamma^\mu e^\lambda) + \frac{2}{3} (u_j^\lambda \gamma^\mu u_j^\lambda) - \\
& \frac{1}{3} (d_j^\lambda \gamma^\mu d_j^\lambda)] + \frac{ig}{4c_w} Z_\mu^0 [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{e}^\lambda \gamma^\mu (4s_w^2 - 1 - \gamma^5) e^\lambda) + (u_j^\lambda \gamma^\mu (\frac{3}{3} s_w^2 - \\
& 1 - \gamma^5) u_j^\lambda) + (d_j^\lambda \gamma^\mu (1 - \frac{3}{3} s_w^2 - \gamma^5) d_j^\lambda)] + \frac{ig}{2\sqrt{2}} W_\mu^+ [(\bar{\nu}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (u_j^\lambda \gamma^\mu (1 + \\
& \gamma^5) C_{\lambda\kappa} d_j^\lambda)] + \frac{ig}{2\sqrt{2}} W_\mu^- [(\bar{e}^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (d_j^\lambda C_{\lambda\kappa}^1 \gamma^\mu (1 + \gamma^5) u_j^\lambda)] + \frac{ig}{2\sqrt{2}} \frac{m_\lambda^1}{M} [-\phi^+ (\bar{\nu}^\lambda (1 - \\
& \gamma^5) e^\lambda) + \phi^- (\bar{e}^\lambda (1 + \gamma^5) \nu^\lambda)] - \frac{g}{2} \frac{m_\lambda^1}{M} [H (\bar{e}^\lambda \nu^\lambda) + i \phi^0 (\bar{e}^\lambda \gamma^\mu e^\lambda)] + \frac{ig}{2M\sqrt{2}} \phi^+ [-m_u^0 (\bar{u}_j^\lambda C_{\lambda\kappa} (1 - \\
& \gamma^5) d_j^\lambda) + m_u^\lambda (\bar{u}_j^\lambda C_{\lambda\kappa} (1 + \gamma^5) d_j^\lambda)] + \frac{ig}{2M\sqrt{2}} \phi^- [m_d^0 (\bar{d}_j^\lambda C_{\lambda\kappa}^1 (1 + \gamma^5) u_j^\lambda) - m_d^\lambda (\bar{d}_j^\lambda C_{\lambda\kappa}^1 (1 - \\
& \gamma^5) u_j^\lambda) - \frac{g}{2} \frac{m_\lambda^1}{M} H (\bar{u}_j^\lambda u_j^\lambda) - \frac{g}{2} \frac{m_\lambda^1}{M} H (\bar{d}_j^\lambda d_j^\lambda)] + \frac{ig}{2} \frac{m_\lambda^1}{M} \phi^0 (\bar{u}_j^\lambda \gamma^\mu u_j^\lambda) - \frac{ig}{2} \frac{m_\lambda^1}{M} \phi^0 (\bar{d}_j^\lambda \gamma^\mu d_j^\lambda) + \\
& \bar{X}^+ (\beta^2 - M^2) X^+ + \bar{X}^- (\beta^2 - M^2) X^- + \bar{X}^0 (\partial^2 - \frac{M^2}{c_w^2}) X^0 + \bar{Y} \partial^2 Y + ig c_w W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \\
& \partial_\mu \bar{X}^+ X^0) + ig s_w W_\mu^+ (\partial_\mu \bar{Y} X^- - \partial_\mu \bar{X}^+ Y) + ig c_w W_\mu^- (\partial_\mu \bar{X}^- X^0 - \partial_\mu \bar{X}^+ X^+) + \\
& ig s_w W_\mu^- (\partial_\mu \bar{X}^- Y - \partial_\mu \bar{Y} X^+) + ig c_w Z_\mu^0 (\partial_\mu \bar{X}^+ X^0 - \partial_\mu \bar{X}^- X^-) + ig s_w A_\mu (\partial_\mu \bar{X}^+ X^+ - \\
& \partial_\mu \bar{X}^- X^-) - \frac{1}{2}gM [\bar{X}^+ X^0 H + \bar{X}^- X^- H + \frac{1}{c_w^2} \bar{X}^0 X^0 H] + \frac{1-2c_w^2}{2c_w^2} igM [\bar{X}^+ X^0 \phi^+ - \\
& \bar{X}^- X^- \phi^-] + \frac{1}{2}igM [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + ig M s_w [\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-] + \\
& \frac{1}{2}igM [\bar{X}^+ X^+ \phi^- - \bar{X}^- X^- \phi^+]
\end{aligned}$$

Bladmuziek (J.S. Bach)

bladmuziek



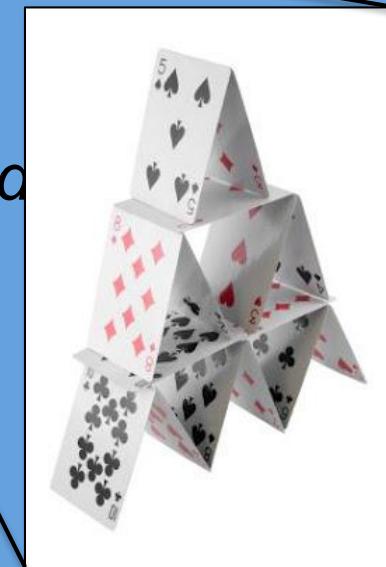
Nieuwe fenomenen, nieuwe krachten ??



deeltjes



Het Standaard Model





Theorie

Het Standaard Model

$$SU(2)_L \otimes U(1)_Y \otimes SU(3)_C$$

+ Higgs boson



Realiteit



Deeltjes hebben **gee**



- September 1964 -

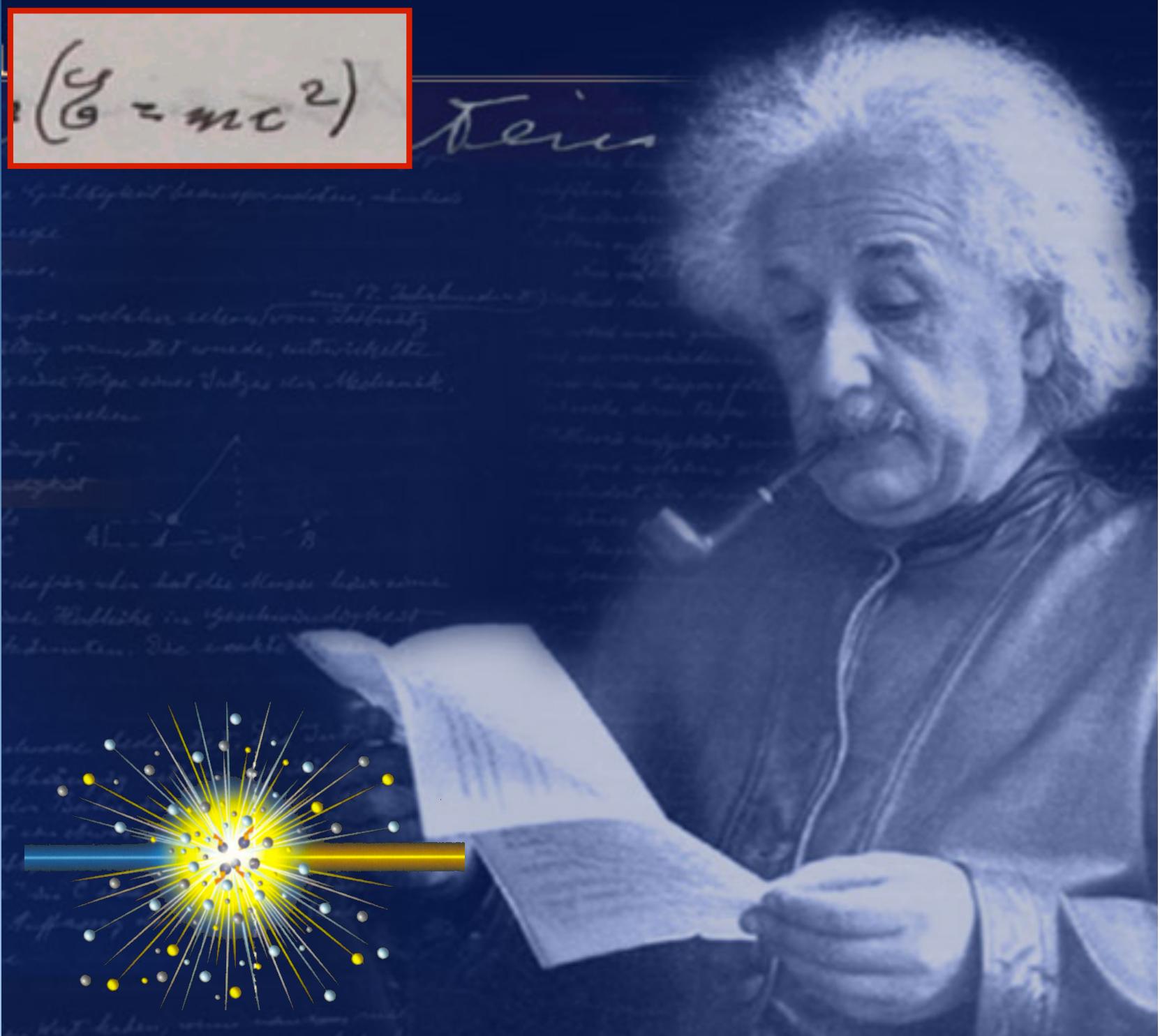
Het Higgs mechanism

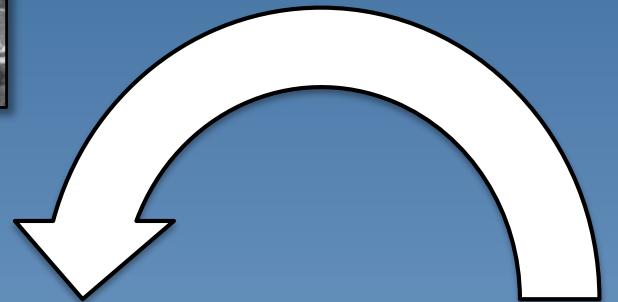
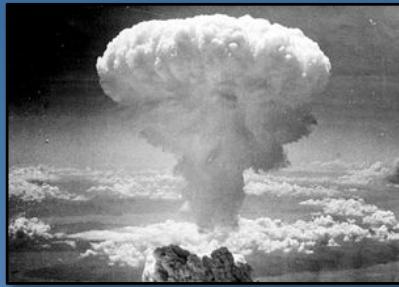
Deeltjes hebben **wel** massa

Higgs veld in het vacuum

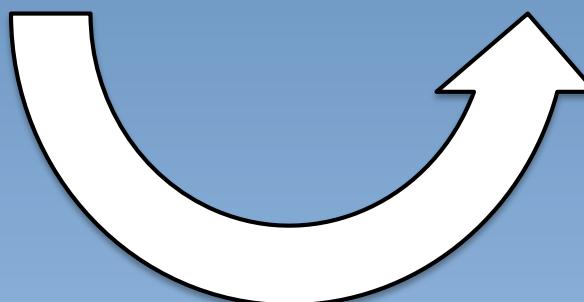
“Als ik gelijk heb moet er ook een nieuw deeltje bestaan: **het Higgs boson**”

Albert Einstein





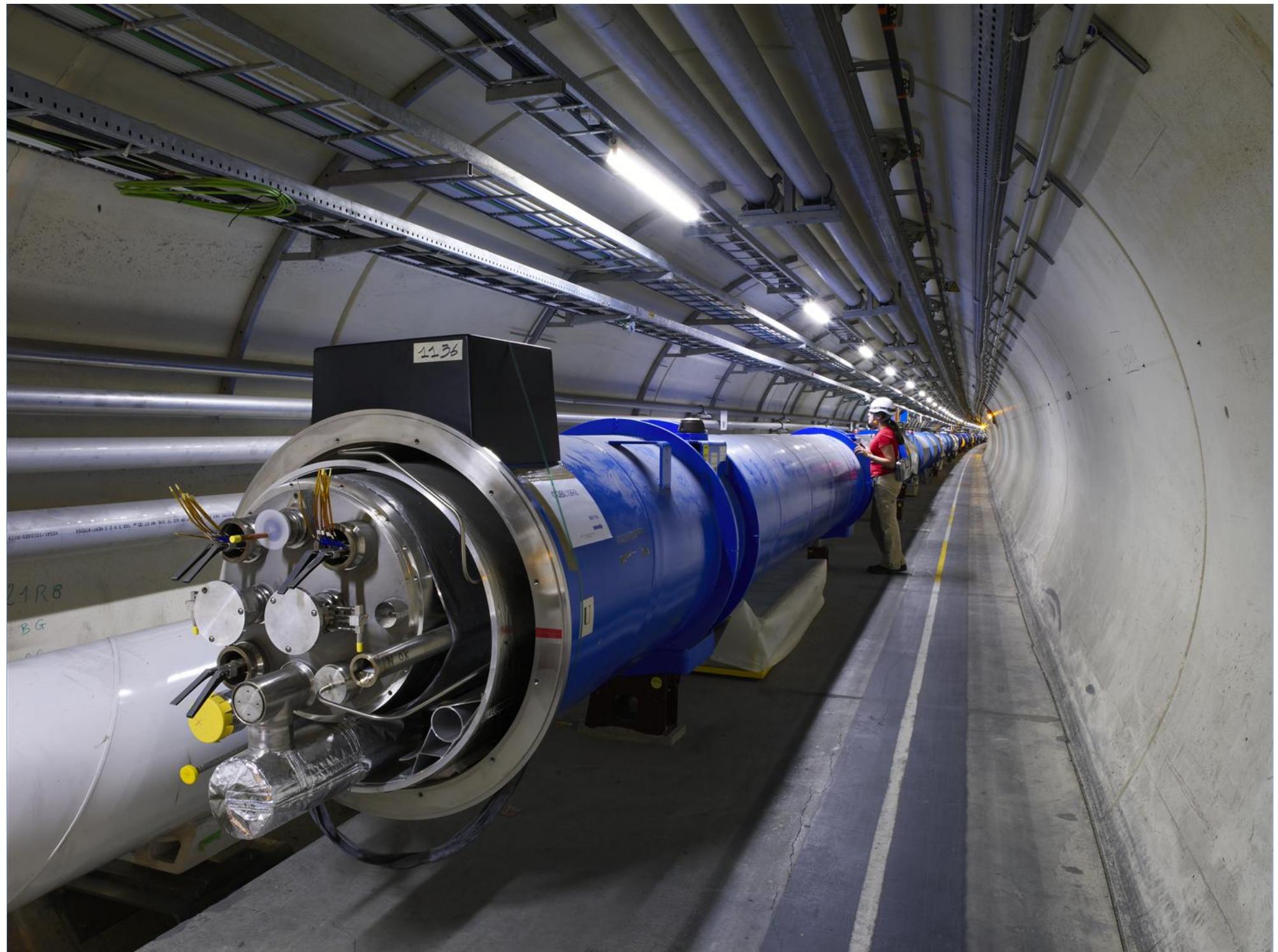
$$E = mc^2$$





Nieuwe deeltjes maken

De Large Hadron Collider (LHC)
op CERN bij Genève

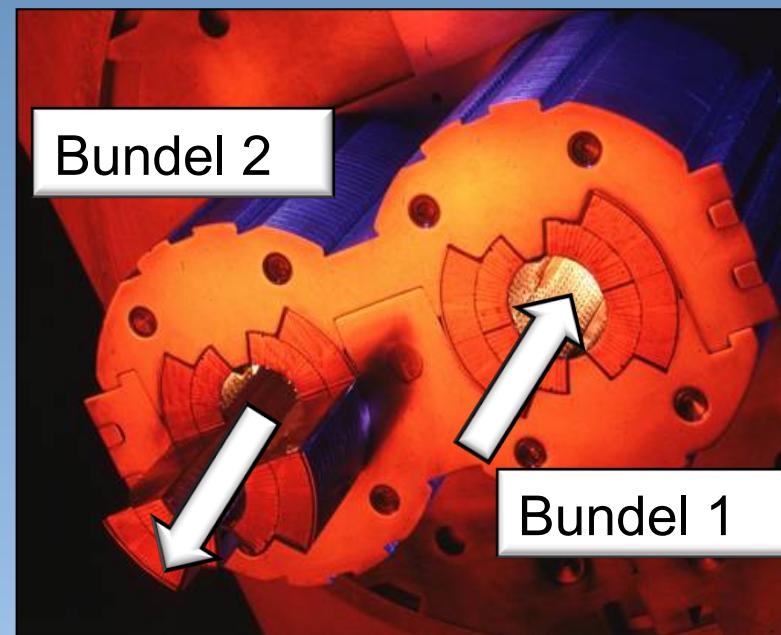


Energie 1 proton (LHC)



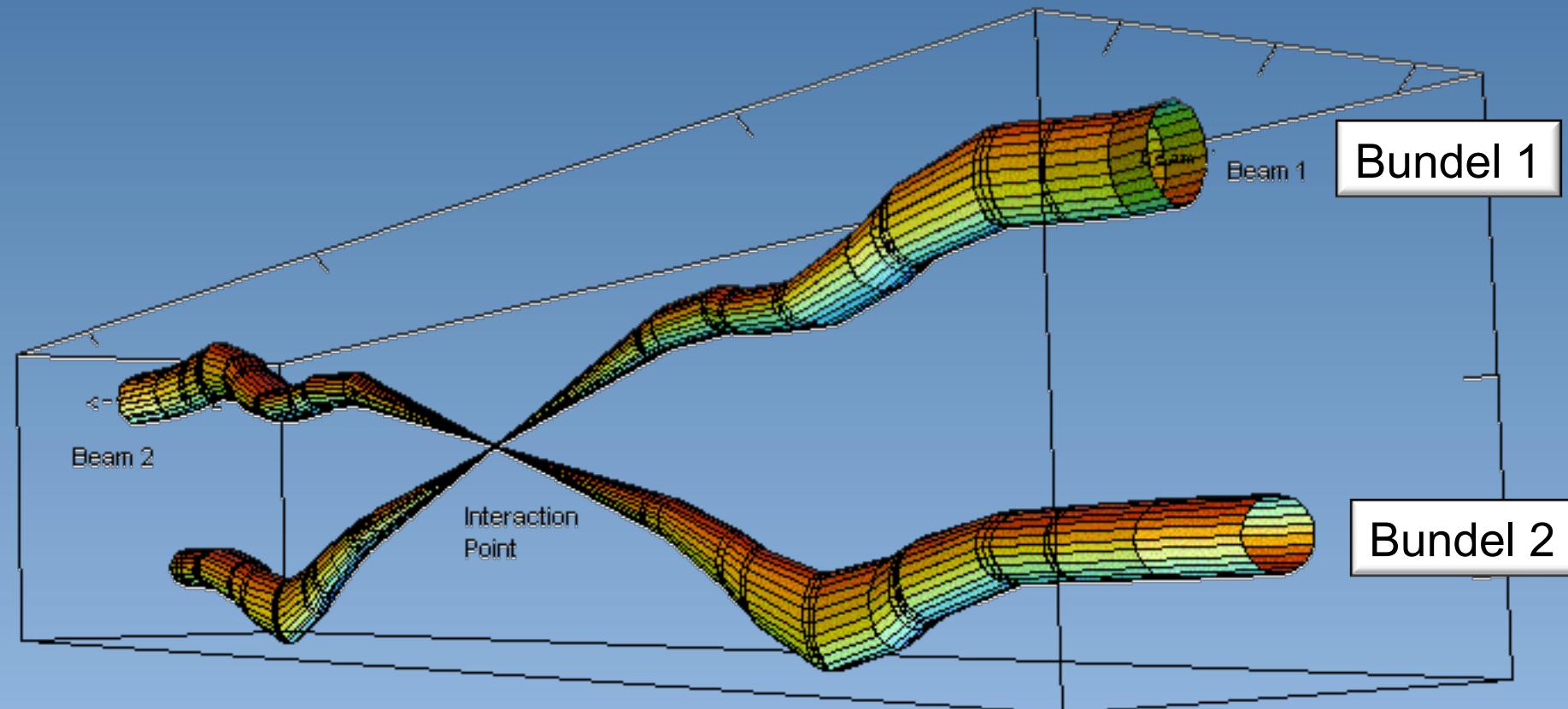
Energie LHC bundel

$3000 \times 100.000.000.000$ protonen



*Doorsnede LHC
bundelpijp*

40 miljoen botsingen per seconde

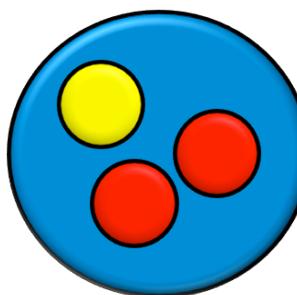




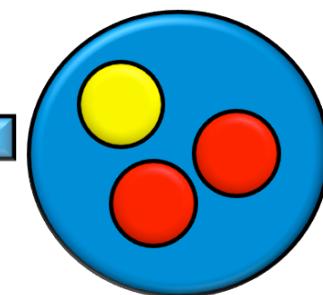
Klassiek botsen

Quantummechanisch botsen

proton



proton



Wat verwacht je ?

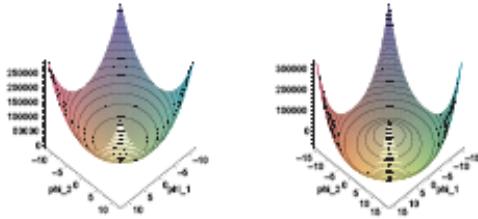


FIG. 1: The potential $V(\phi)$ ($\phi = \phi_1 + i\phi_2$) plotted for an arbitrary positive value of λ and for an arbitrary positive (right) or negative (left) value of μ^2 .

broken, and new interesting features emerge. To be more specific, let's pick the following ϕ_0 minimum (along the direction of the real part of ϕ , as traditional) and shift the ϕ field accordingly:

$$\phi_0 = \left(-\frac{\mu^2}{2\lambda} \right)^{1/2} = \frac{v}{\sqrt{2}} \quad \rightarrow \quad \phi(x) = \phi_0 + \frac{1}{\sqrt{2}}(\phi_1(x) + i\phi_2(x)) . \quad (5)$$

The Lagrangian in Eq. (3) can then be rearranged as follows:

$$\mathcal{L} = -\underbrace{\frac{1}{4}F_{\mu\nu}F_{\nu\rho}}_{\text{massive vector field}} + \underbrace{\frac{1}{2}g^2v^2A^\mu A_\mu}_{\text{massive scalar field}} + \underbrace{\frac{1}{2}(\partial^\mu\phi_1)^2 + \mu^2\phi_1^2 + \frac{1}{2}(\partial^\mu\phi_2)^2 + gvA_\mu\partial^\mu\phi_2}_{\text{Goldstone boson}} + \dots \quad (6)$$

and now contains the correct terms to describe a massive vector field A^μ with mass $m_A^2 = g^2v^2$ (originating from the kinetic term of \mathcal{L}_ϕ), a massive real scalar field ϕ_1 with mass $m_{\phi_1} = -2\mu^2$, that will become a Higgs boson, and a massless scalar field ϕ_2 , a so called Goldstone boson which couples to the gauge vector boson A^μ . The terms omitted contain couplings between the ϕ_1 and ϕ_2 fields irrelevant to this discussion. The gauge symmetry of the theory allows us to make the particle content more transparent. Indeed, if we parameterize the complex scalar field ϕ as:

$$\phi(x) = \frac{e^{i\chi(x)}}{\sqrt{2}}(v + H(x)) \xrightarrow{U(1)} \frac{1}{\sqrt{2}}(v + H(x)) , \quad (7)$$

the χ degree of freedom can be rotated away, as indicated in Eq. (7), by enforcing the $U(1)$ gauge invariance of the original Lagrangian. With this gauge choice, known as *unitary gauge* or *unitarity gauge*, the Lagrangian becomes:

$$\mathcal{L} = \mathcal{L}_A + \frac{g^2v^2}{2}A^\mu A_\mu + \frac{1}{2}\left(\partial^\mu H\partial_\mu H + 2\mu^2H^2\right) + \dots \quad (8)$$

a factor of two, making therefore almost a factor of four at the rate level. All other mass corrections, in the matrix element and phase space entering the calculation of the $H \rightarrow Q\bar{Q}$ decay rate, can in first approximation be safely neglected.

4. Loop induced Higgs boson decays: $H \rightarrow \gamma\gamma, \gamma Z, gg$

a factor of two, making therefore almost a factor of four at the rate level. All other mass corrections, in the matrix element and phase space entering the calculation of the $H \rightarrow Q\bar{Q}$ decay rate, can in first approximation be safely neglected.

Higgs boson decays: $H \rightarrow \gamma\gamma, \gamma Z, gg$

the $H\gamma\gamma$ and $H\gamma Z$ couplings are induced at one loop via both a

. At the lowest order the decay rate for $H \rightarrow \gamma\gamma$ can be written

$$= \frac{G_F\alpha^2 M_H^2}{128\sqrt{2}\pi^2} \left| \sum_f N_c^f Q_f^2 A_f^H(\tau_f) + A_W^H(\tau_W) \right|^2 , \quad (102)$$

($f = l, q$ respectively), Q_f is the charge of the f fermion species, and $f(\tau)$ is defined as:

$$(\tau) = \begin{cases} \arcsin^2 \frac{1}{\sqrt{\tau}} & \tau \geq 1 \\ -\frac{1}{4} \left[\ln \frac{1+\sqrt{1-\tau}}{1-\sqrt{1-\tau}} - i\pi \right]^2 & \tau < 1 \end{cases} , \quad (103)$$

and A_W^H are given by:

$$A_f^H = 2\tau [1 + (1 - \tau)f(\tau)] , \quad (104)$$

$$A_W^H(\tau) = -[2 + 3\tau + 3\tau(2 - \tau)f(\tau)] .$$

decay rate for $H \rightarrow \gamma Z$ is given by:

$$\frac{\alpha M_H^2}{\pi^4} \left(1 - \frac{M_Z^2}{M_H^2} \right)^2 \left| \sum_f A_f^H(\tau_f, \lambda_f) + A_W^H(\tau_W, \lambda_W) \right|^2 , \quad (105)$$

$\lambda_i = 4M_i^2/M_Z^2$ ($i = f, W$), and the form factors $A_f^H(\tau, \lambda)$ and

$$2N_c^f \frac{Q_f(I_3f - 2Q_f \sin^2 \theta_W)}{\cos \theta_W} [I_1(\tau, \lambda) - I_2(\tau, \lambda)] , \quad (106)$$

$\cos \theta_W \left\{ \left[\left(1 + \frac{2}{\tau} \right) \tan^2 \theta_W - \left(5 + \frac{2}{\tau} \right) \right] I_1(\tau, \lambda) \right. \right.$

$$+ 4(3 - \tan^2 \theta_W) I_2(\tau, \lambda) \Big\} , \quad (107)$$

and after Eq. (102), and I_3^f is the weak isospin of the f fermion

$$2N_c^f \frac{Q_f(I_3f - 2Q_f \sin^2 \theta_W)}{\cos \theta_W} [I_1(\tau, \lambda) - I_2(\tau, \lambda)] , \quad (106)$$

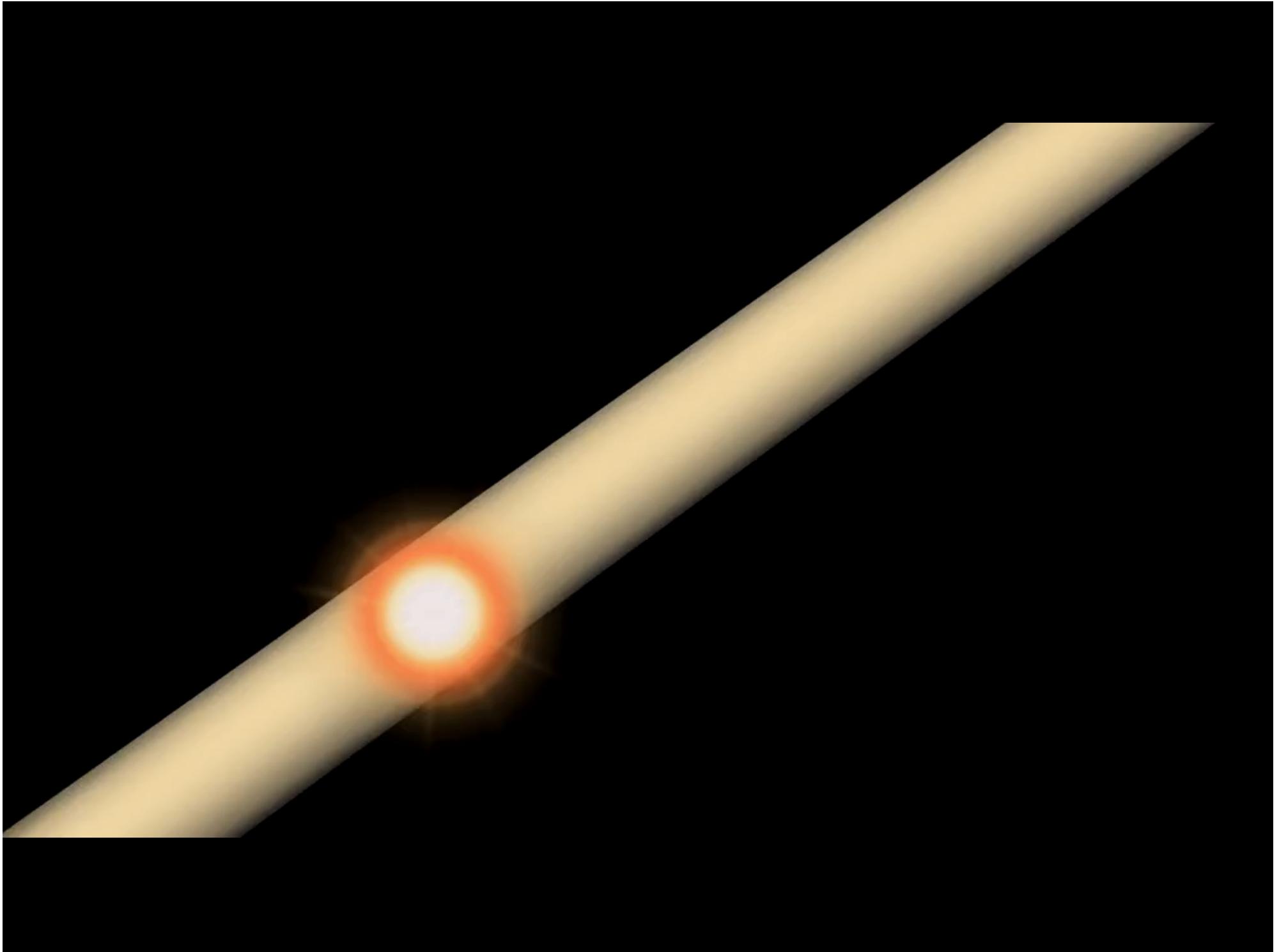
$\cos \theta_W \left\{ \left[\left(1 + \frac{2}{\tau} \right) \tan^2 \theta_W - \left(5 + \frac{2}{\tau} \right) \right] I_1(\tau, \lambda) \right. \right.$

$$+ 4(3 - \tan^2 \theta_W) I_2(\tau, \lambda) \Big\} , \quad (107)$$

and after Eq. (102), and I_3^f is the weak isospin of the f fermion

$$\frac{\tau^2 \lambda^2}{\lambda} + \frac{\tau^2 \lambda^2}{2(\tau - \lambda)^2} [f(\tau) - f(\lambda)] + \frac{\tau^2 \lambda}{(\tau - \lambda)^2} [g(\tau) - g(\lambda)] ,$$

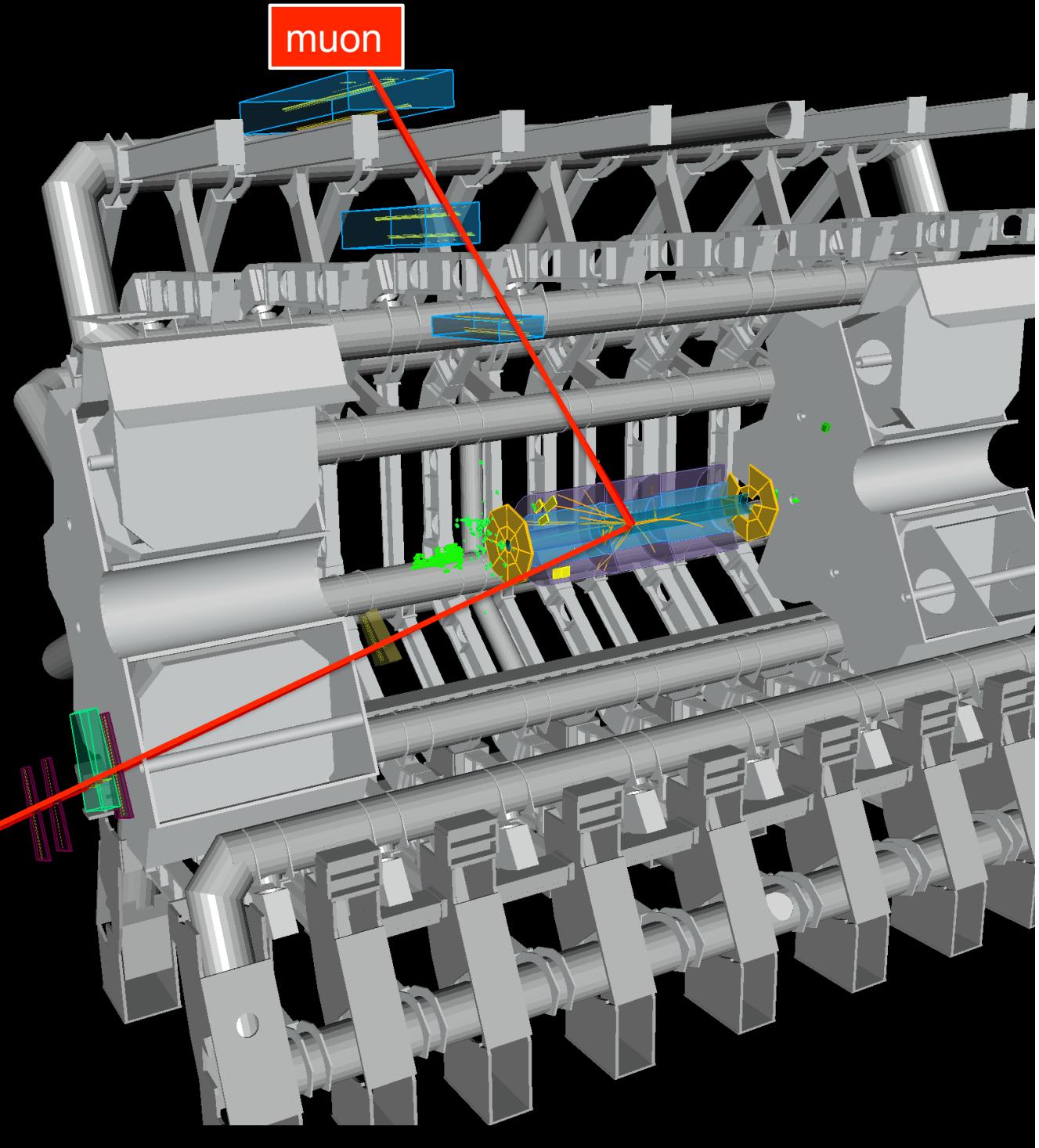
Inderdaad nieuwe deeltjes ?



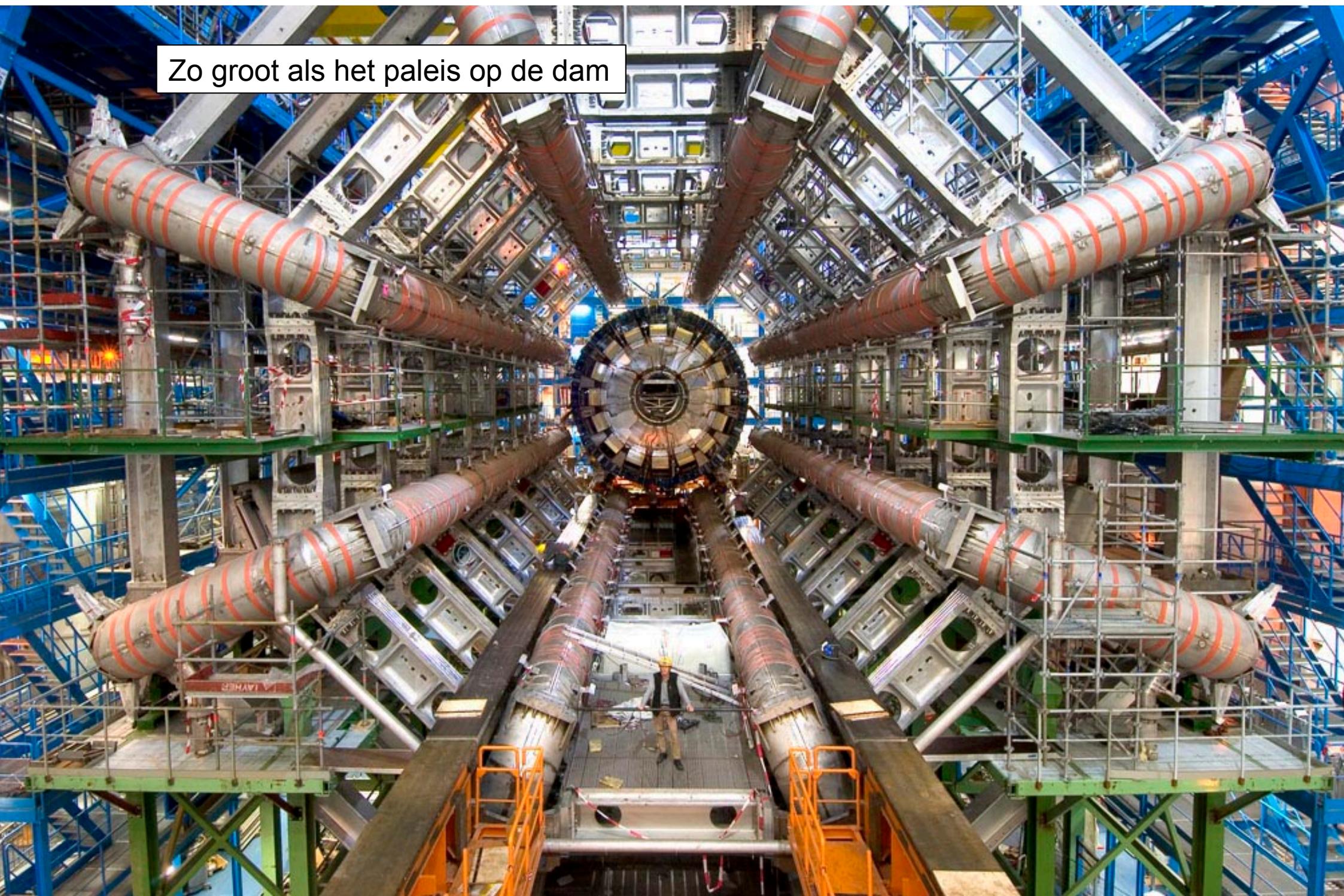
Onze ‘fototoestellen’



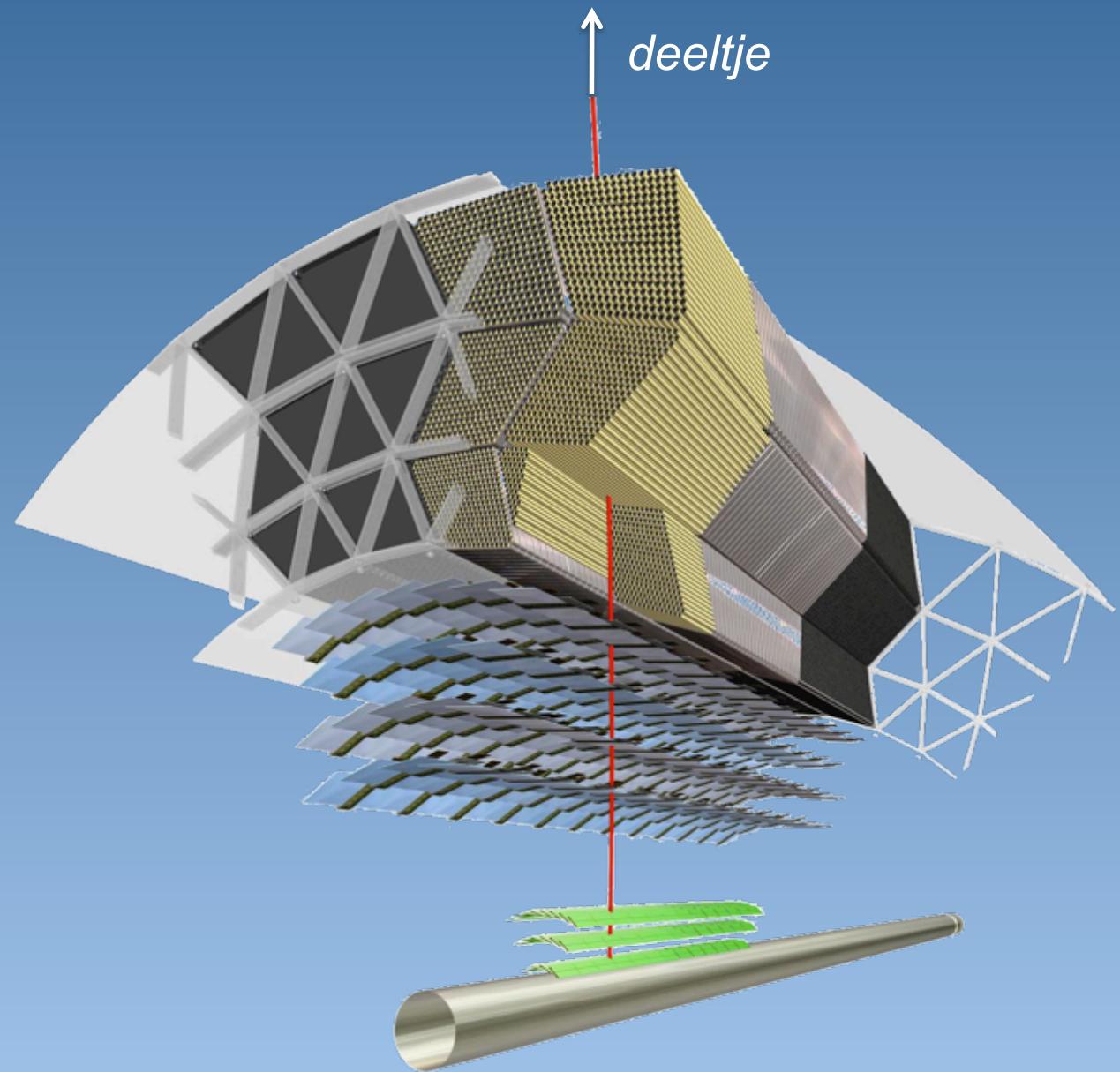
'foto' van een
echte botsing



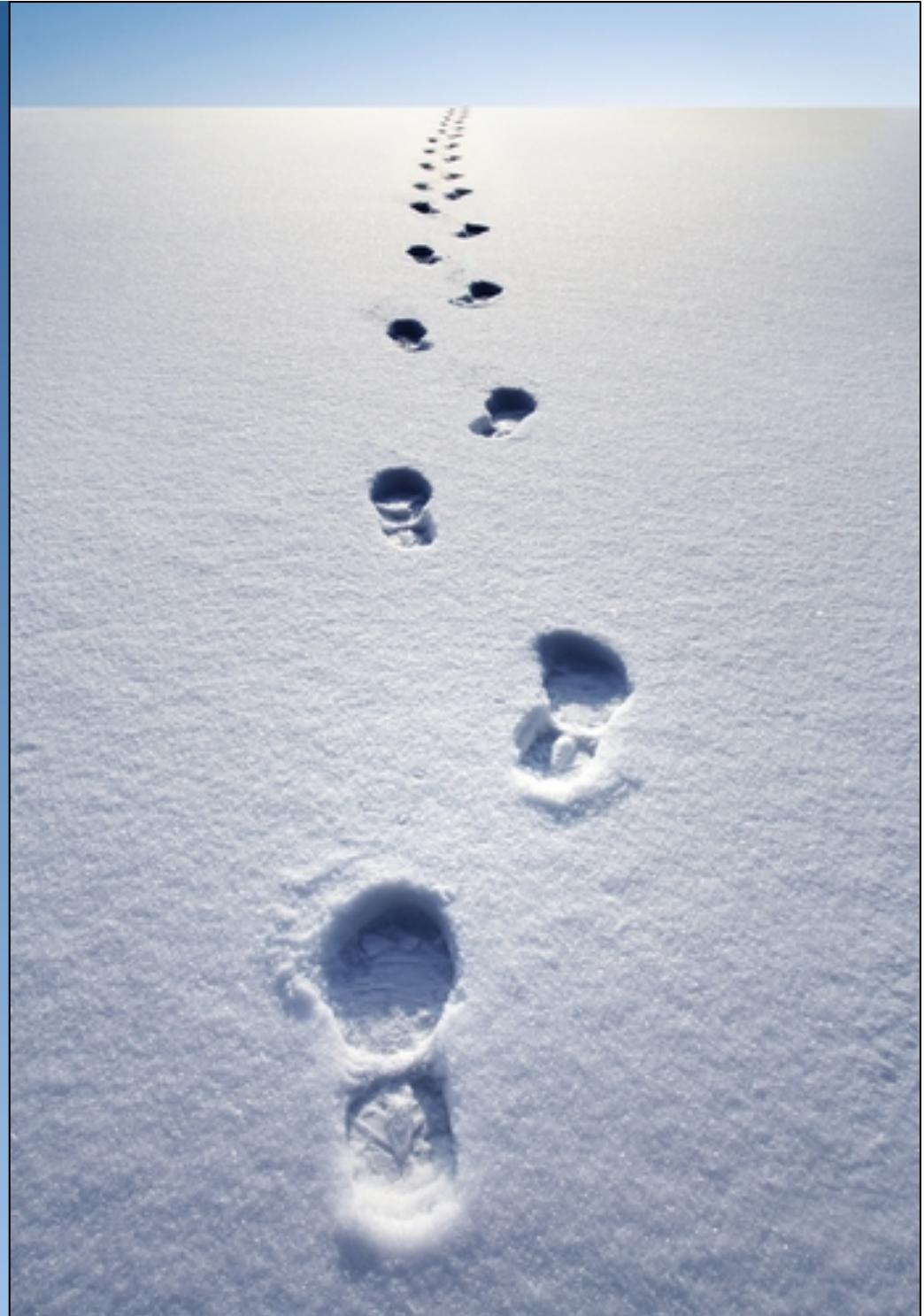
Zo groot als het paleis op de dam



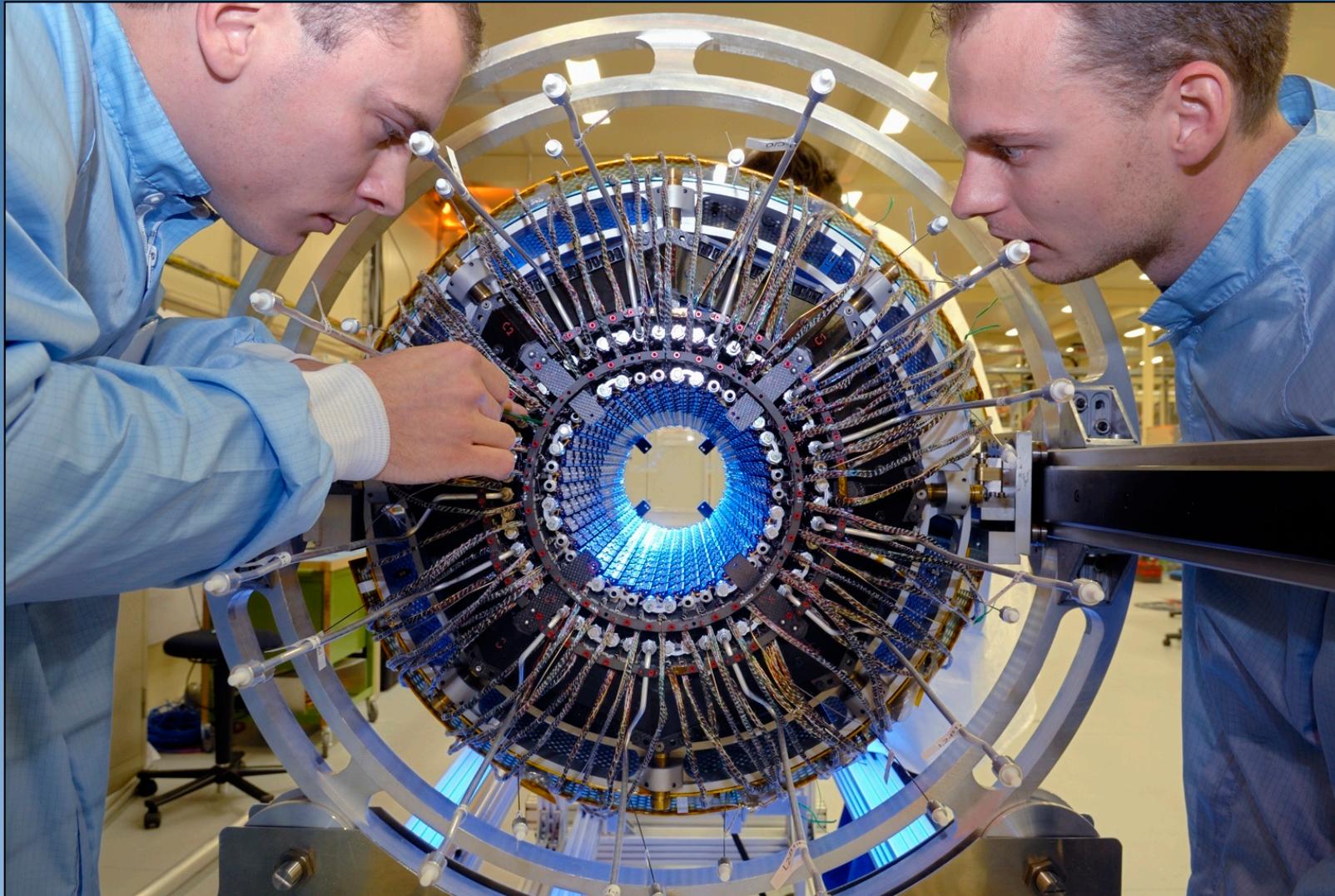
Waarom is die ATLAS detector nou zo complex ?



Konijn, auto of mens ?

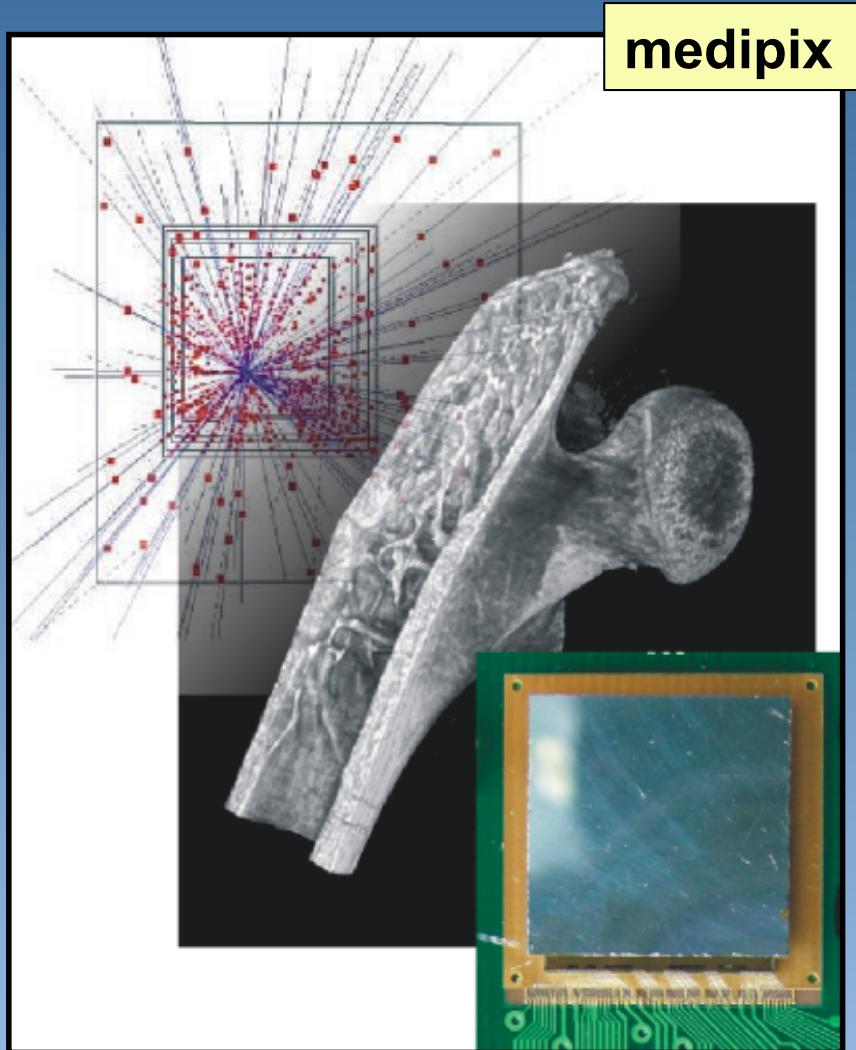


De Atlas pixel detector



80 MegaPixel camera 40.000.000 foto's per seconde

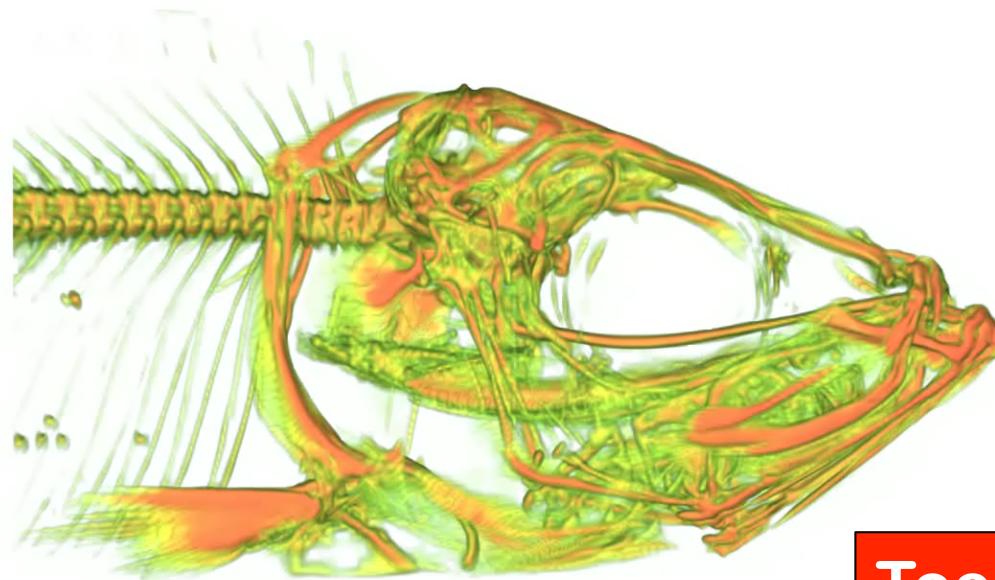
Medische toepassingen



Normale CT scan ~ 200 Röntgen foto's

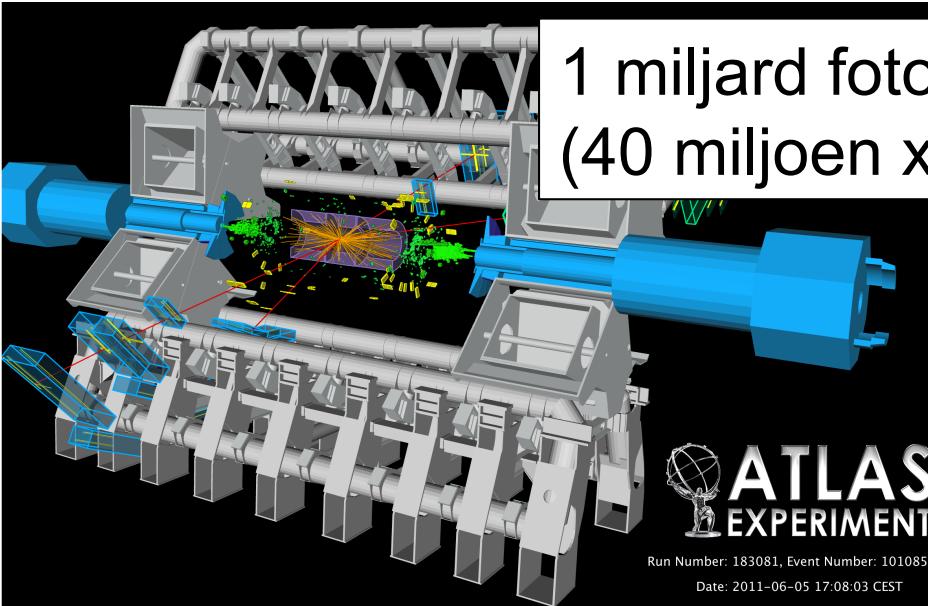
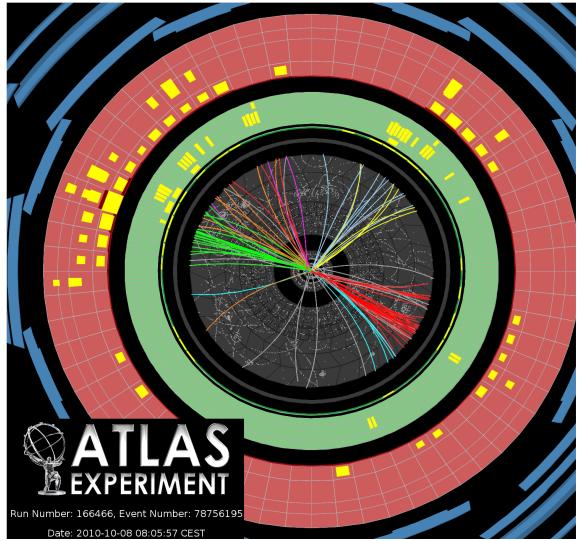
Kleinere pixel grootte: 1 mm → 50 µm

CT scan of a (spiering) of 2.4 cm length using MedPix

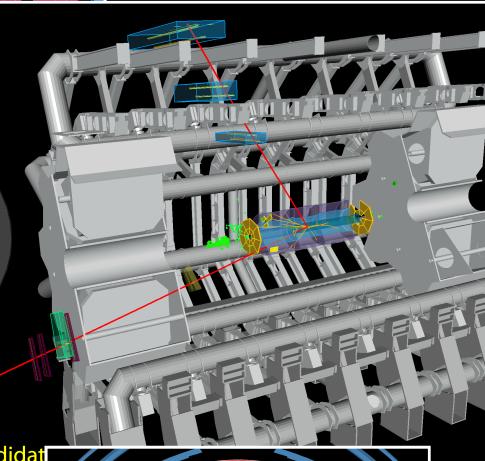
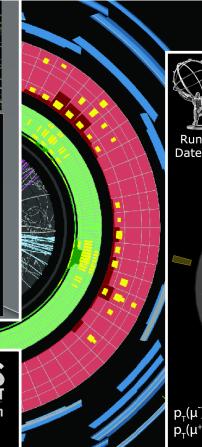
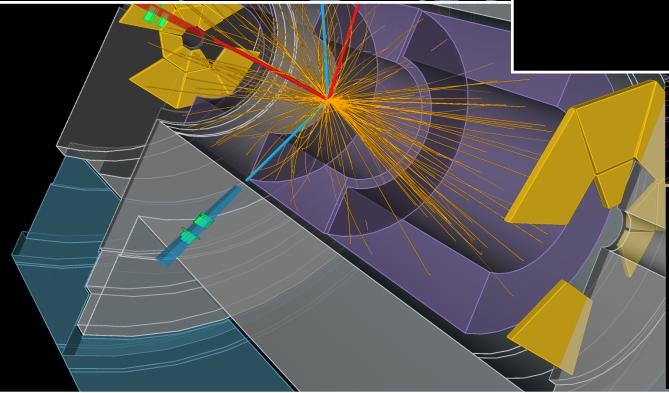


Toekomst:
kleuren-Xrays

Toepassing: minder X-rays voor zelfde resolutie

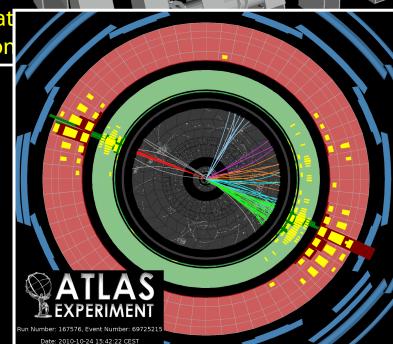
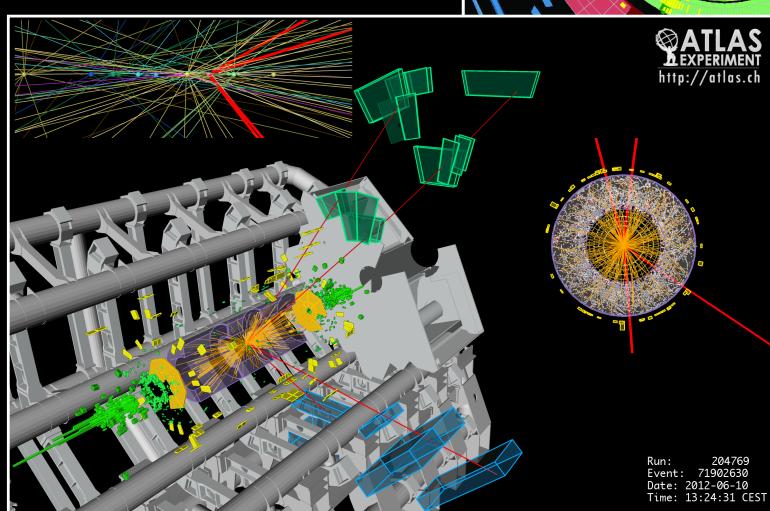


1 miljard foto's per seconde
(40 miljoen x 20 botsingen)



$p_t(\mu) = 27 \text{ GeV}$ $\eta(\mu) = 0.7$
 $p_t(\mu') = 45 \text{ GeV}$ $\eta(\mu') = 2.2$
 $M_{\mu\mu} = 87 \text{ GeV}$

Z- $\mu\mu$ candidate
in 7 TeV collision



Computing

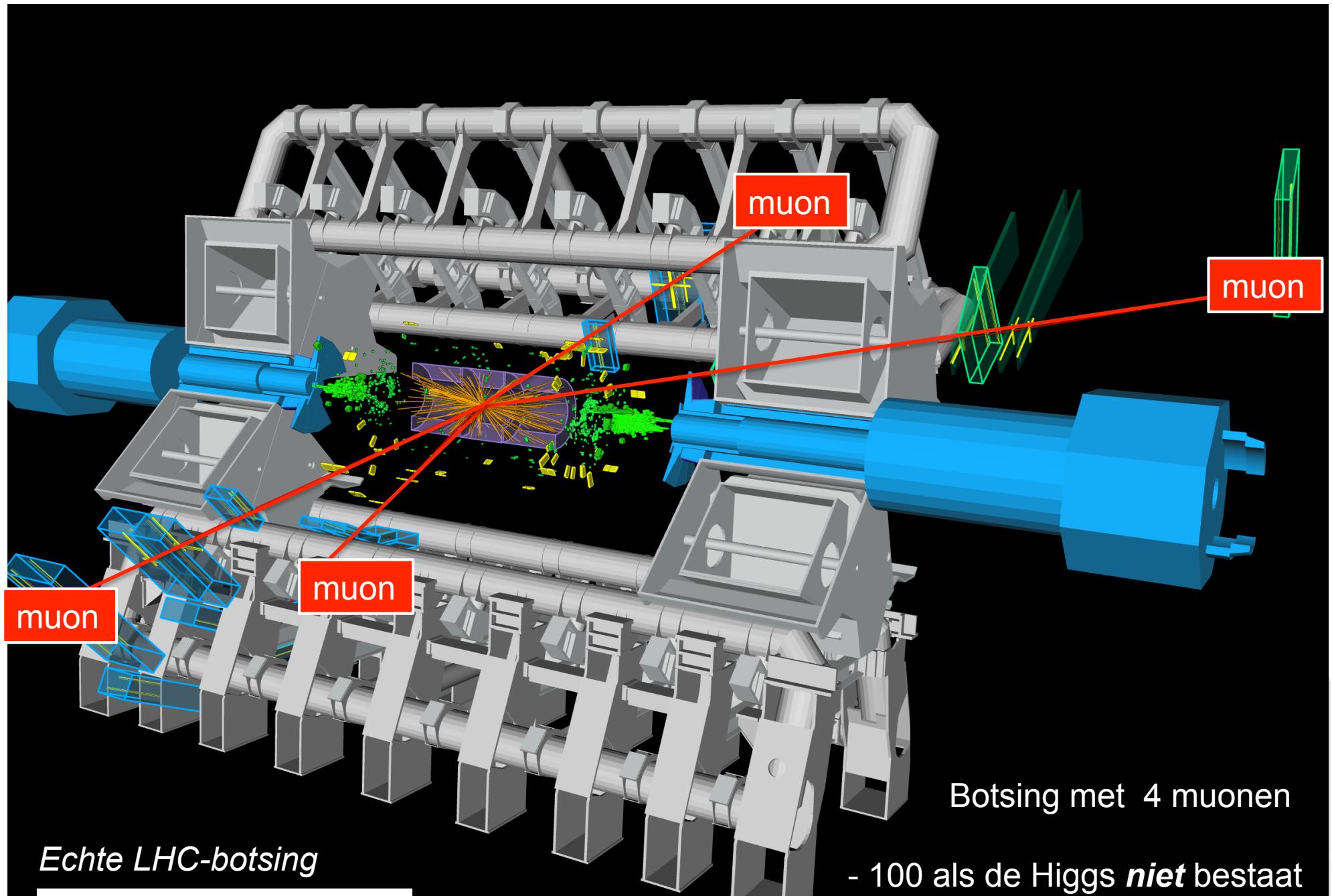
Lastig: Per seconde:

- 1 miljard botsingen (foto's)
- je mag er maar 200 bewaren (diskruimte)

Big data: 15 Peta Bytes en 100,00 CPU's

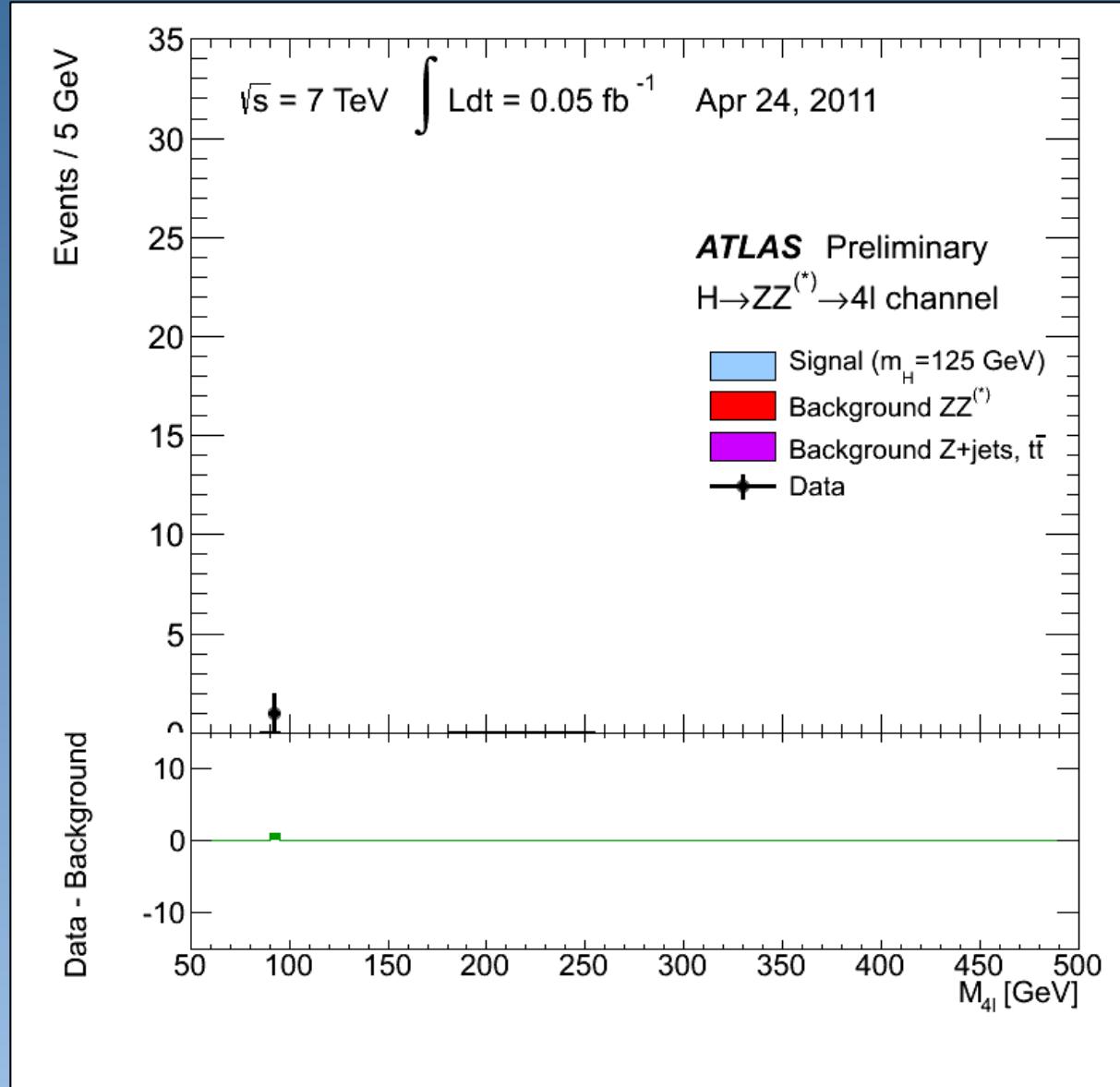
... het echte werk!





- 100 als de Higgs *niet* bestaat
- 200 als de Higgs *wel* bestaat

Kijk hier ($m_h = 125$ GeV)

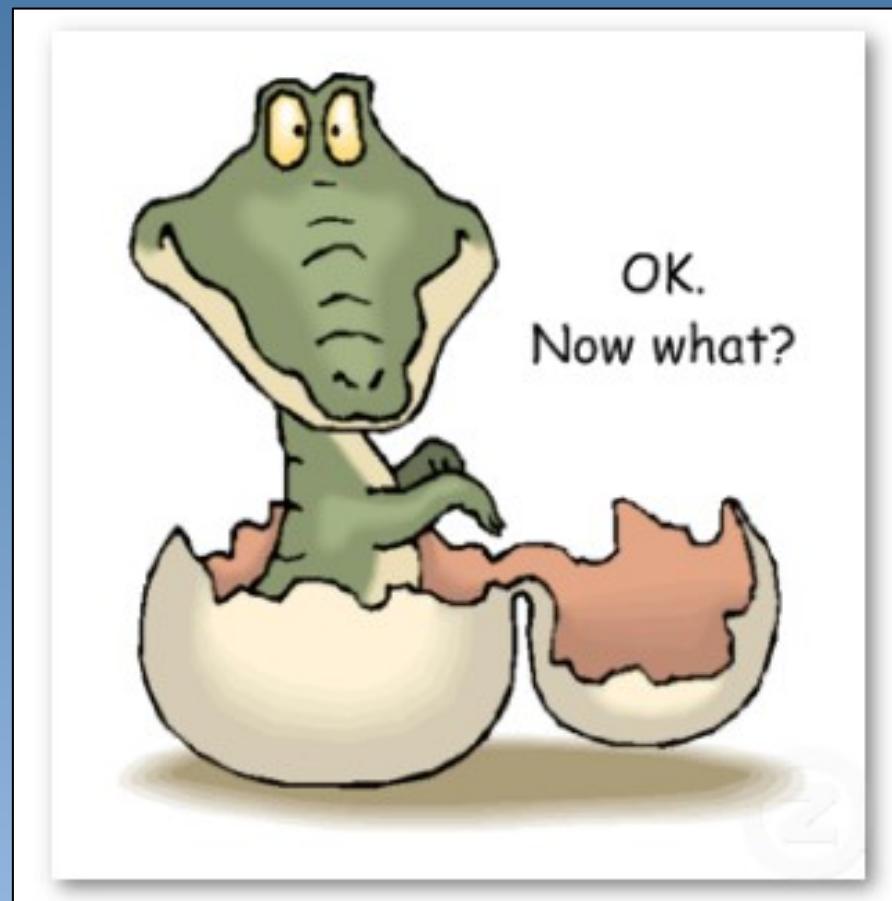


4 juli 2012: feest !!



Dinsdag 8 oktober 11:45 uur
bekendmaking Nobelprijs natuurkunde

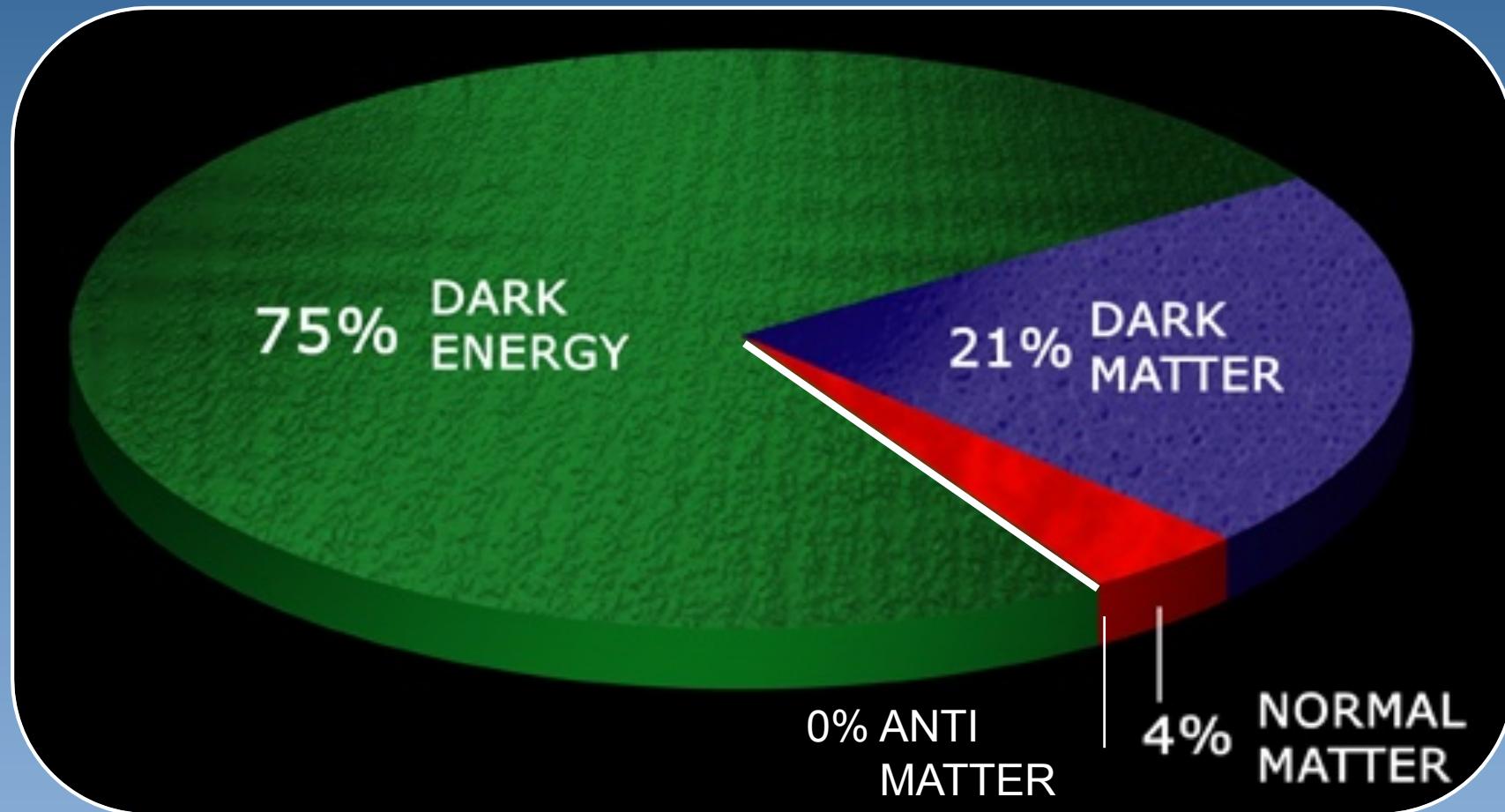
5 juli 2012: wat nu ?



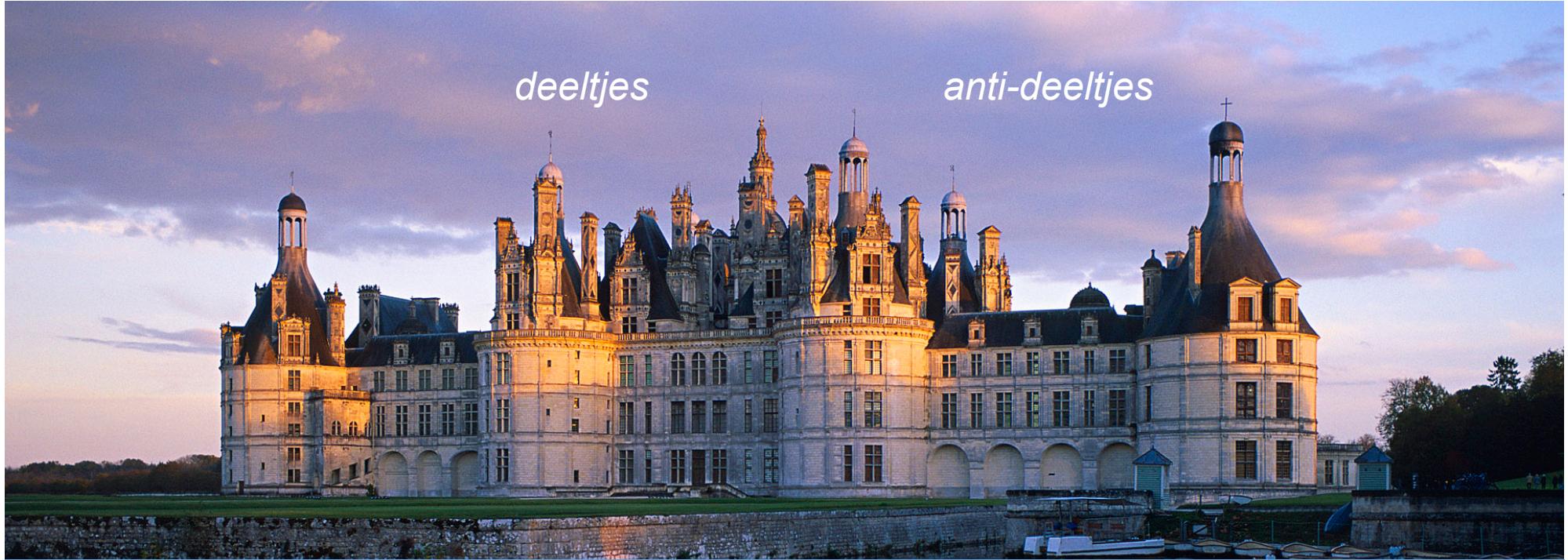
Wat snappen we nog niet:







Wat is de donkere materie ?

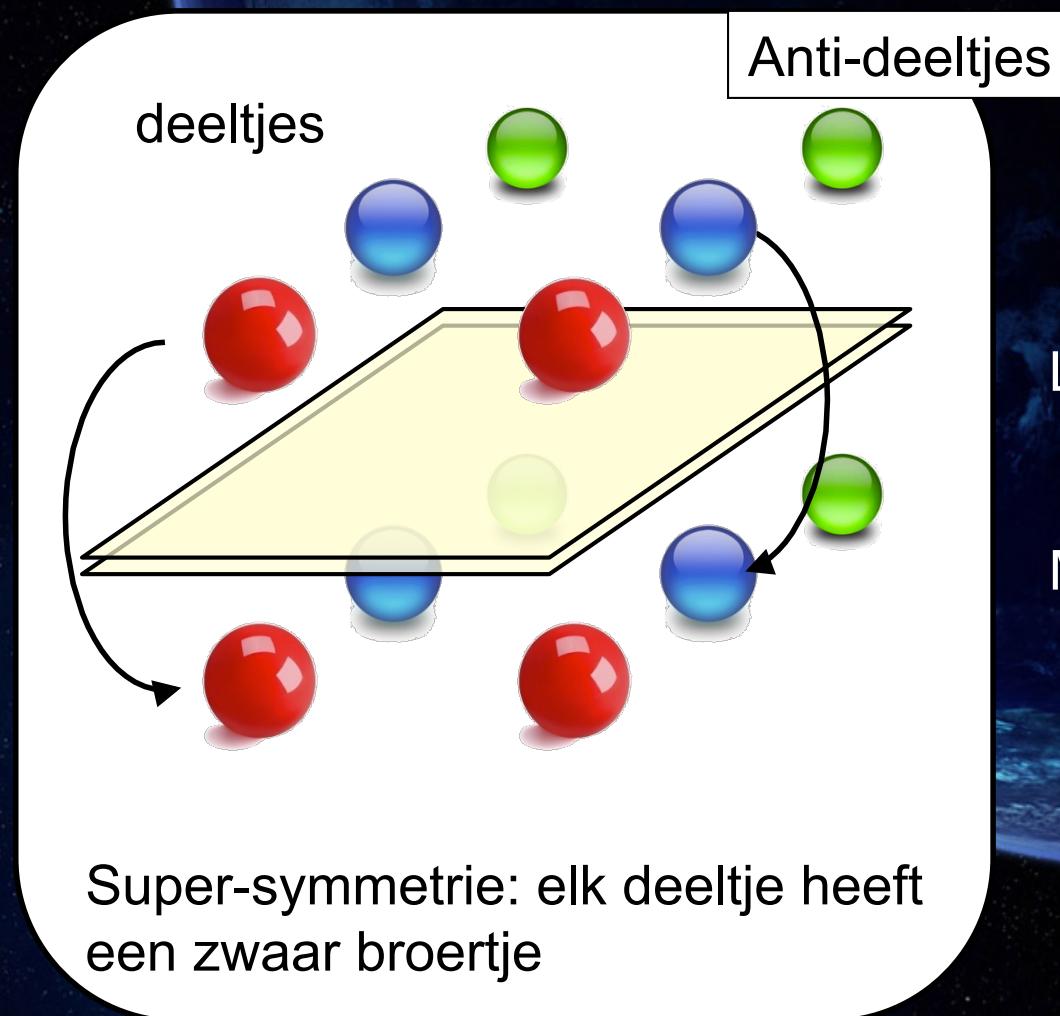


deeltjes

anti-deeltjes

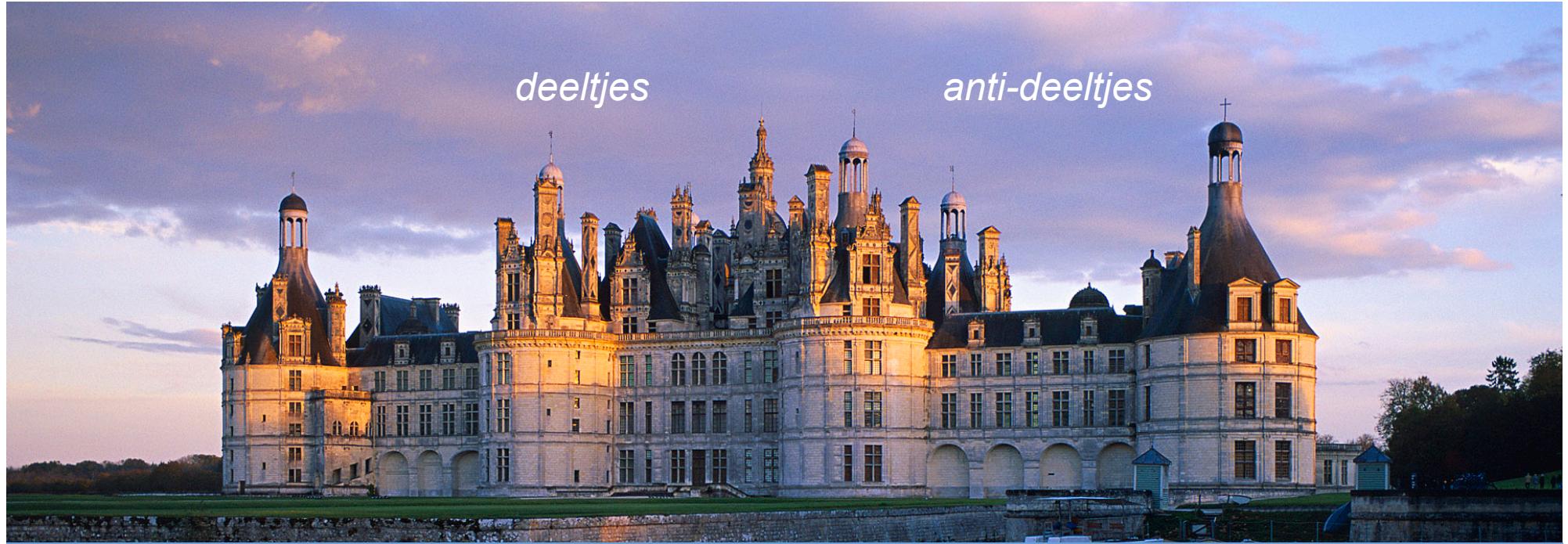
Donkere Materie zelf maken

Theoretisch idee: Supersymmetrie:



Lichtste is donkere materie deeltje ?

Maak ze bij de LHC !



Natuurkunde

De natuur zit vol vreemde spelregels

Het Higgs boson bestaat

Wat is donkere materie / energie ?

Waarom is de zwaartekracht zo raar ?

Spin-off

Slimme jonge mensen voor de industrie

Computing, sensors, versnellers

Medische apparatuur

avontuur

BACKUP

Relativiteitstheorie



Relativiteitstheorie

“structuur van tijd en ruimte”

De snelheid van een tennisbal en een lichtstraal gezien vanuit de ogen van een bewegende waarnemer



$$v = 299.792.458 \text{ [m/s]}$$



$$v = 20 \text{ [m/s]}$$



Snelheid ?

bal lichtstraal

20

299.792.458

Hoe verandert dit als de waarnemer zelf beweegt ?

relatieve snelheid
bal lichtstraal

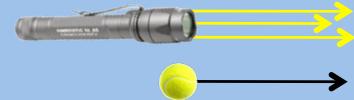


$$v = 299.792.458 \text{ [m/s]}$$
$$v = 20 \text{ [m/s]}$$

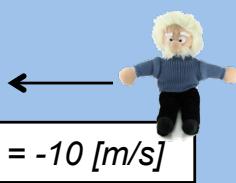


20

299.792.458



$$v = 299.792.458 \text{ [m/s]}$$
$$v = 20 \text{ [m/s]}$$



30

299.792.458

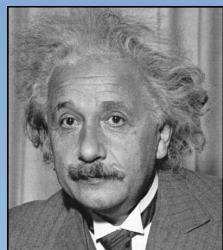


$$v = 299.792.458 \text{ [m/s]}$$
$$v = 20 \text{ [m/s]}$$



10

299.792.458



Idee (en observatie)

“de snelheid van het licht is constant voor elke waarnemer”

Albert Einstein Waarom: geen idee!

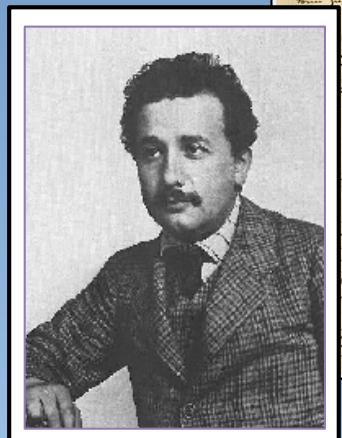
Consequenties:

Ruimte en tijd verbonden, gelijktijdigheid niet uniek gedefinieerd, ...

Speciale/algemene relativiteitstheorie:
tijd langzamer in een gravitatieveld

Relativiteitstheorie ... nuttig ?

Relativiteitstheorie ?



Es muss scharf hervorgehoben werden, dass diese Gleichungen
 die Maxima von Wellen aufweisen. Um es getrennt zu haben
 kann man Ringe, die aus einer oder zwei Wellen gebildet
 sind, ohne hohle oder grüne Abteilungen enthalten, und ein drittes
 hohles Ringe.
 Nur diese sind aus der Forderung der allgemeinen Relativität
 auf eine mathematische Theorie der Schwingungen in Beziehung
 mit den Gravitationsgleichungen (8) in einer Weise, die die Newtonsche
 Atmungszeit, in einer Weise, die die Zeitänderung des voraus
 Lernenden erfordert (nach Schrödinger's Theorie der Schwingungen
 überlappende) Periodizierung des Atoms legt, muss man eine
 Anzahl von den physikalischen Bedingungen des Theorems schwingen,
 um sie zu erfüllen.

8.5. Die Schwingungen der Ringe, welche hohlen sind für die gravitatorische
 Atmungszeit.

Um zu zeigen, dass die Totalwellenzahl eines Impulsen Kreiszeit
 entsprechend ist mit dem beginnenden, wie in folgenden Raum-Zeichen
 Form zu schreiben

$$\left. \begin{array}{l} S[\Gamma_{\mu\nu}] = 0 \\ R \stackrel{\text{def}}{=} T_{\mu\nu}^{\mu\nu} - \frac{T_{\mu\nu}}{2} \end{array} \right\} (74n)$$

und das Verhältnis an den Graden des determinaten logischen
 elementaren Schwingungsvermögens (die ist gleichzeitig gegeben, dass
 bei den Gleichungen (8) symmetrisch ist). In diesem Zweck können
 Funktionen des \vec{q} einzuführen ($\vec{q} = \frac{\partial}{\partial x^\mu}$). Dann ist gleichzeitig

$$S[\vec{q}] = T_{\mu\nu}^{\mu\nu} + 2T_{\mu}^{\mu\nu} - 2T_{\mu\nu}^{\mu\nu} (\vec{q} \cdot \vec{q})$$

$$= -T_{\mu\nu}^{\mu\nu} + 2T_{\mu}^{\mu\nu} + 2T_{\mu\nu}^{\mu\nu} (\vec{q} \cdot \vec{q})$$

oder

$$S[\vec{q}] = -\frac{1}{2} \left[q_\mu q^\mu \left(\frac{\partial^2}{\partial x_\mu^2} (q_\nu^2 + 2q_\mu^2 - 2q_\mu q_\nu) \right) \right]$$

die beiden letzten Formen sind die allgemeinen Wirkungsgrößen
 der Schwingungen, welche die logischen elementaren durch die des Bezeichnung
 aktions- und zeitabhangig (beweglich) ist. Durch Verhältnisse der Potenz
 können sie leicht erreicht werden für $S[\vec{q}] = 0$, welche
 logische der Potenz \vec{q} und $\vec{q} \cdot \vec{q}$ symmetrischen Graden \vec{q}
 werden. Sie bleibt also nur das erste Glied der weiteren
 zu bestimmen, welches nun mit Wirkungsgrößen verknüpft

ist, welche die nur von den Tieren q_μ und q_ν ($q_{\mu\nu}$) abhängen, werden
 die ist sehr einfach, nämlich gleich null, so kommt man wieder
 $q_\mu = 0$.



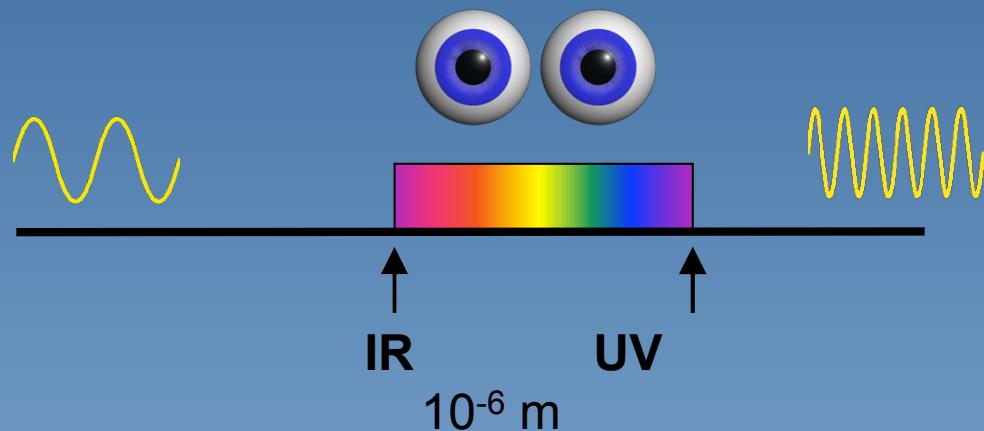
“Zonder relativiteitstheorie zit de GPS er 10 km per dag naast!”

Kijken zonder je ogen te gebruiken



1] Kijken met licht

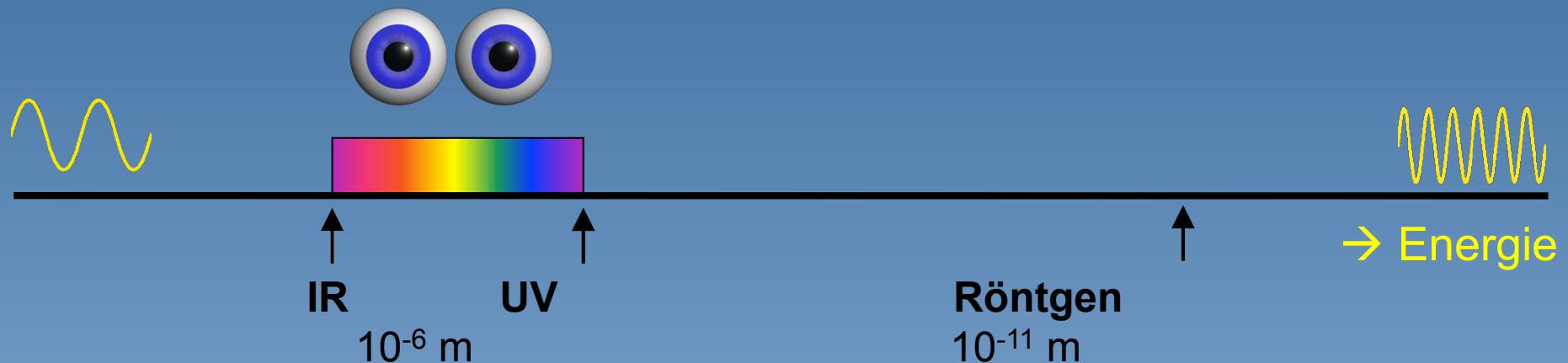
Licht verstrooit aan objecten die groter zijn dan zijn golflengte



Met je oog kan je niks zien
dat kleiner is dan 10^{-6} m

1] Kijken met licht

Licht verstrooit aan objecten die groter zijn dan zijn golflengte.



2] Kijken met deeltjes

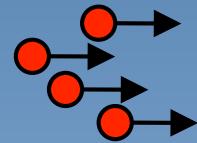
Quantummechanica: deeltjes zijn golven. Hoe meer energie hoe kleiner

→ microscoop voor zeer kleine afstanden:

- klein: elektronen microscoop
- kleinst: protonen micsoscoop

onbekend voorwerp achter een gordijn
... en 100 kogeltjes

bovenaanzicht

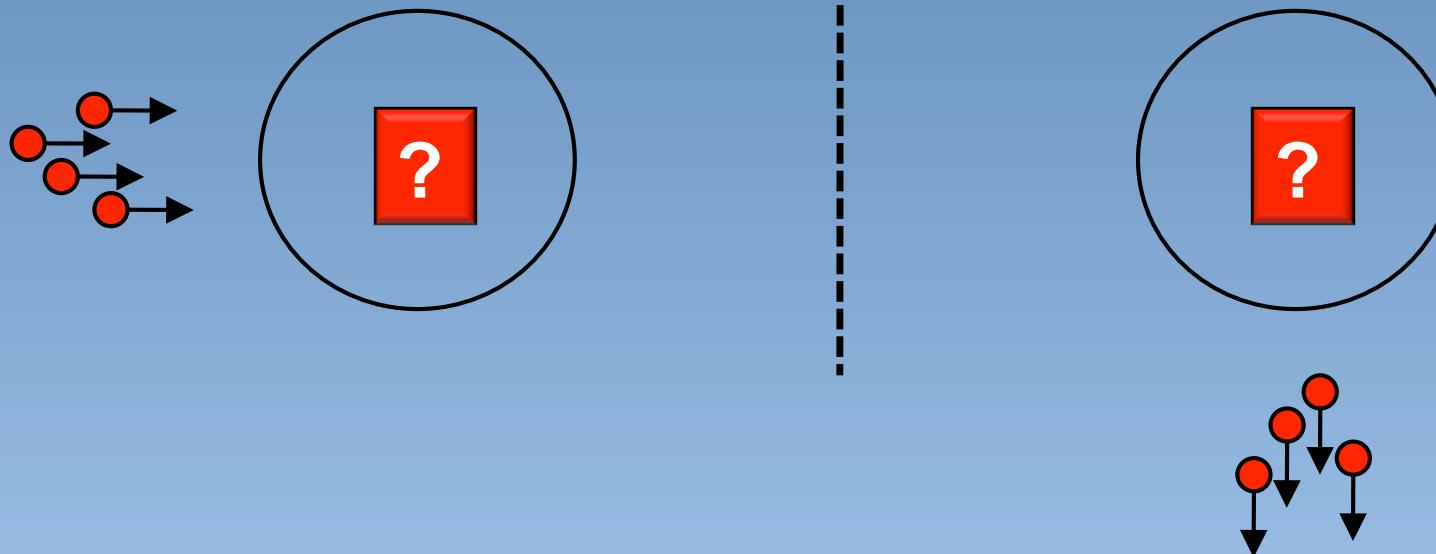


Hoe ketsen de kogels af ?

Situatie 1



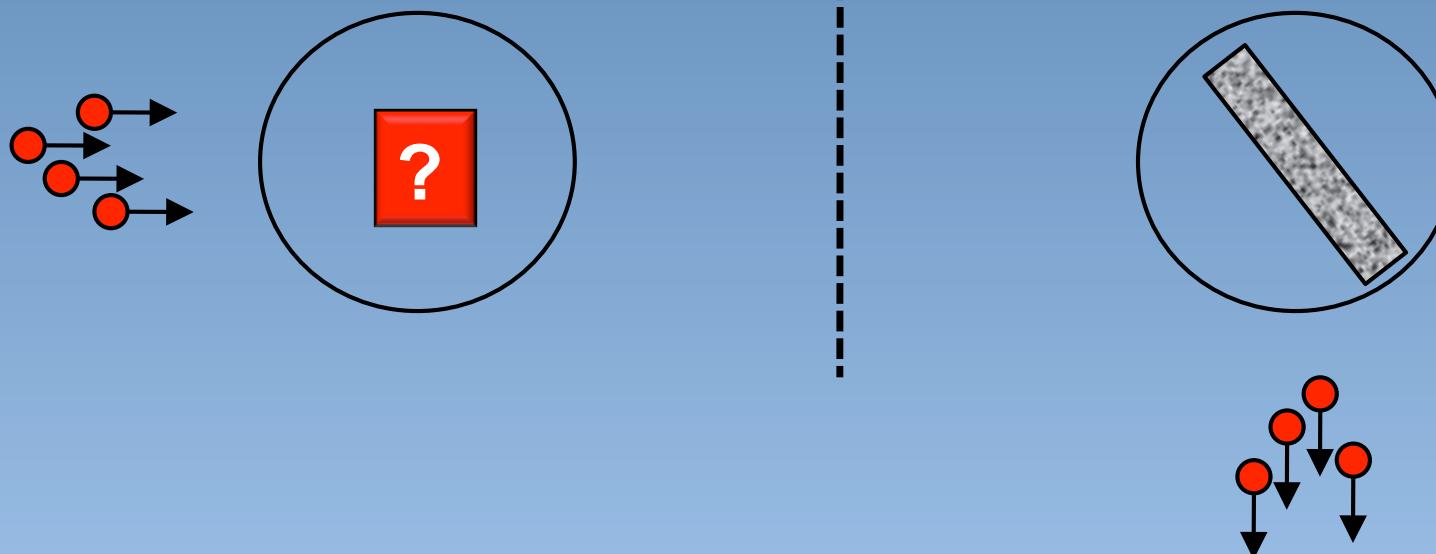
Situatie 2



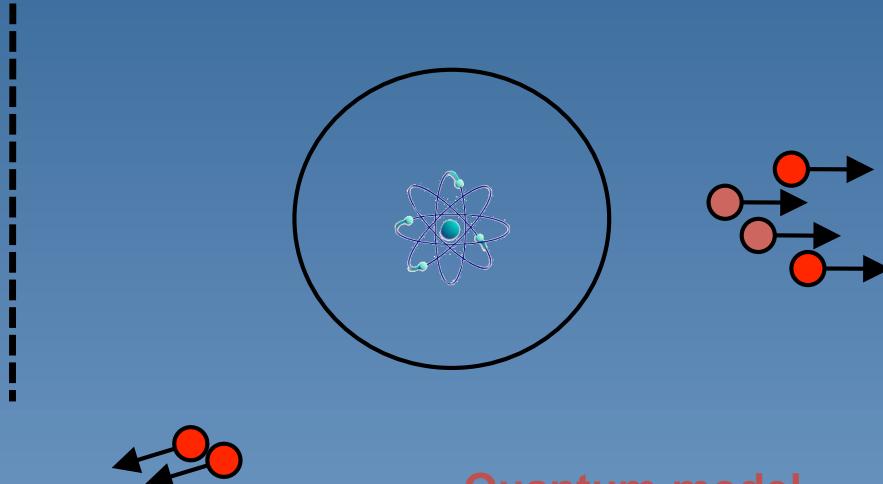
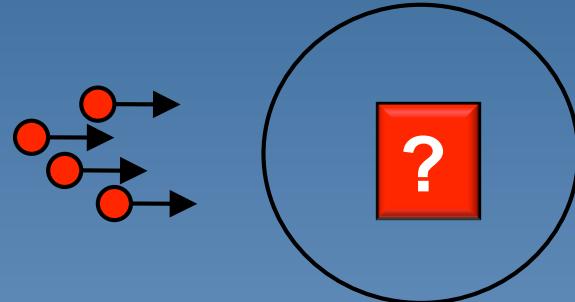
Situatie 1



Situatie 2



Experiment 3



Quantum-model
Einstein: $E=mc^2$

'Kijken' naar sub-atomaire结构uren:

- 1) Kleine kogels maken: Deeltjesversneller
- 2) Berekenen wat je verwacht voor een hypothese: Theorie
- 3) Afgeketste deeltjes kunnen bekijken: Detector

Distribution of All CERN Users by Nationality on 20 January 2010

