



THE DISCOVERY OF THE HIGGS BOSON

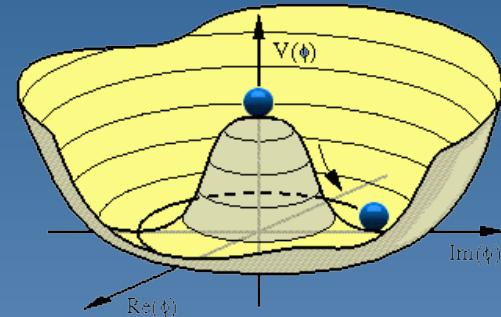
Ivo van Vulpen (Uva/Nikhef)

CERN in Geneva, Switzerland

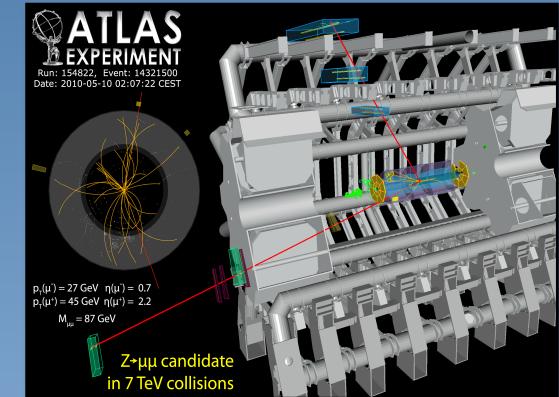


Things to remember

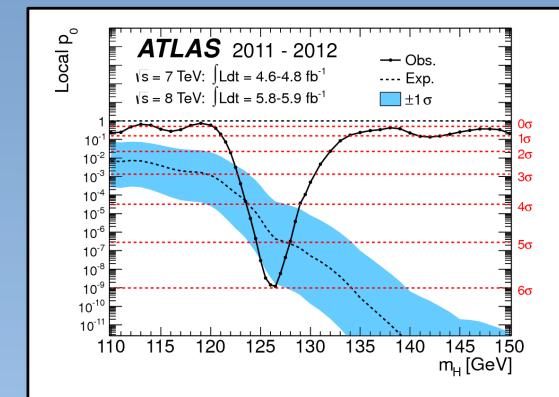
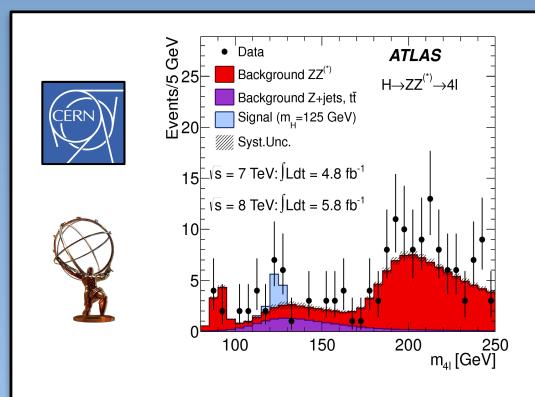
1) Higgs mechanism is at the heart of the Standard Model



2) LHC and ATLAS detector operating fine!

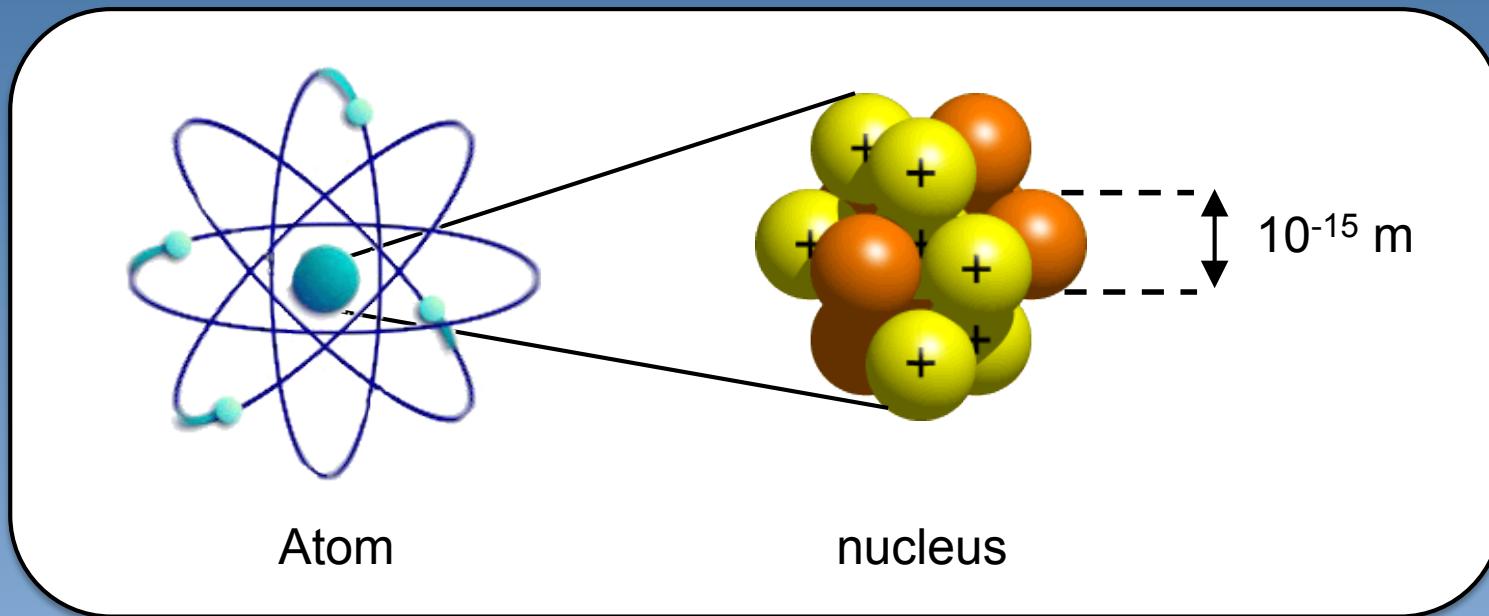


3) Discovery of the Higgs boson (interpretation)



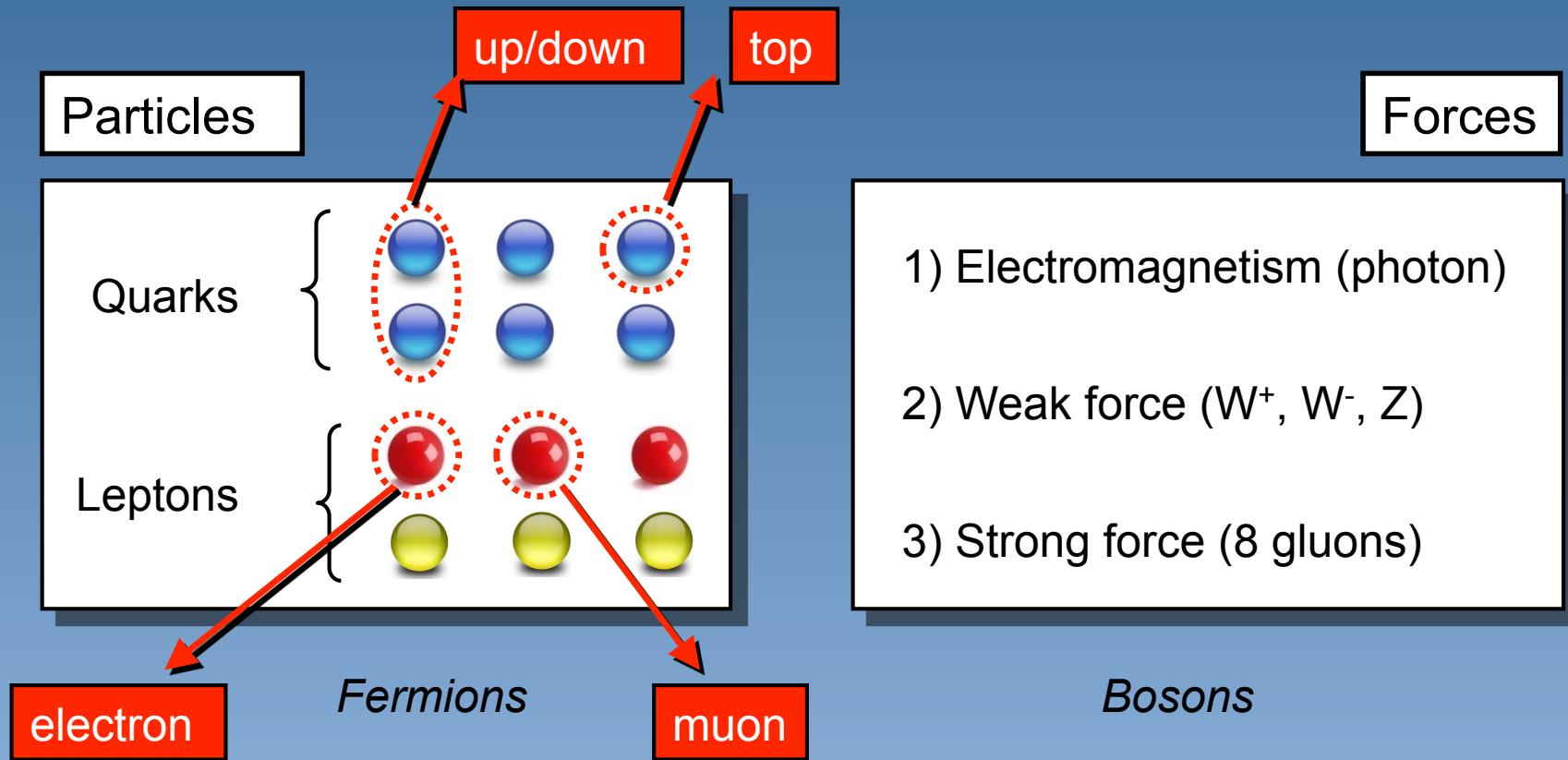
Particle Physics

Studies nature at distance scales $< 10^{-15} \text{ m}$



Standard Model:
Quantum field theory that describes phenomena down to 10^{-18} m

The Standard Model



The Standard Model

$$SU(2)_L \otimes U(1)_Y \otimes SU(3)_C$$

Weak iso-spin, hypercharge, colour

Standard Model based on local gauge invariance (symmetries)

Free electron

$$\mathcal{L}_e = i\bar{\psi}\gamma_\mu\partial^\mu\psi - m\bar{\psi}\psi$$

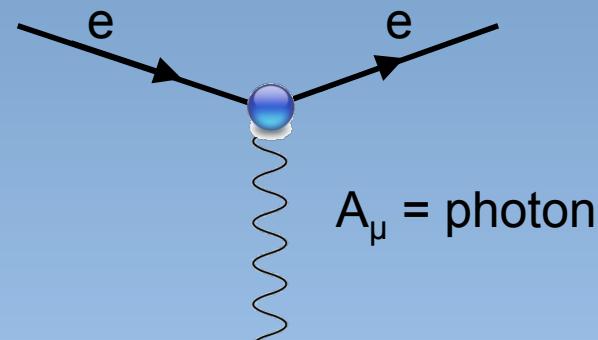
Extra symmetry requirement

Local gauge invariance

$$\psi(x) \rightarrow e^{i\alpha(x)}\psi(x)$$

Introduce covariant derivative (vector-field)

$$\partial_\mu \rightarrow D_\mu \equiv \partial_\mu - iqA_\mu$$



'automatic' force carriers and interactions'

In the Standard Model

$$SU(2)_L \otimes U(1)_Y \otimes SU(3)_C$$

Electroweak:
 W^+ , W^- , Z , γ

QCD:
8 gluons

What is missing in our Standard Model ?

GOOD THINGS:

- ➊ 'understand' origin of forces
- ➋ Excellent agreement data
- ➌ Connection EM / Weak force

NOT-SO-GOOD THINGS:

- ➊ No massive gauge bosons (W^\pm, Z)
- ➋ No massive fermions (all particles)
- ➌ Vector boson scattering diverges

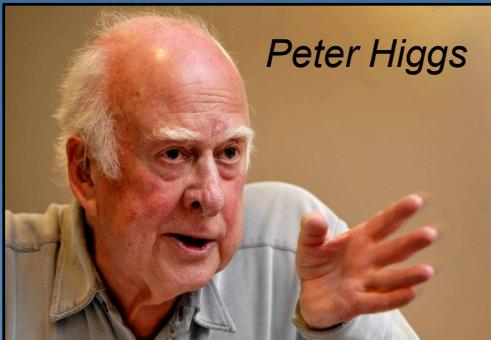


Higgs mechanism



Dark matter
Dark energy
Matter/anti-matter
Nature of gravity
Beginning of time

The Higgs mechanism



Massive gauge bosons in a local gauge invariant theory

$$SU(2)_L \otimes U(1)_Y \otimes SU(3)_C$$

There has to be a Higgs boson

BROKEN SYMMETRIES, MASSLESS PARTICLES AND GAUGE FIELDS

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Received 21 July 1964

Recently a number of people have discussed the Goldstone theorem^{1,2,3} that any solution of a Lorentz-invariant theory which violates an internal symmetry operation of that theory must contain a massless scalar particle. Kibble and Lee² showed that this theorem does not necessarily apply in non-relativistic theories and argued that their considerations would apply equally well to Lorentz-invariant field theories. Gell-Mann⁴, however,

gave a proof that the failure of the Goldstone theorem in the non-relativistic case is of a type which cannot occur when Lorentz-invariance is imposed on a theory. The purpose of this note is to show that Gell-Mann's argument fails for an important class of field theories, that in which the conserved currents are coupled to gauge fields.

Following the procedure used by Gell-Mann⁴, let us consider a theory of two hermitian scalar fields

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$\psi_1(0, \vec{x}, t)$ which is invariant under the phase transformation

$$\begin{aligned} \psi_1 &\rightarrow e^{i\theta} \psi_1 \quad \theta = \text{constant} \\ \psi_2 &\rightarrow e^{-i\theta} \psi_2 \quad \theta = \text{constant} \end{aligned} \quad (1)$$

Then there is a conserved current j_μ such that

$$(j^\mu \partial_\mu) j_\mu = \psi_2 \partial_\mu \psi_1. \quad (2)$$

We assume that the Lagrangian in such that symmetry is broken by the non-vanishing of the vacuum expectation value of ψ_2 . Goldstone's theorem is proved by showing that the Fourier transform of $(j^\mu \partial_\mu) j_\mu$ contains a term

$i(\vec{k}_2) k_\mu j_\mu / (2\pi)^3$, where k_μ is the momentum, as a consequence of Lorentz-invariance, the conservation law and eq. (1).

Kibble and Lee² avoided this result in the non-relativistic case by showing that the most general form of this Fourier transform is now, in Gell-Mann's notation,

$F(T) = k_{12} \psi_2 (T^2, m) + i k_2 \psi_2 (T^2, m) + C \delta_T \psi_2 (T^2, m)$, where ψ_{12} , which may be taken as $(1, 0, 0, 0)$, depicts not a special Lorentz frame. The conservation law then reduces eq. (2) to the less general form

$$\begin{aligned} F(T) &= k_{12} \psi_2 (T^2, m) + [k_2^2 \psi_{12} + k_{12} (m)] \psi_2 (T^2, m) \\ &\quad + C \delta_T \psi_2 (T^2, m). \end{aligned} \quad (3)$$

It turns out, on applying eq. (3), that all three terms in eq. (3) can contribute to ψ_{12} . Thus the Goldstone theorem fails if $k_{12} = 0$, which is possible only if the other term is zero. Gell-Mann's remark that no special timelike vector n_μ is available in a Lorentz-invariant theory appears to rule out this possibility in such a theory.

There is however a class of non-relativistic field theories in which a vector n_μ does indeed play a part. This is the class of gauge theories, where an auxiliary and timelike vector n_μ must be intro-

duced in order to define a radiation gauge in which the vector gauge fields are well-defined operators. Such theories are nevertheless Lorentz-covariant, as has been shown by Stueckelberg⁵. (This has, of course, long been known of the simplest such theory, quantum electrodynamics.) There seems to be no reason why the vector n_μ should not appear in the Fourier transform under consideration.

It is characteristic of gauge theories that the conservation laws hold in the strong sense, as a consequence of field equations of the form

$$\partial_\mu A_\mu = \partial_\mu V_\mu = \partial_\mu A'_\mu. \quad (4)$$

Except in the case of abelian gauge theories, the fields A_μ , V_μ , A'_μ are not simply the gauge field variables A_μ , V_μ , but contain additional terms with combinations of the structure constants of the group as coefficients. How the structure of the Fourier transforms of $(\psi_{12}(T), \psi_2(T))$ must be given by eq. (3). Applying eq. (3) to this combination gives us as the Fourier transform of $(\psi_{12}(T), \psi_2(T))$ the single term $[k_{12}^2 - k_{12} (m)] \psi_2 (T^2, m)$. We have thus exercised both Goldstone's zero-mass boson and the "gauge" state ($k_{12} = 0$) proposed by Kibble and Lee.

In a subsequent note it will be shown, by considering some classical field theories which display broken symmetries, that the introduction of gauge fields may be expected to produce qualitative changes in the nature of the particles described by such theories after quantization.

References

- 1) J. Goldstone, Nuovo Cimento 33 (1960) 134.
- 2) J. Goldstone, A. Salam and S. Weinberg, Phys. Rev. 187 (1960) 1918.
- 3) J. Goldstone and B. W. Lee, Phys. Rev. Letters 13 (1964) 166.
- 4) R. Gell-Mann, Phys. Rev. Letters 13 (1964) 115.
- 5) J. Stueckelberg, Phys. Rev. 117 (1960) 208.

- September 1964 -

(In)famous Higgs boson



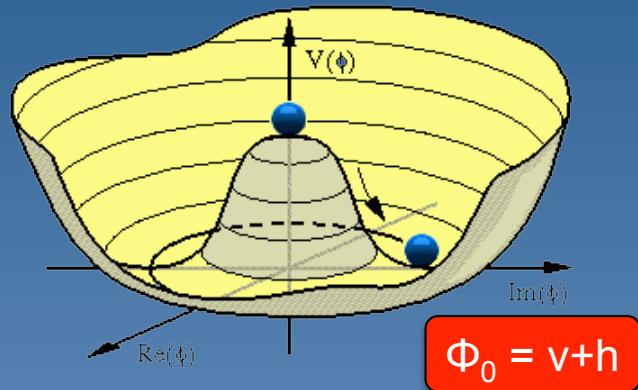
The Higgs boson



Paris Hilton

Being famous is not the same as being important

The Higgs mechanism



- 1) Add iso-spin doublet Φ (4 d.o.f.), with $Y_\Phi = +1$
- 2) Potential: $V(\Phi) = \mu^2\Phi^2 + \lambda\Phi^4$, with $\mu^2 < 0$

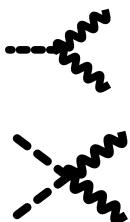
Gauge bosons

$(V = W^+, W^-, Z, \gamma)$

$$L(\phi) = \underbrace{(D_\mu \phi)(D^\mu \phi)}_{\propto \vec{V}^2} - V(\phi)$$

$$\propto \vec{V}^2 \quad \text{mass } V$$

$$\propto h \vec{V}^2$$



$$\propto h^2 \vec{V}^2$$

+ scalar particle

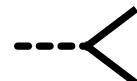
Fermions

$(\psi = e, \mu, \dots)$

$$L(\psi) = \underbrace{\bar{\psi} \phi \psi}_{\propto m_\psi \bar{\psi} \psi} + \dots$$

$$\propto m_\psi \bar{\psi} \psi \quad \text{mass } \psi$$

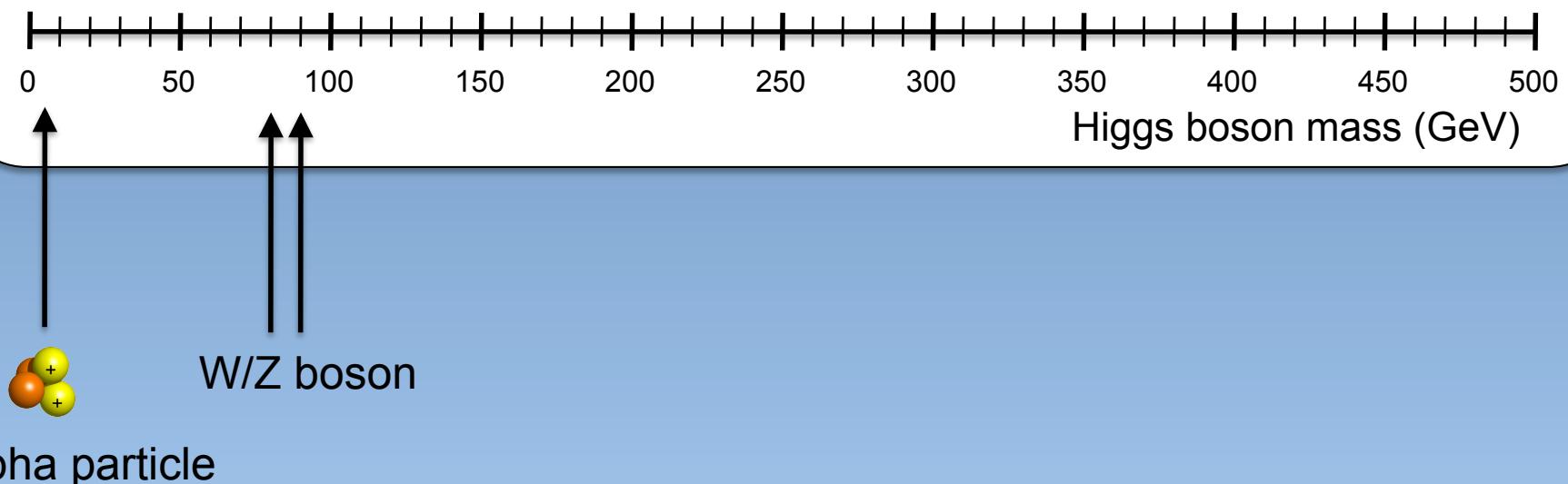
$$\propto \left(\frac{m_\psi}{v} \right) \bar{\psi} h \psi$$



“the Higgs couples to mass”

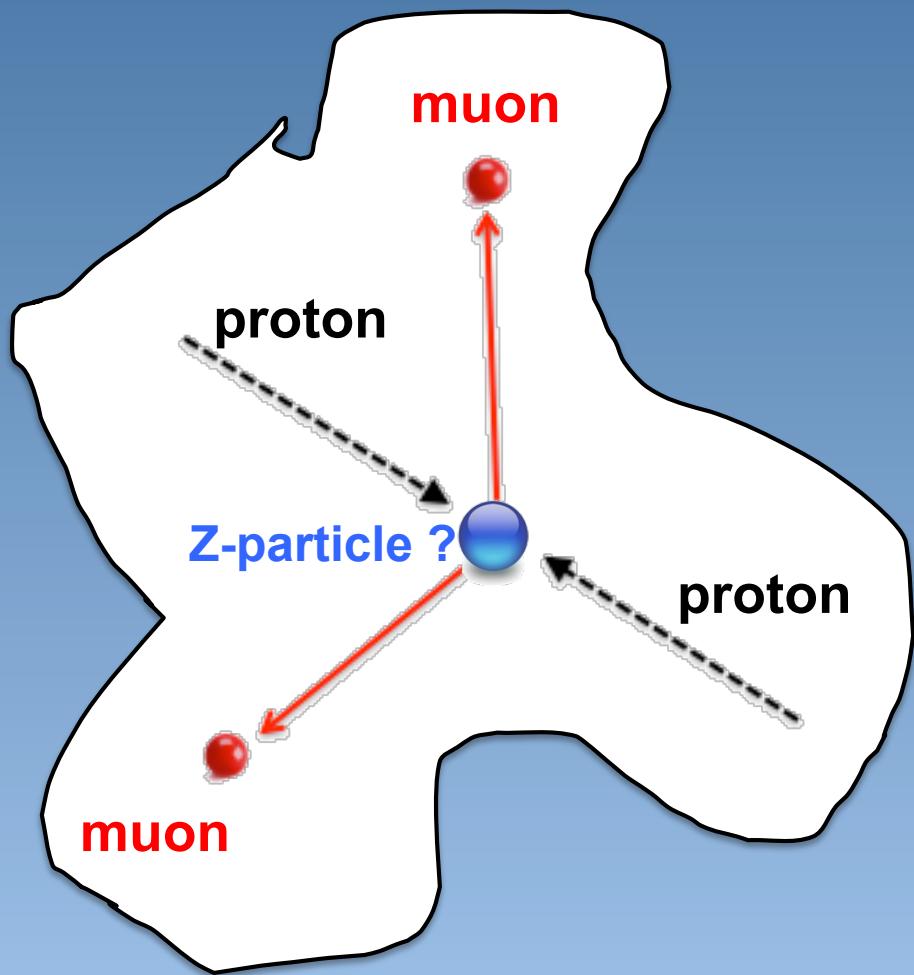
The mass of the Higgs boson

*Properties of the Higgs boson are known ... as a function of it's **unknown** mass*

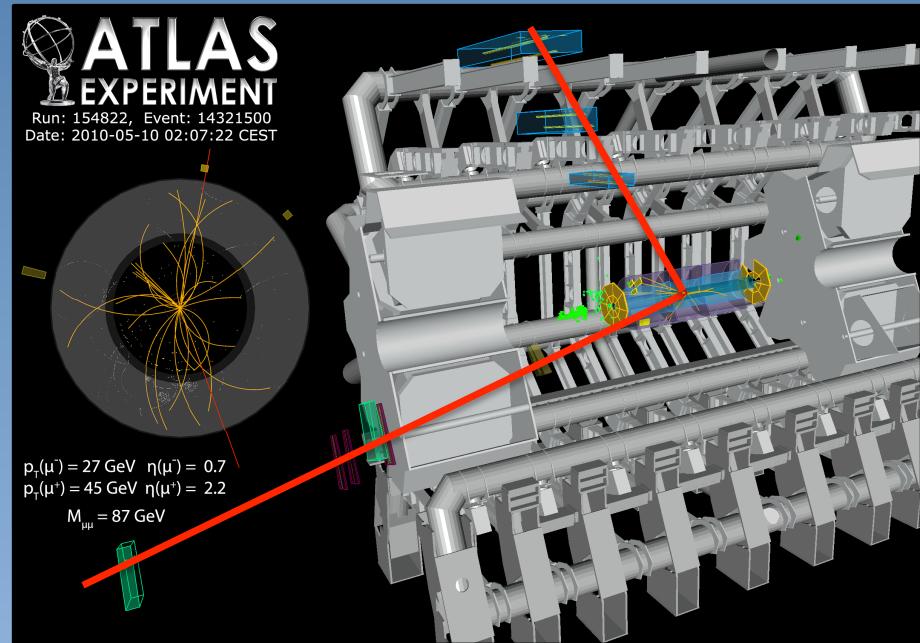


How do you discover a Z particle ?

Deviations from the prediction without a Z boson



2-muons as seen in the ATLAS detector

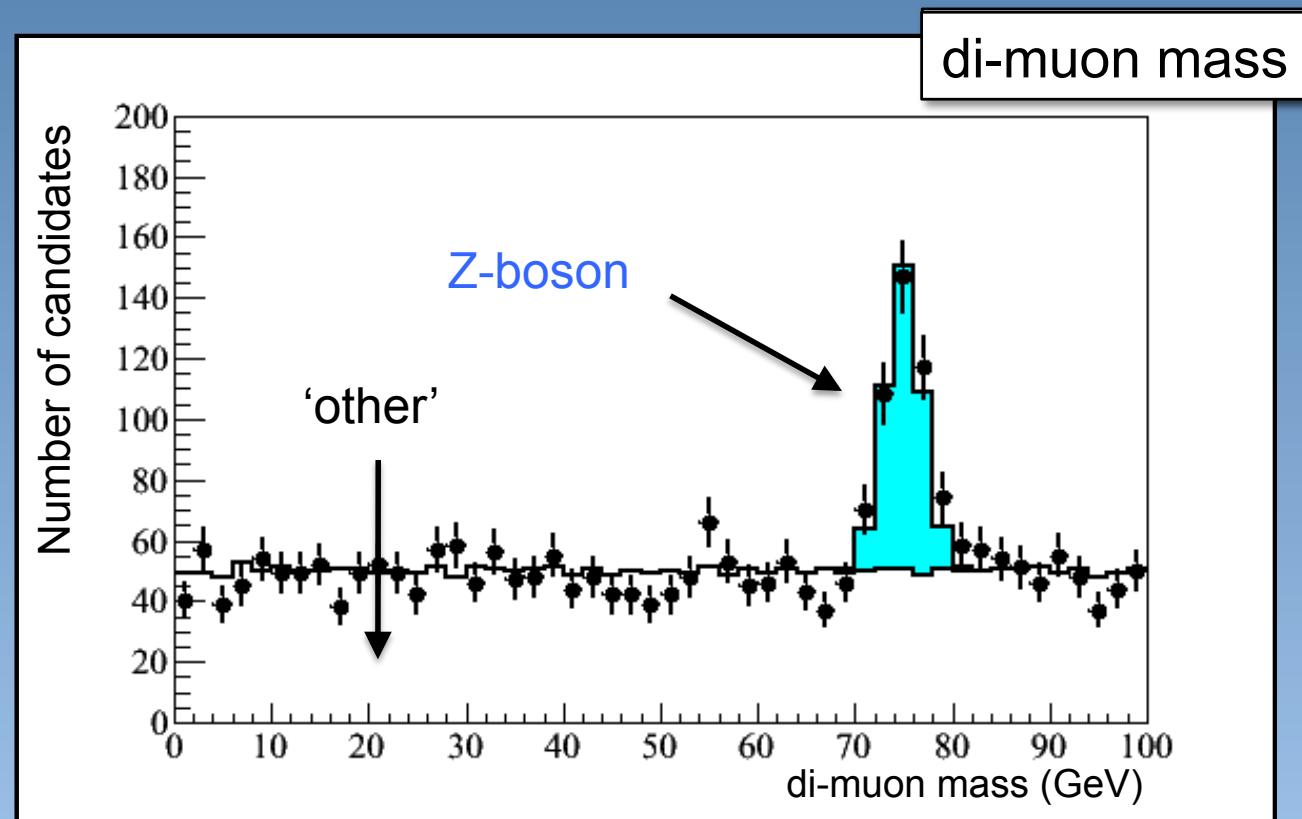


What is the fingerprint of the Z particle:
→ the di-muon mass

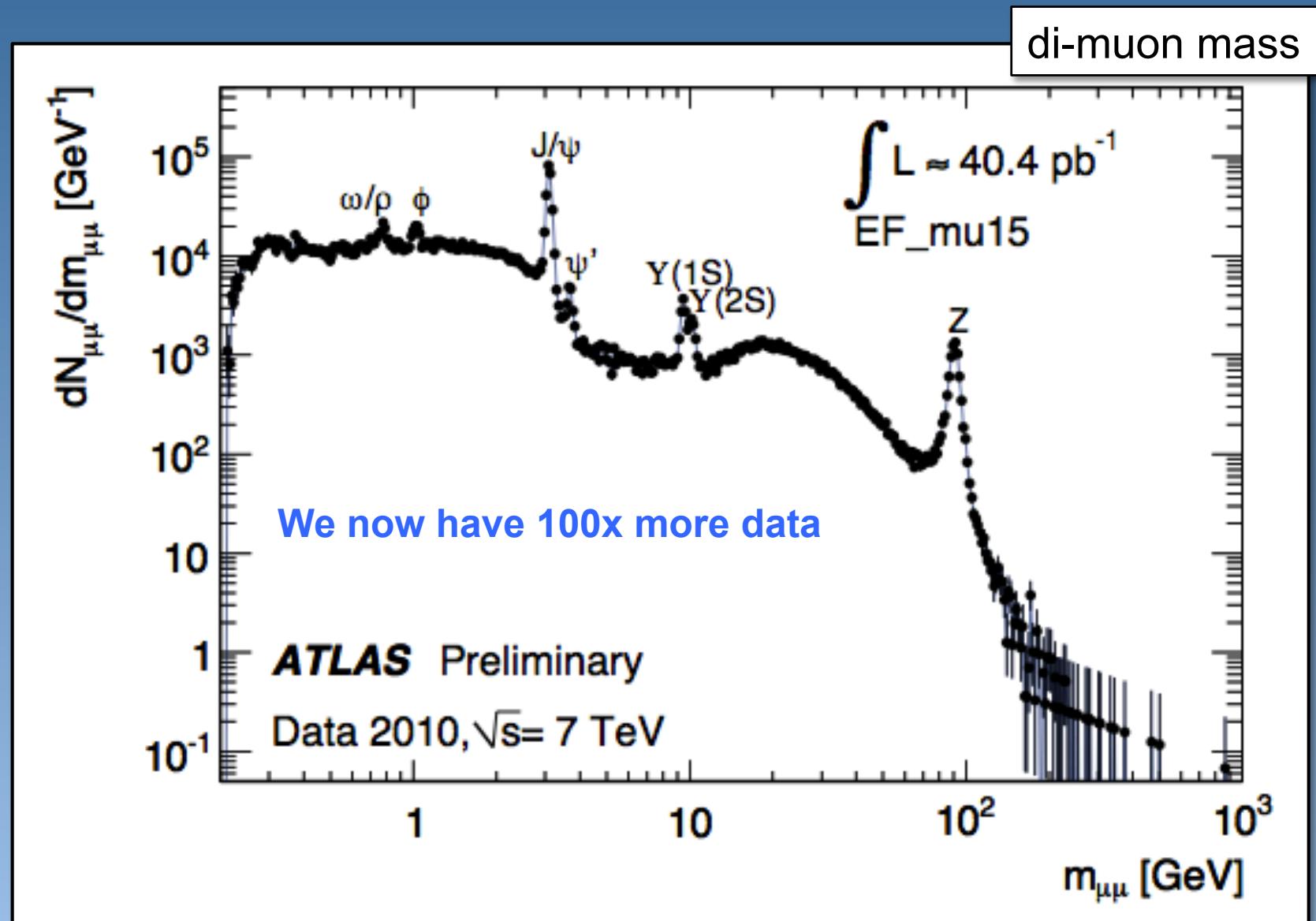


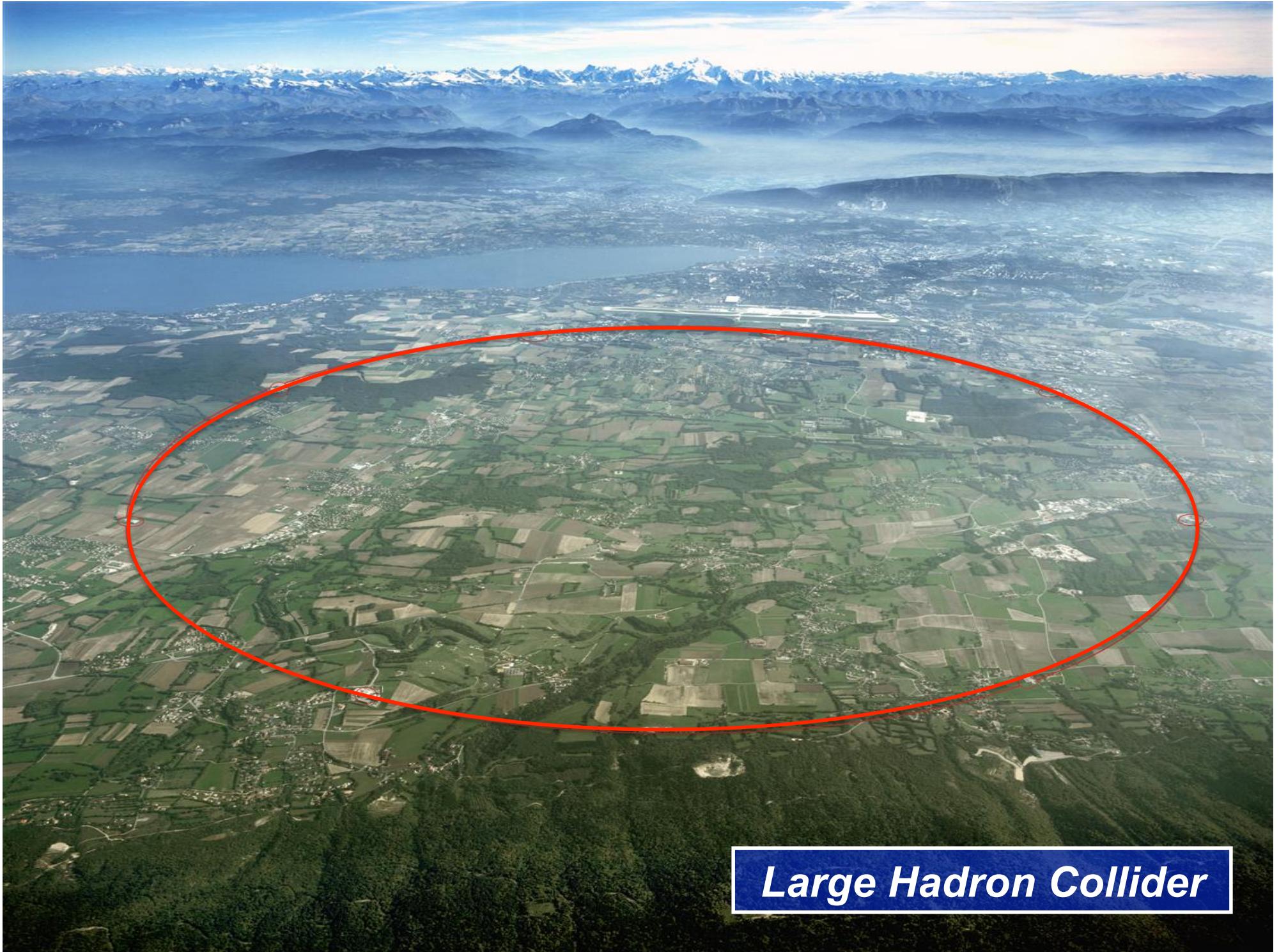
'other': no preference/structure

Z-boson: Breit-Wigner \otimes resolution



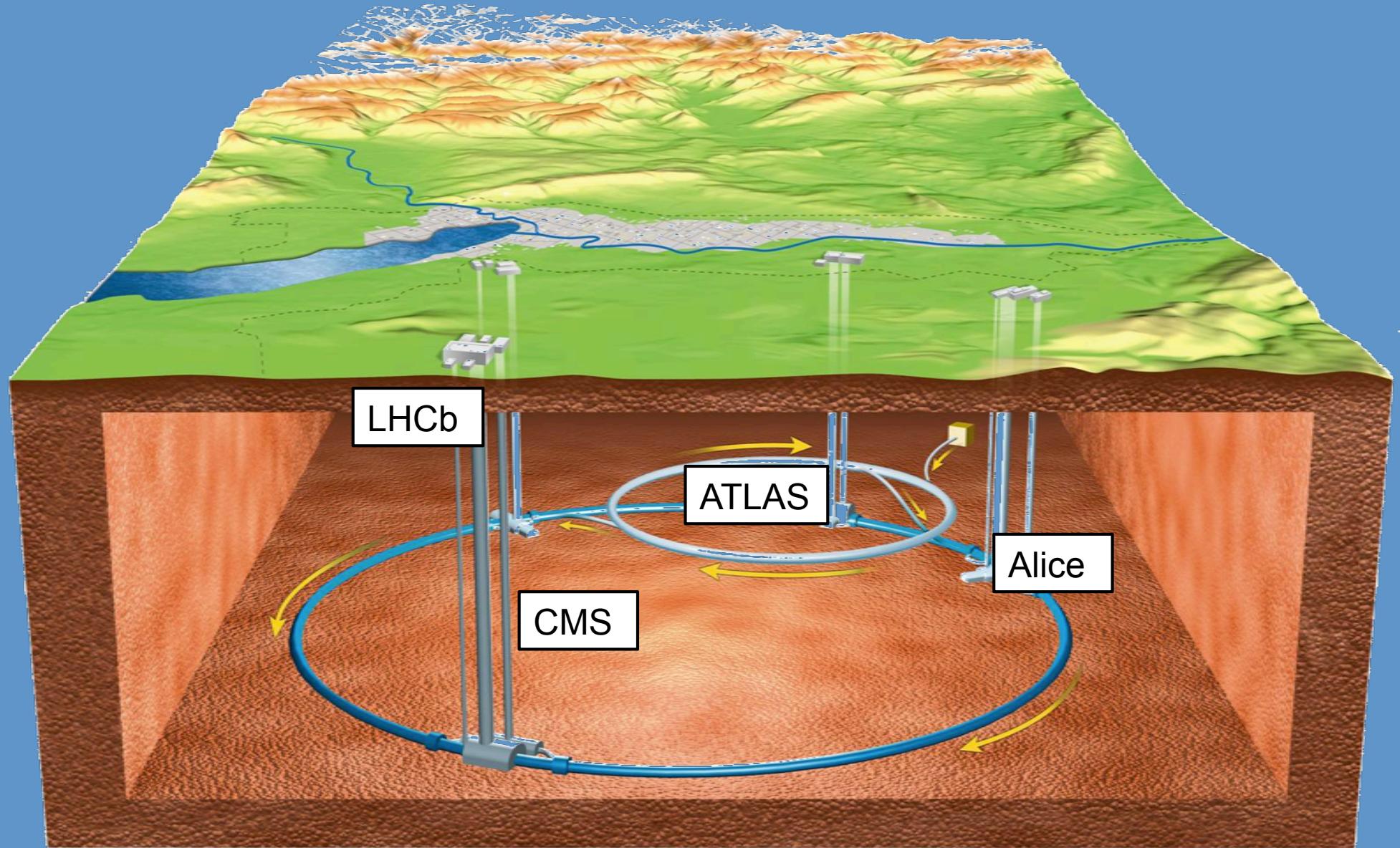
Di-muon resonances seen at the LHC

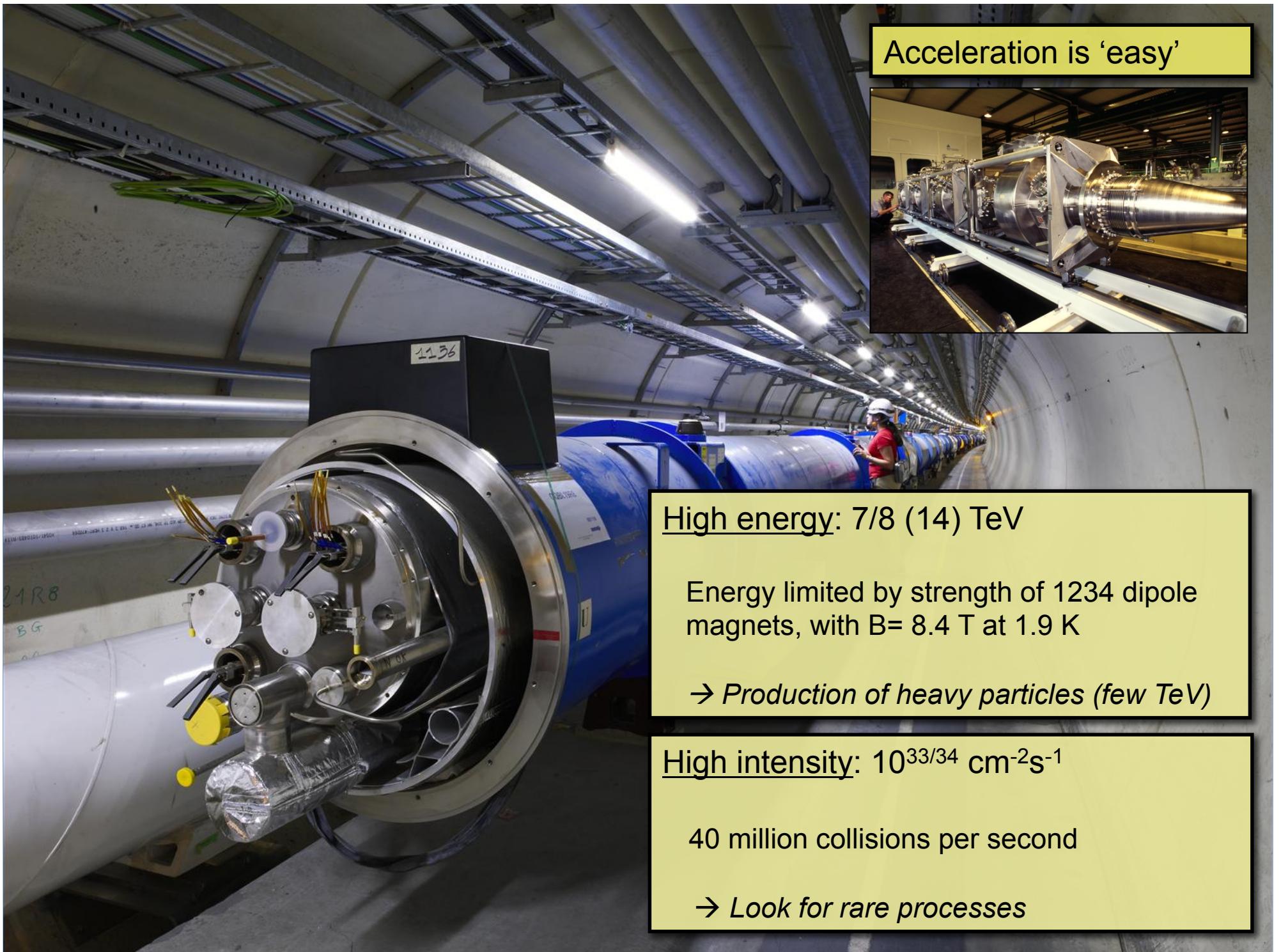




Large Hadron Collider

The Large hadron collider





Acceleration is 'easy'

High energy: 7/8 (14) TeV

Energy limited by strength of 1234 dipole magnets, with $B = 8.4 \text{ T}$ at 1.9 K

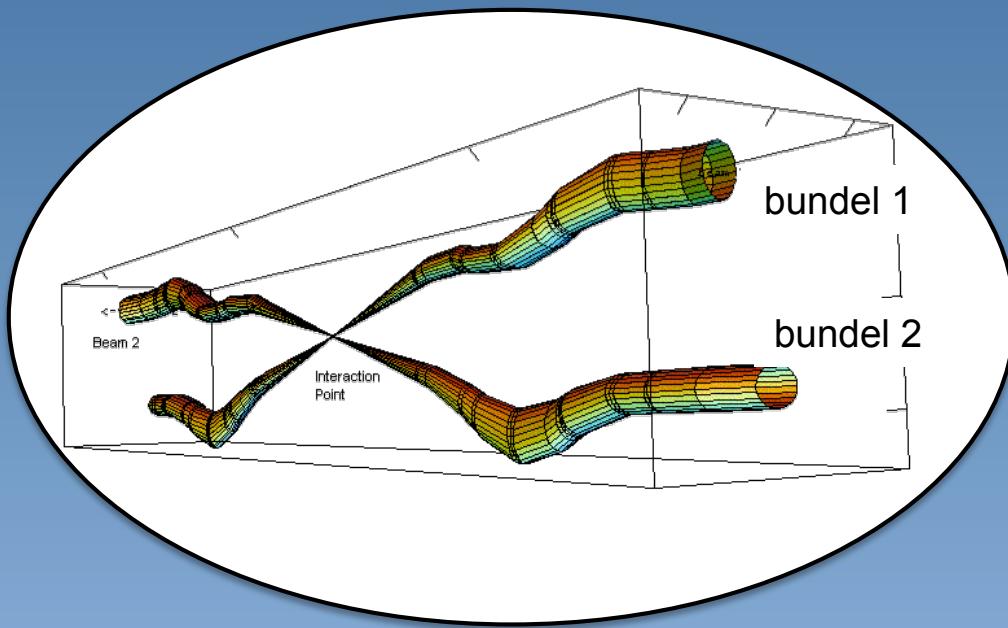
→ *Production of heavy particles (few TeV)*

High intensity: $10^{33/34} \text{ cm}^{-2}\text{s}^{-1}$

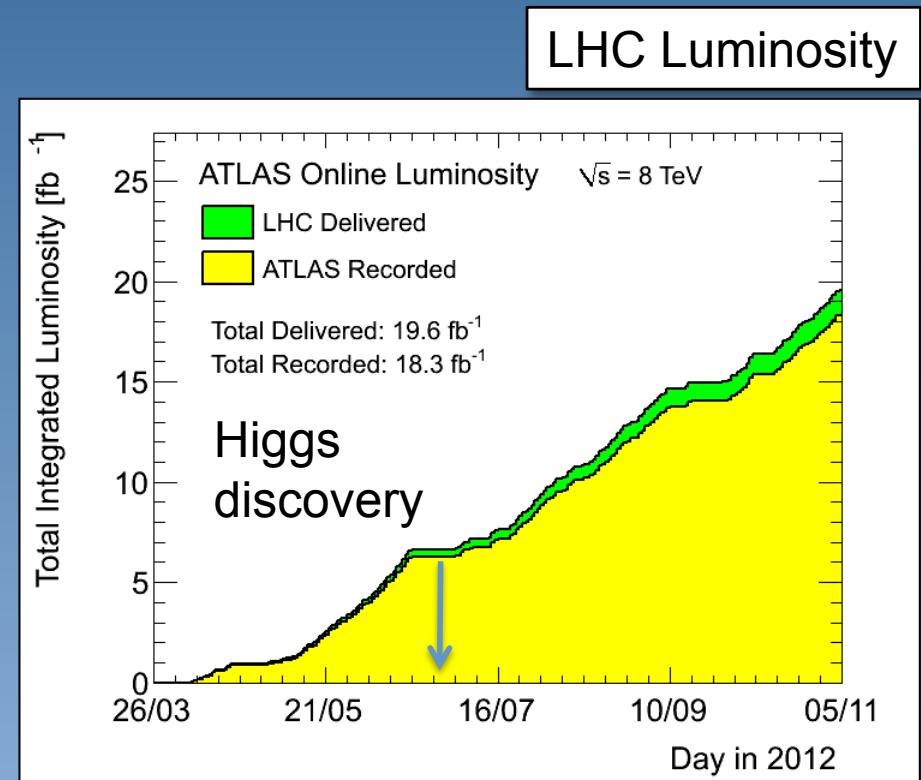
40 million collisions per second

→ *Look for rare processes*

LHC data-set 2011+ 2012



40 million collisions per second
Only 200 Hz to tape

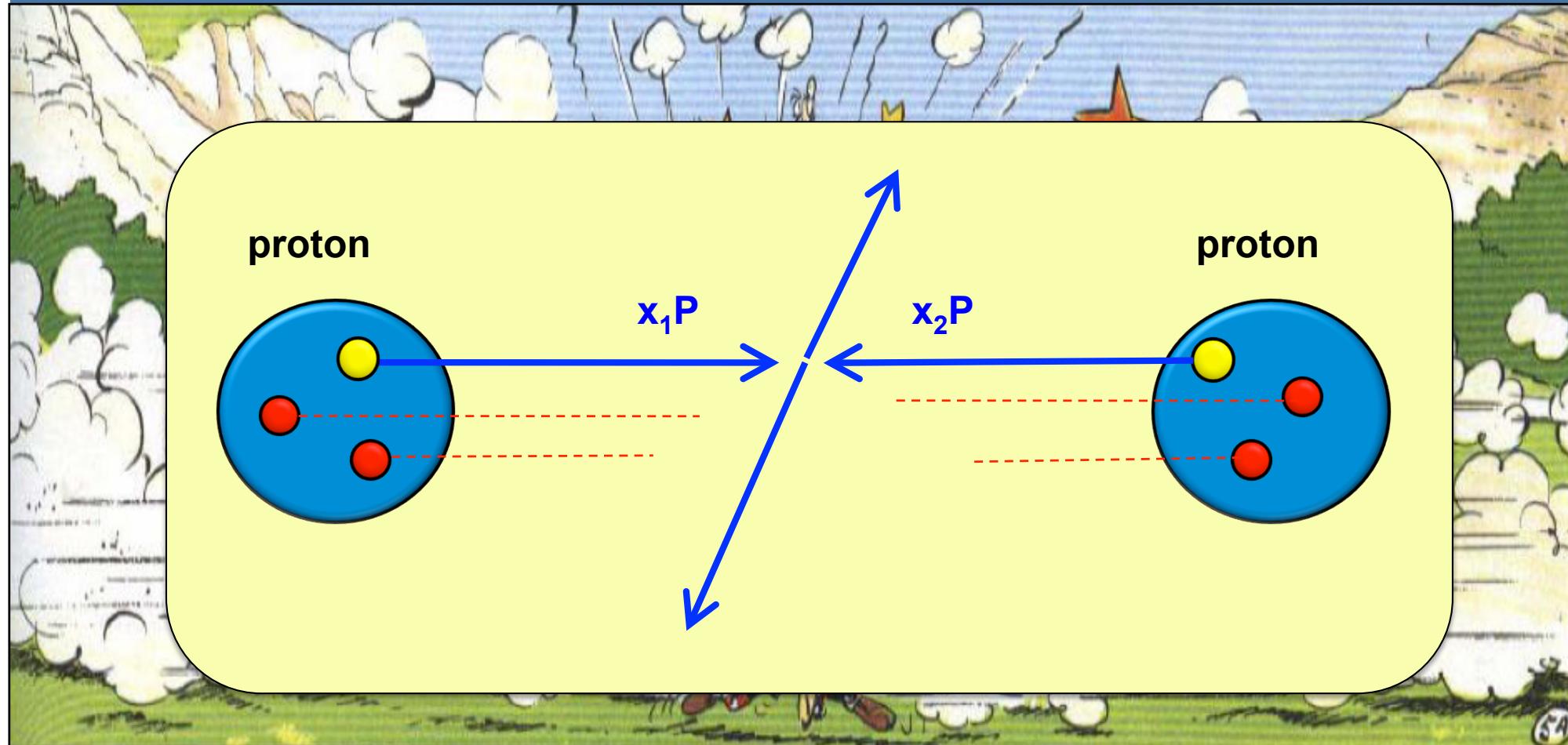


ATLAS (2011): 4.8 fb^{-1} at $\sqrt{s}=7 \text{ TeV}$
ATLAS (2012): 19.6 fb^{-1} at $\sqrt{s}=8 \text{ TeV}$

Note: almost all results in this talk are
from data up to discovery

Phenomenology of proton-proton collisions

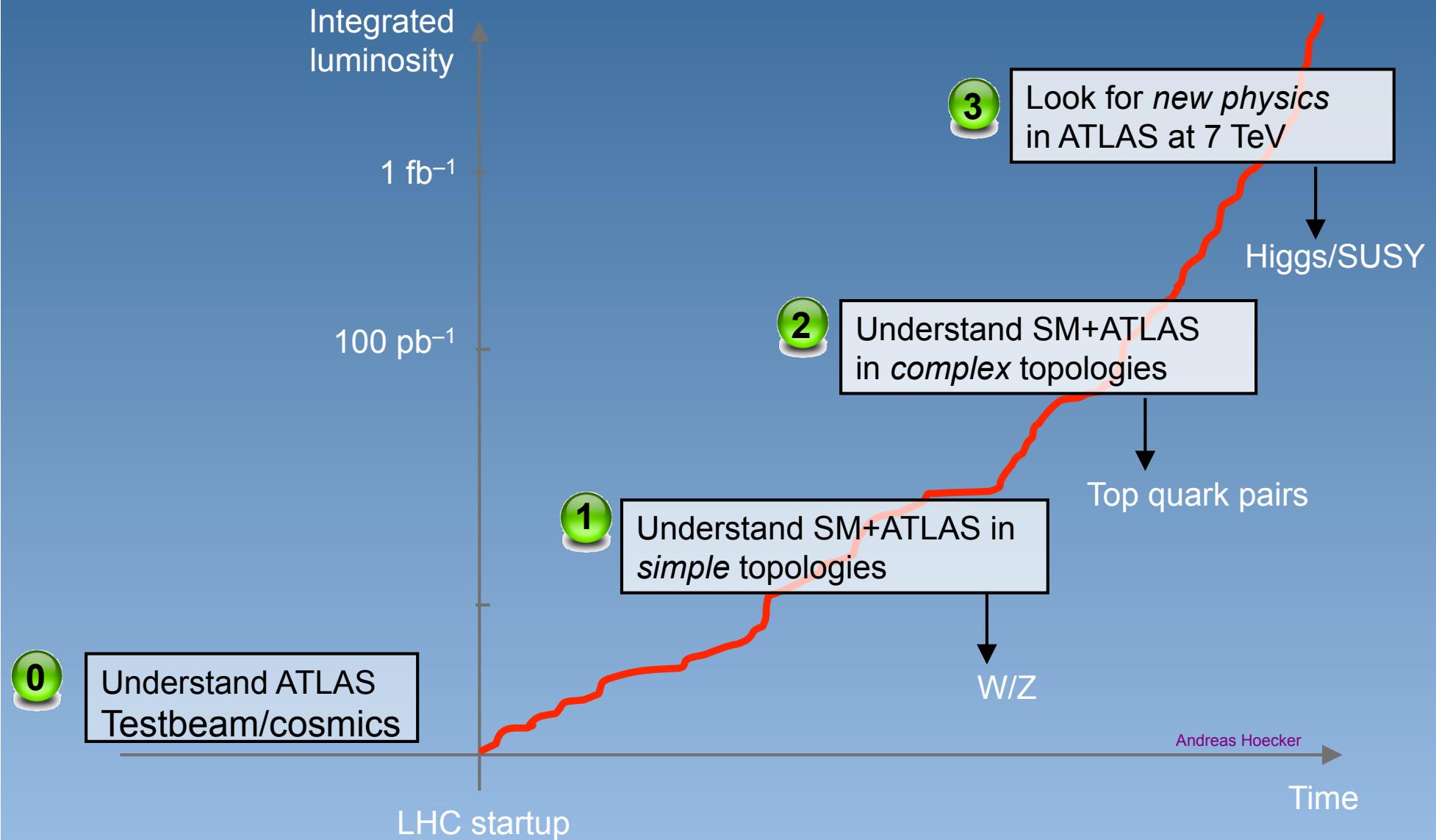
"Proton is a bag full of (anti-)quarks and gluons"



Only a fraction of the collision energy available for ‘hard’ interaction

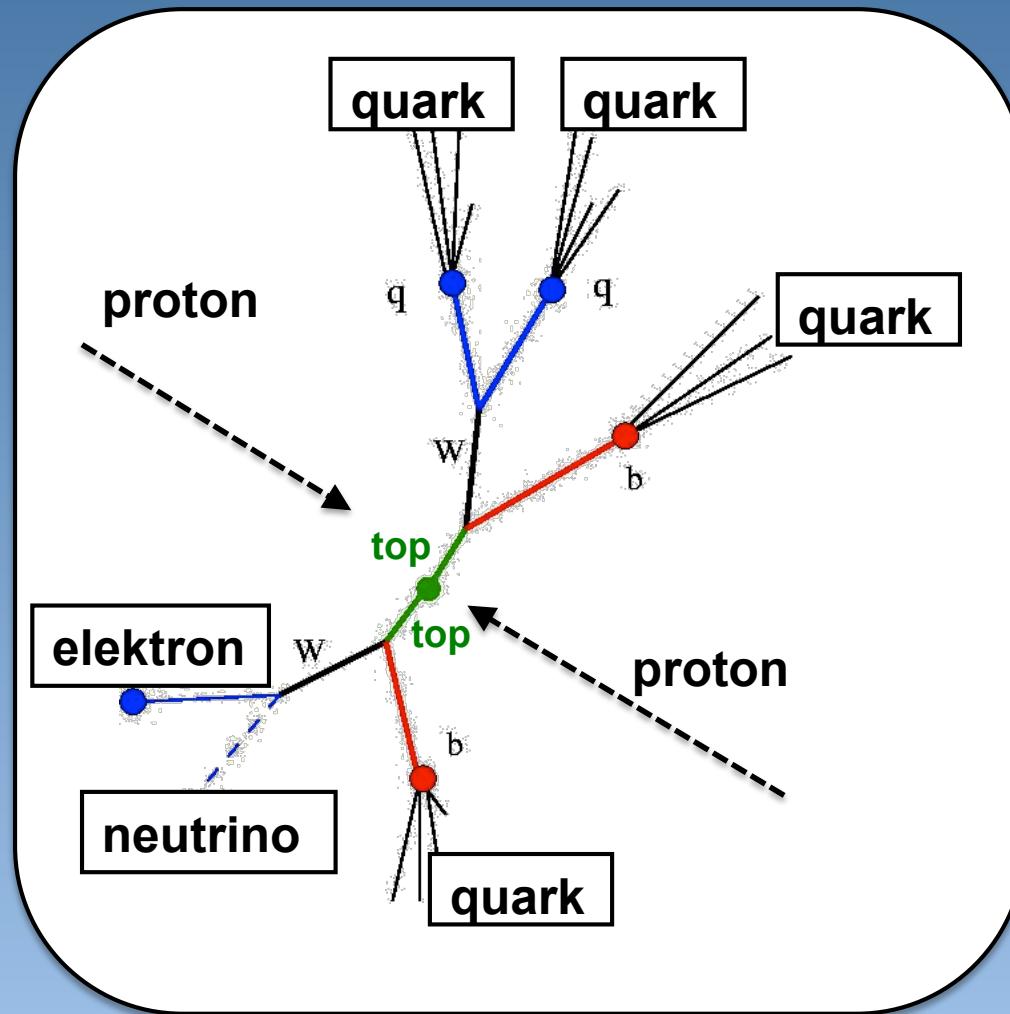
LHC start-up programme

New detector and new energy scale

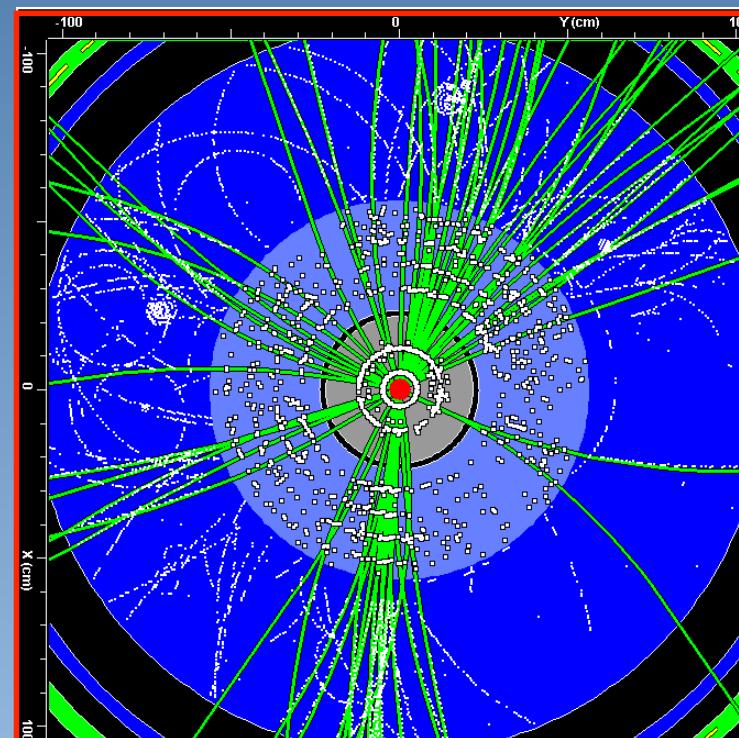


Interpreting LHC events

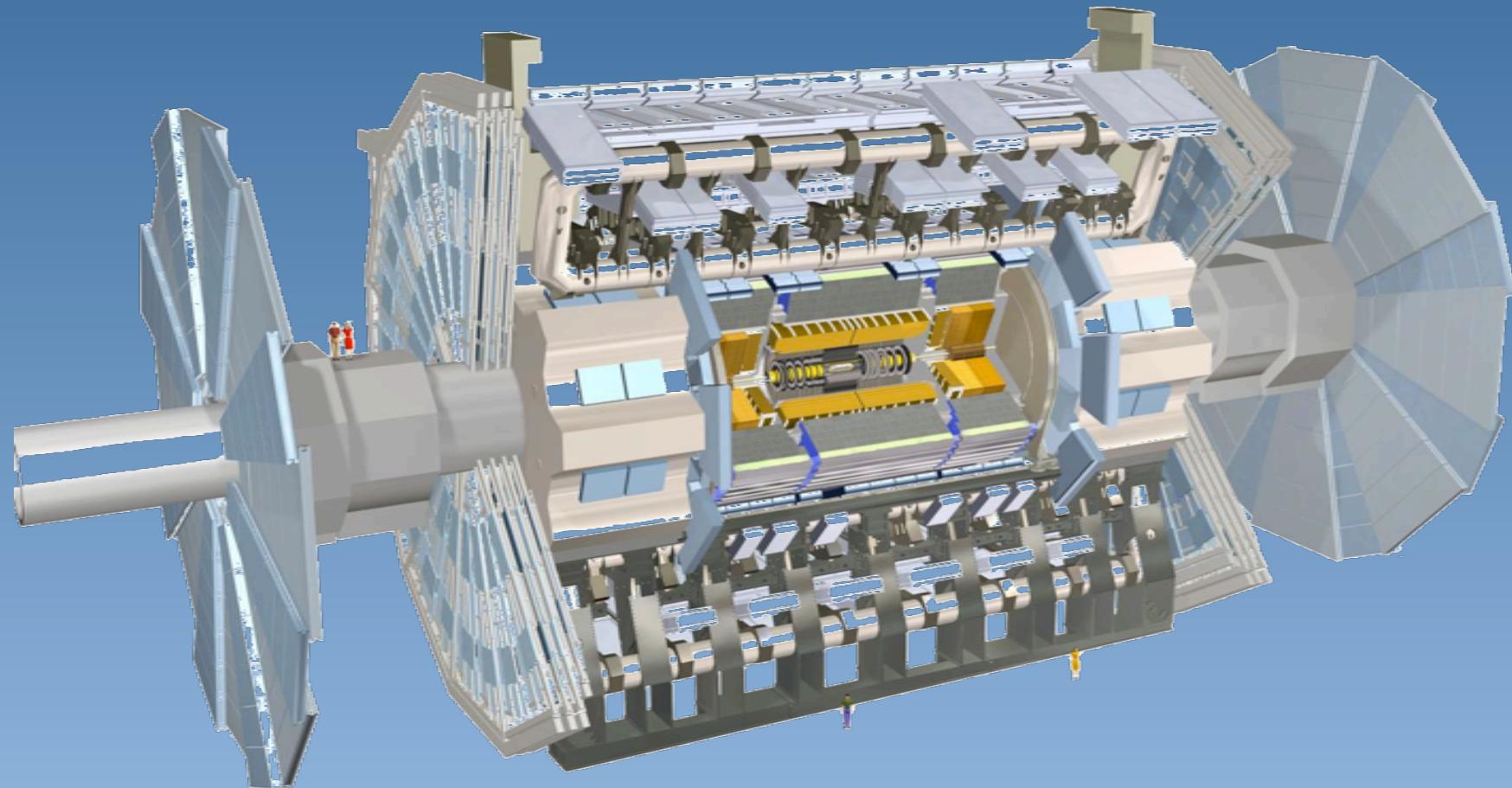
top quark pair-production



Simulation top quark production



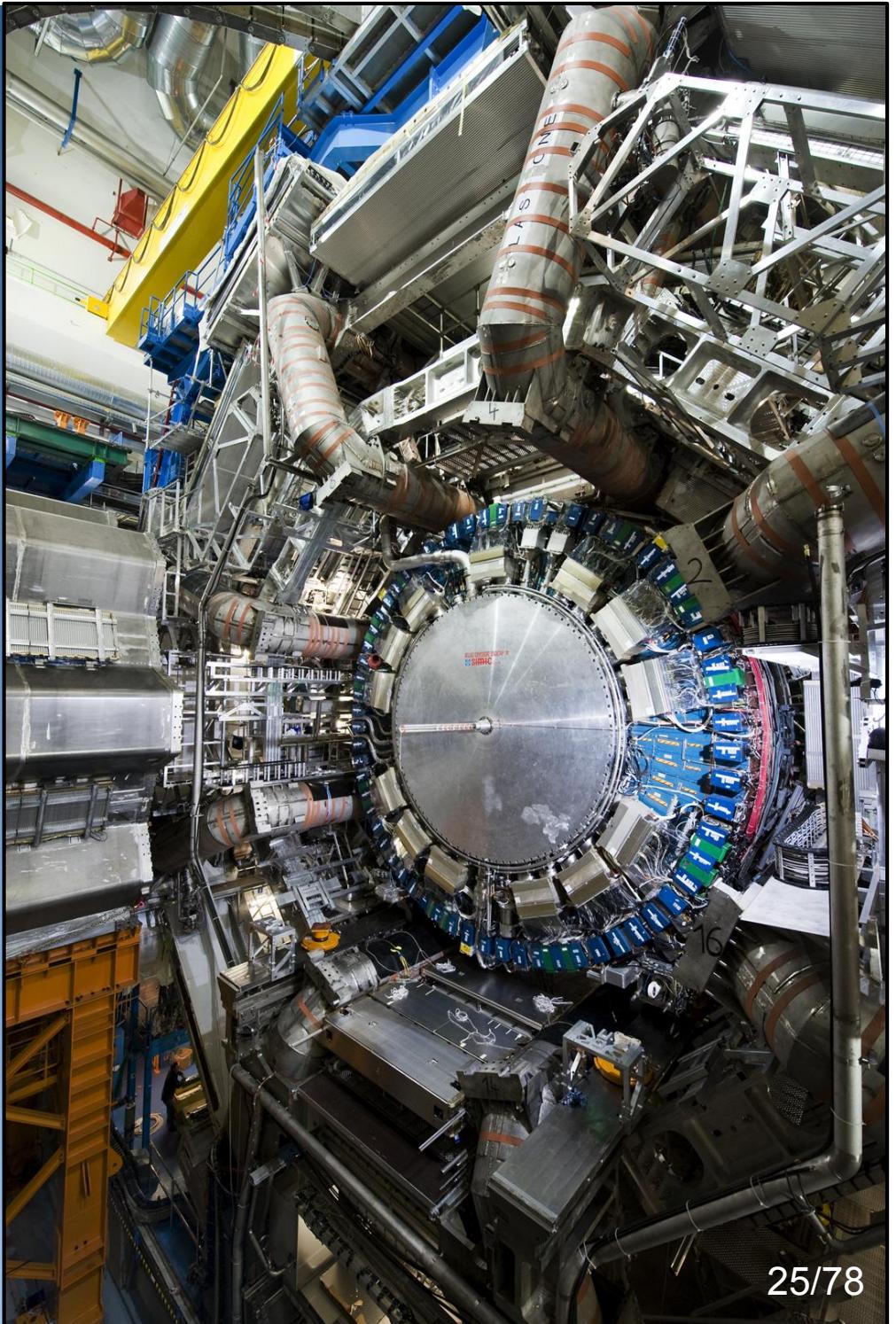
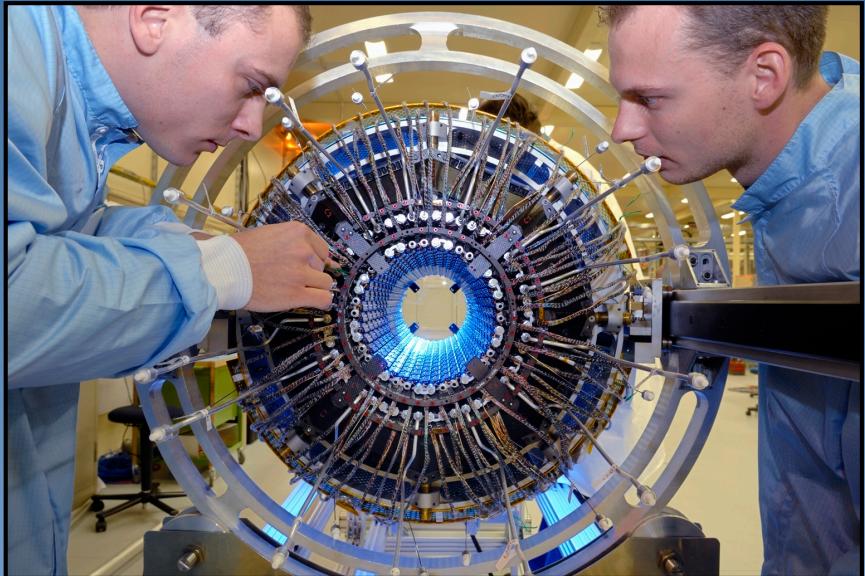
The ATLAS detector (as a computer model)



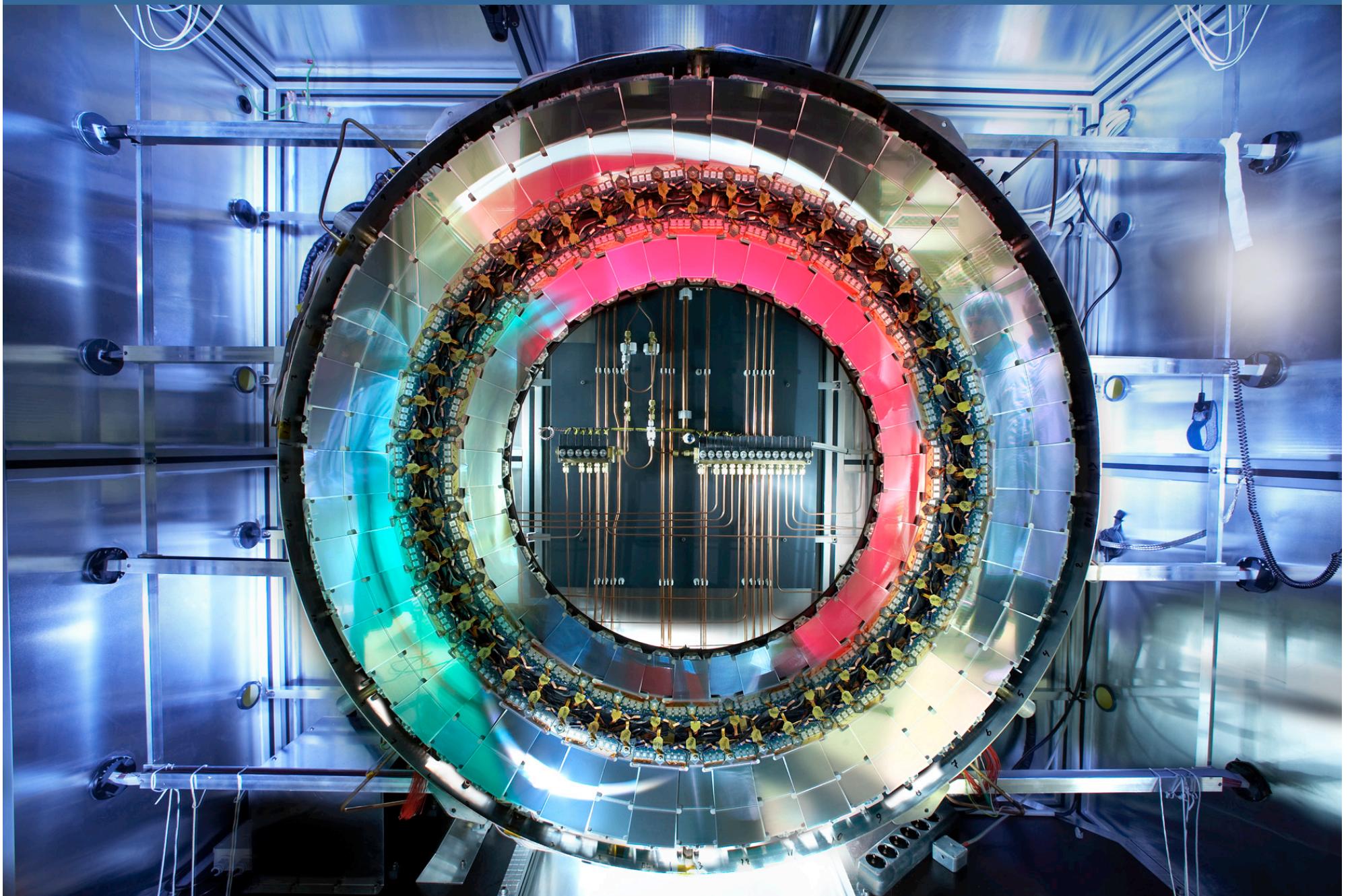
The Atlas detector really exists



The ATLAS detector (in real life)



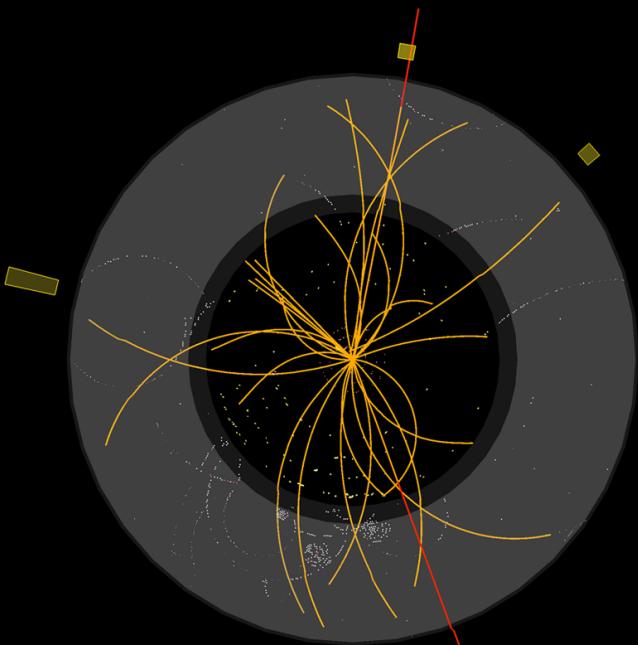
Part of the ATLAS detector has been designed and built at Nikhef





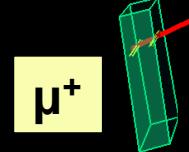
ATLAS EXPERIMENT

Run: 154822, Event: 14321500
Date: 2010-05-10 02:07:22 CEST

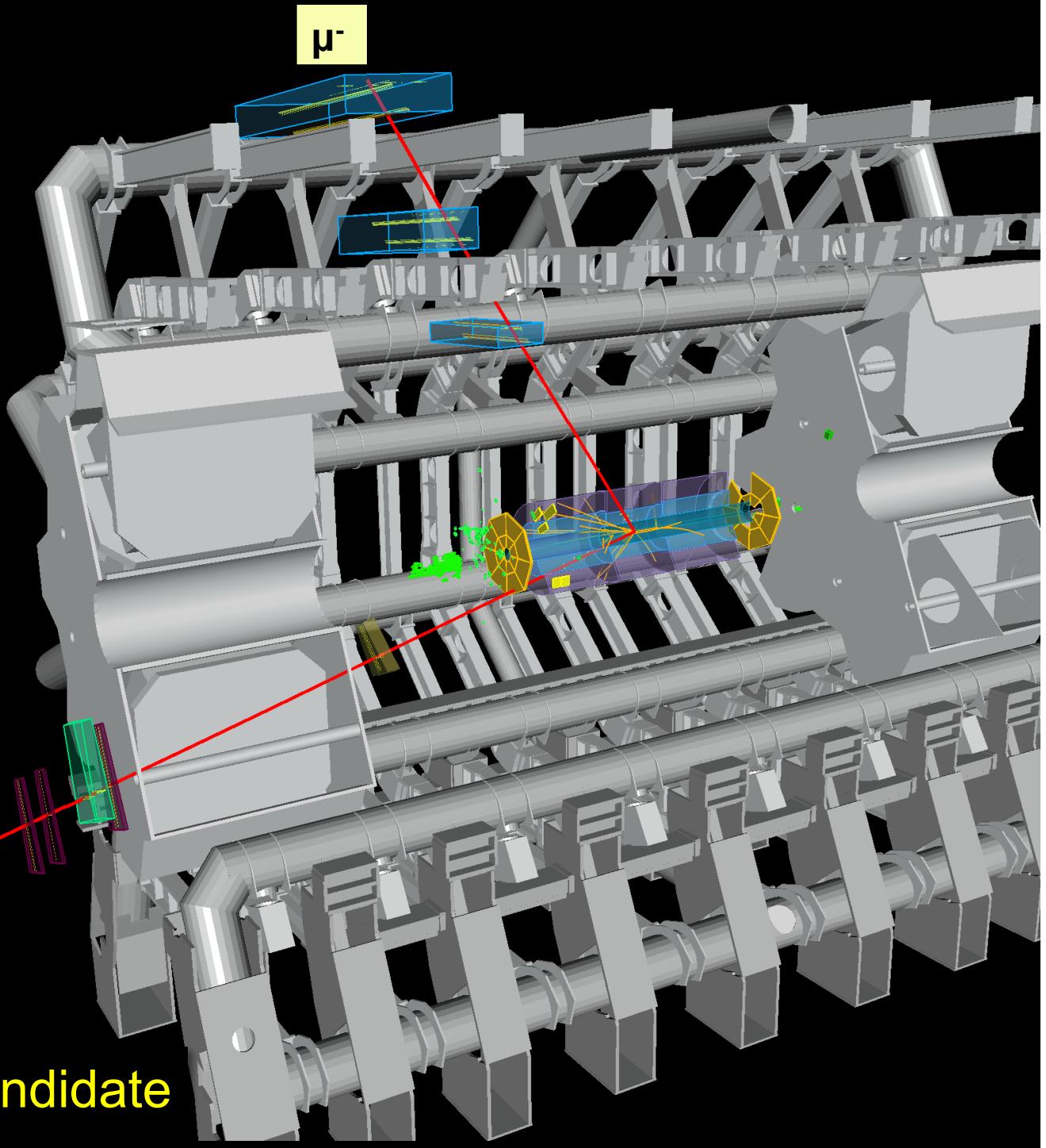


$$p_T(\mu^-) = 27 \text{ GeV} \quad \eta(\mu^-) = 0.7 \\ p_T(\mu^+) = 45 \text{ GeV} \quad \eta(\mu^+) = 2.2$$

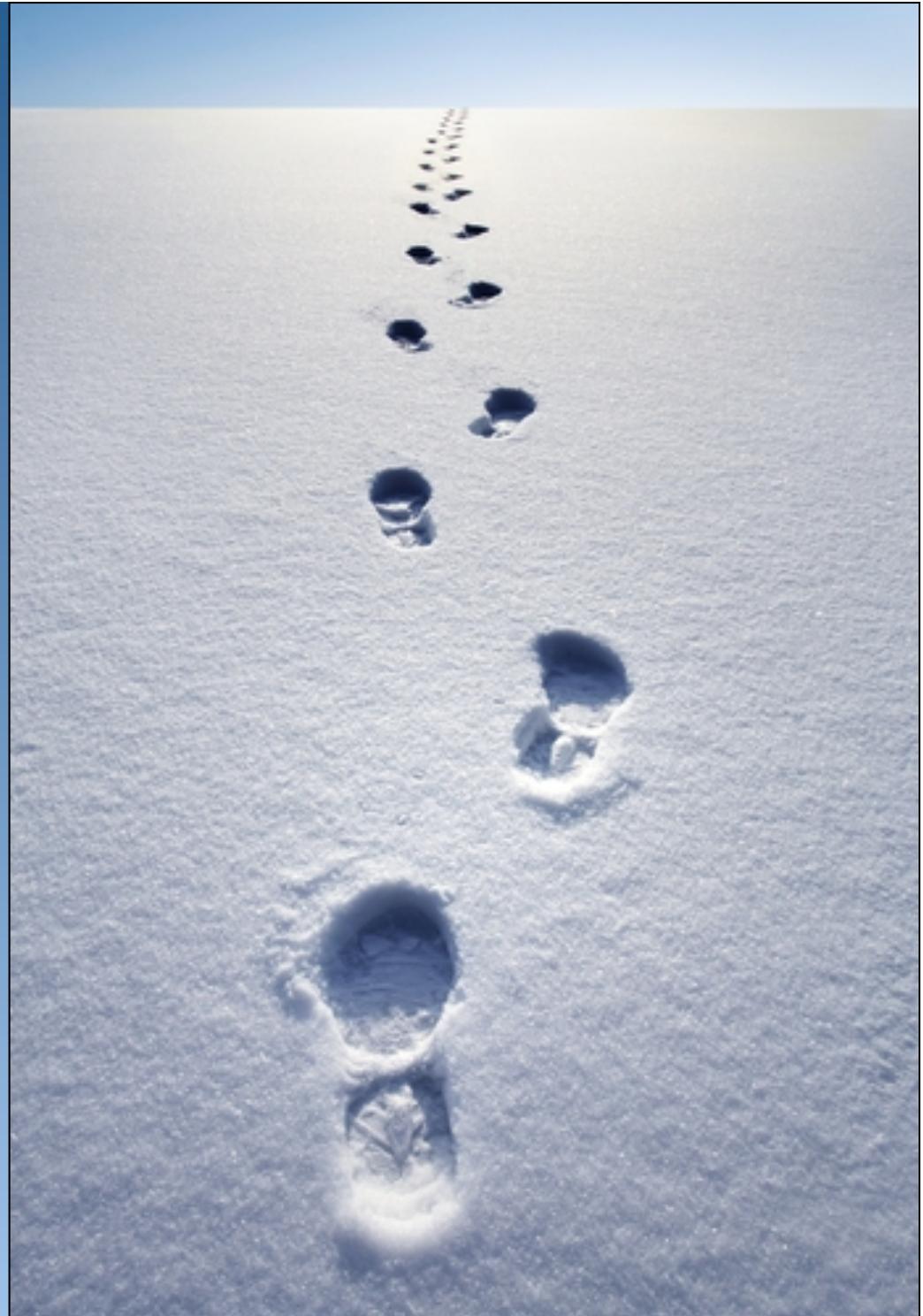
$$M_{\mu\mu} = 87 \text{ GeV}$$

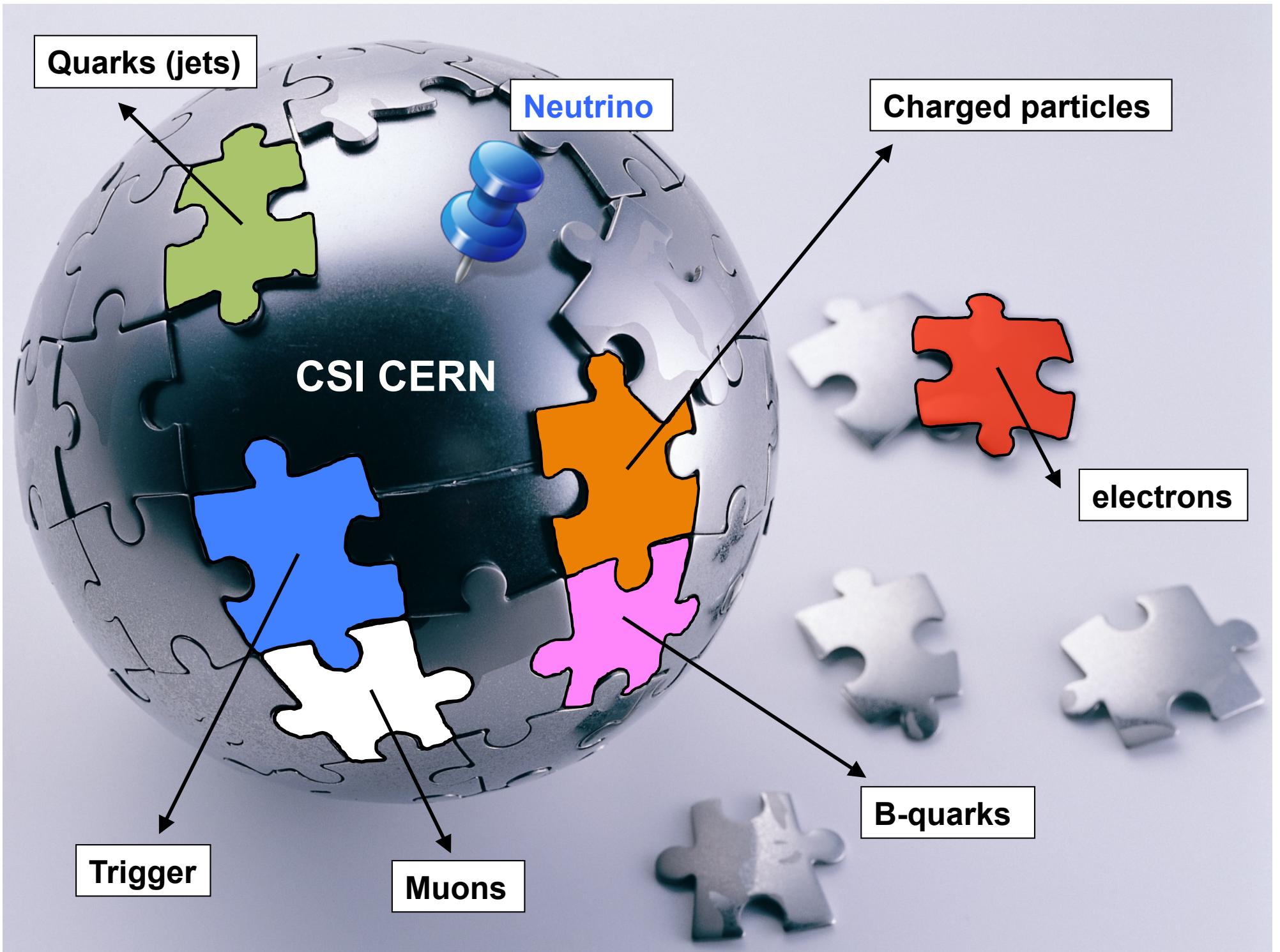


$Z \rightarrow \mu^+\mu^-$ candidate



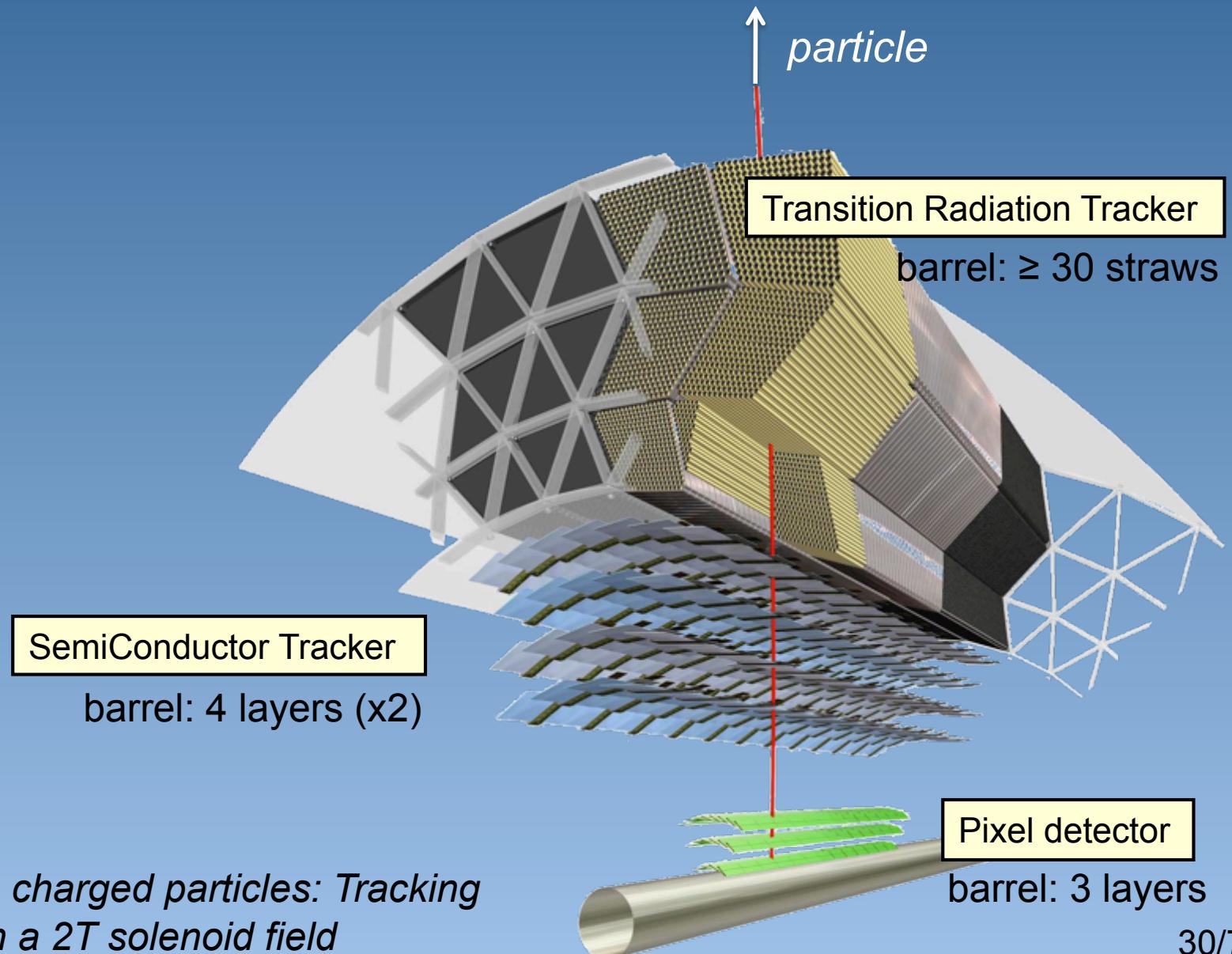
Rabbit, car or human ?



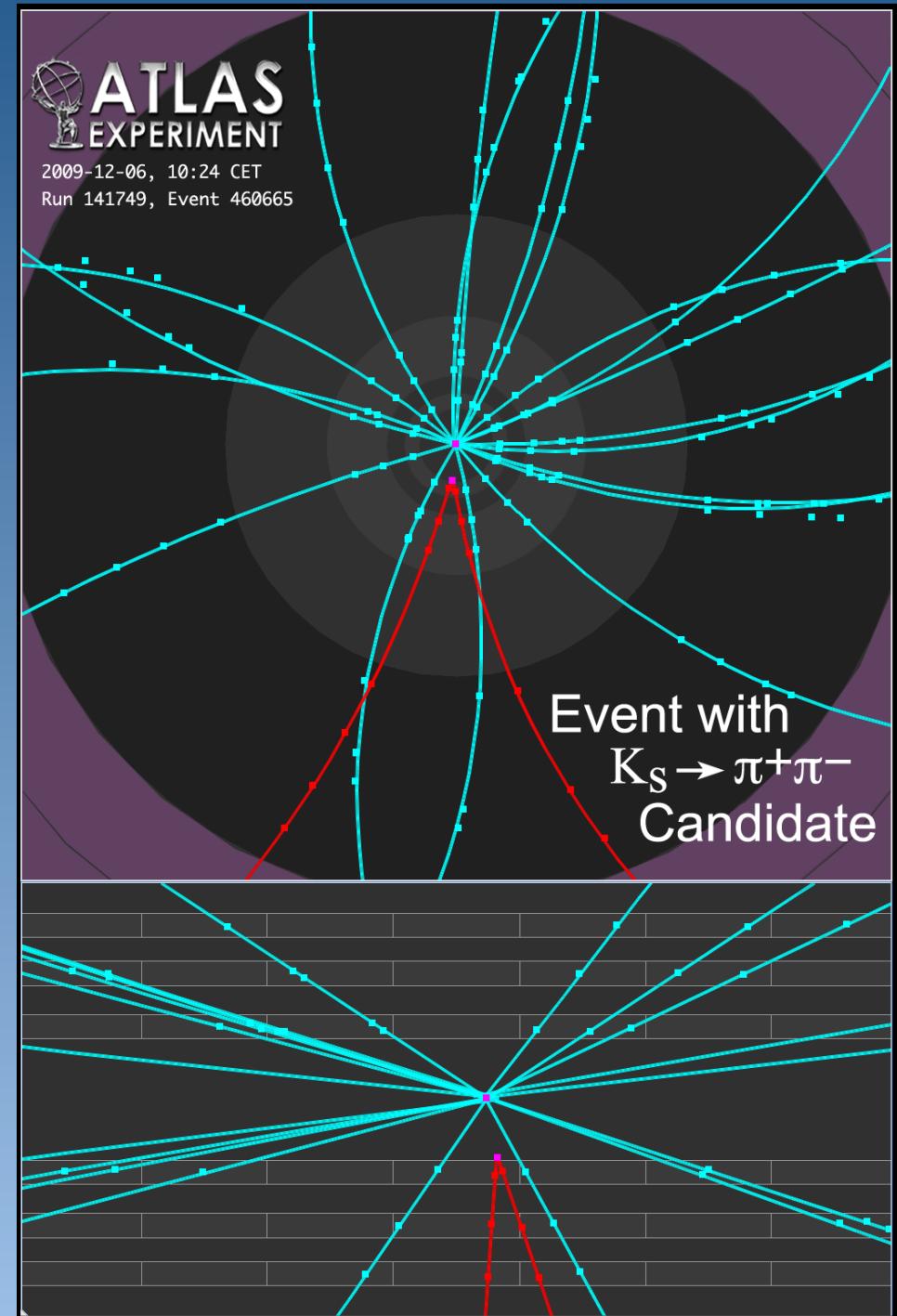
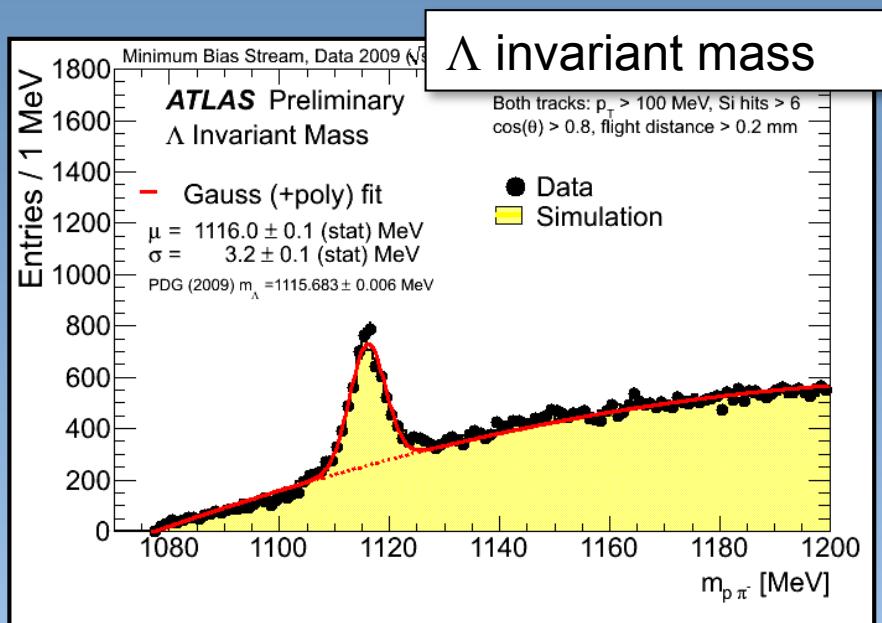
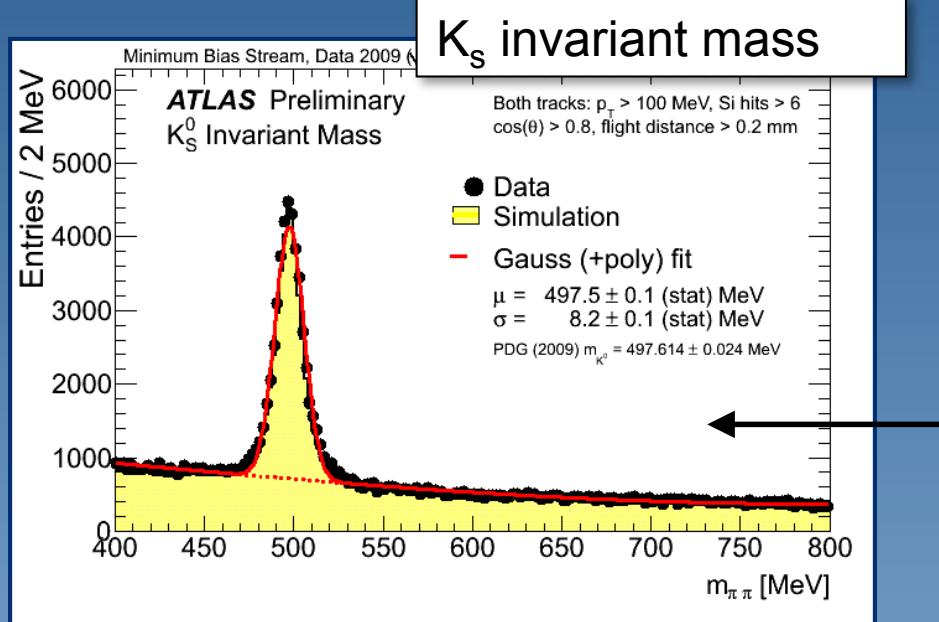


Complex structure ATLAS detector ?

Wide range of particle types / characteristics

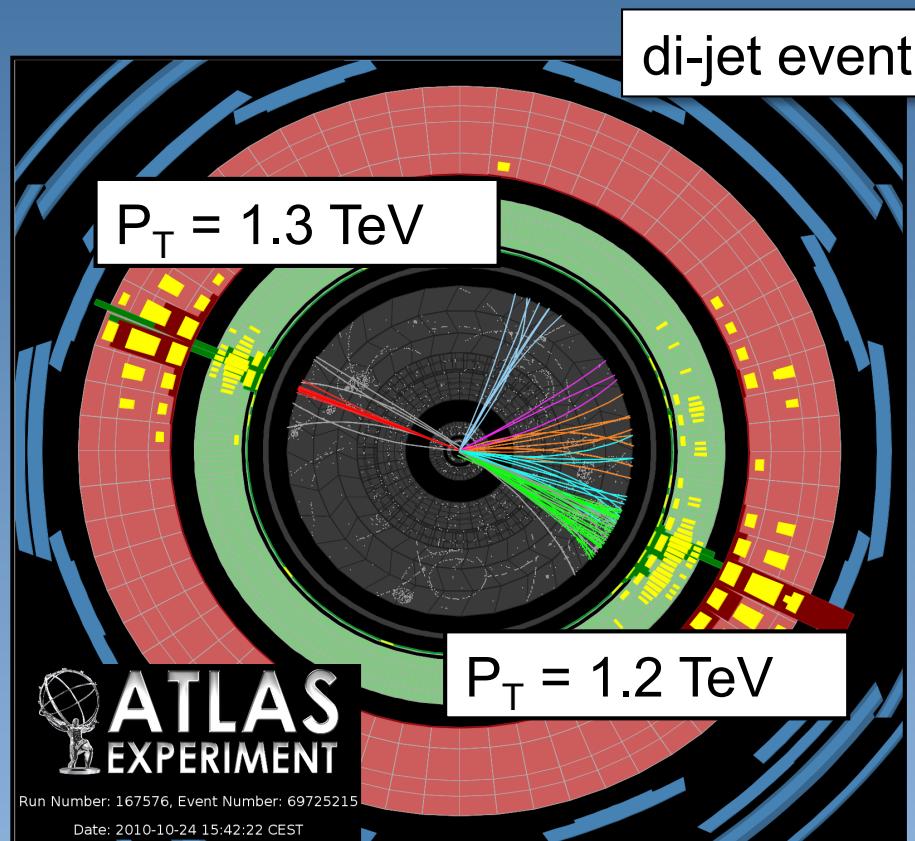


Secondary vertices

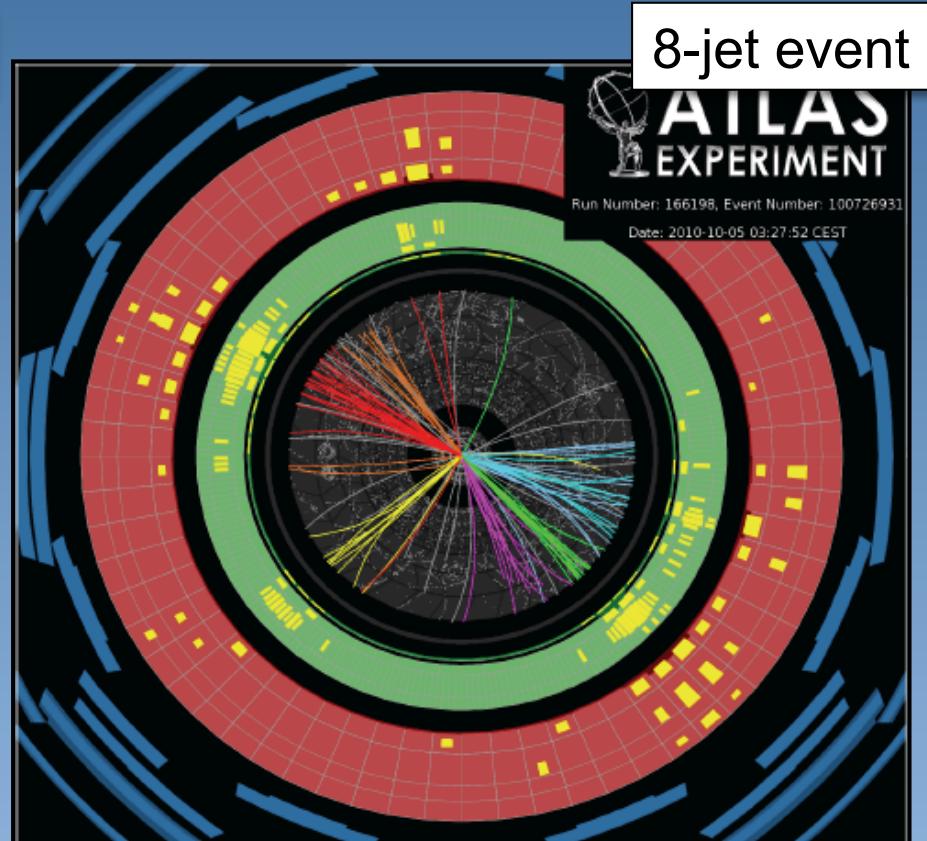


Quark (jet) reconstruction

Tracker and calorimeters



High transverse momentum

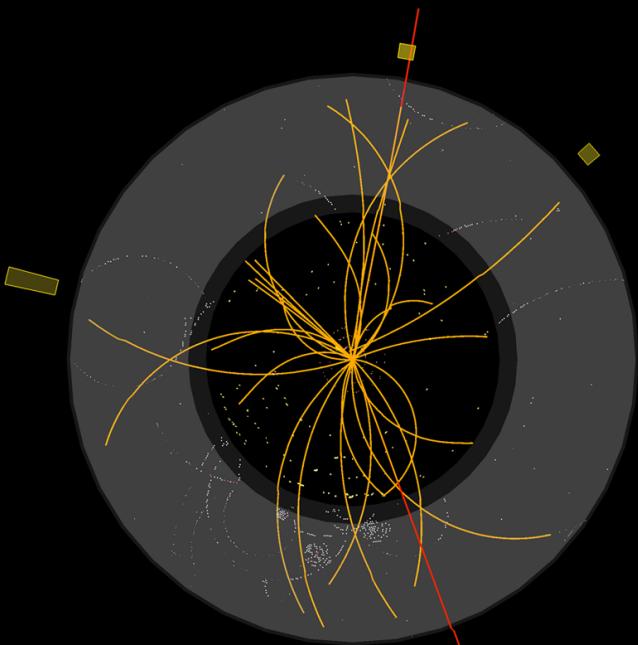


Multi-jet: 8 jets with $P_T > 60 \text{ GeV}$



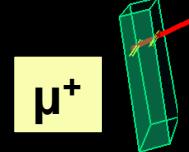
ATLAS EXPERIMENT

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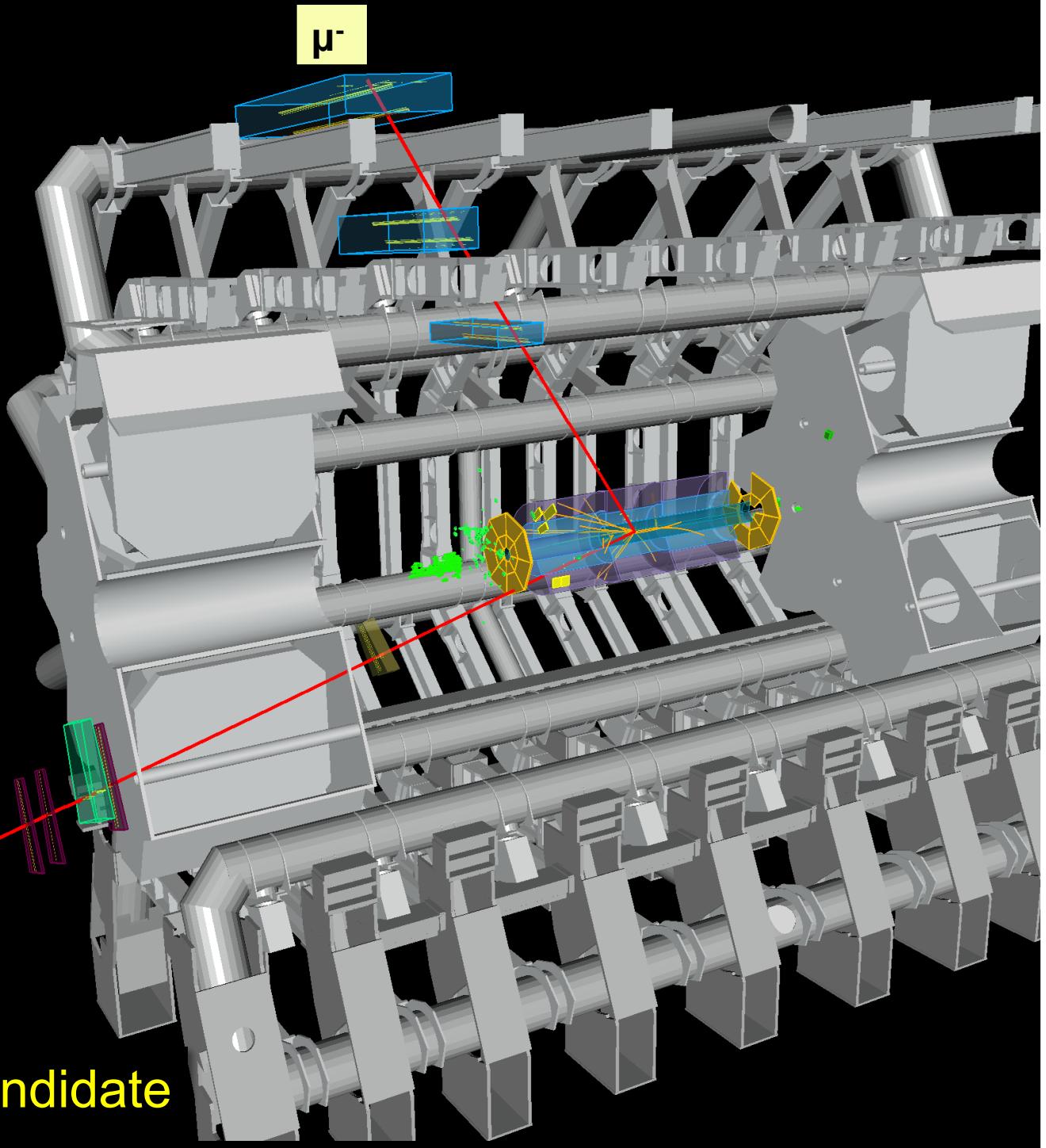


$$p_T(\mu^-) = 27 \text{ GeV} \quad \eta(\mu^-) = 0.7 \\ p_T(\mu^+) = 45 \text{ GeV} \quad \eta(\mu^+) = 2.2$$

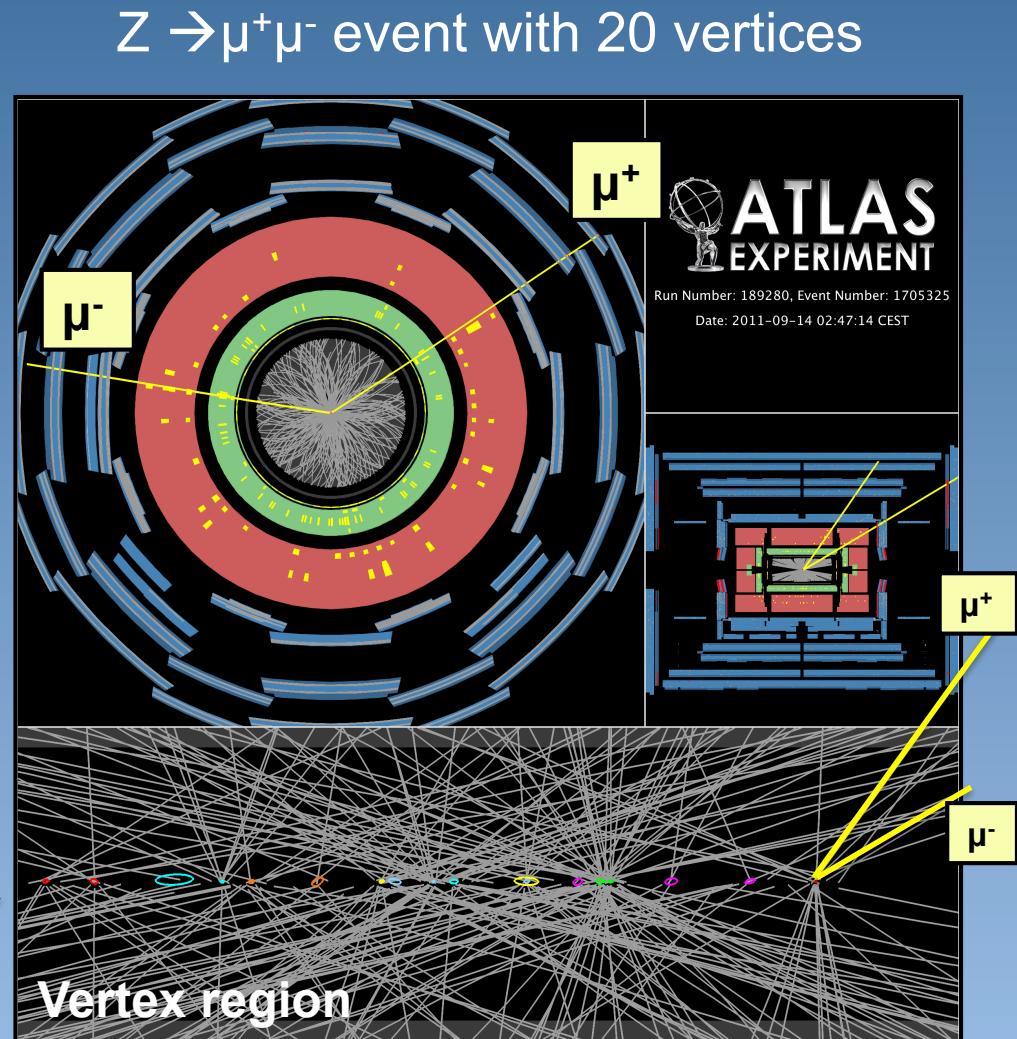
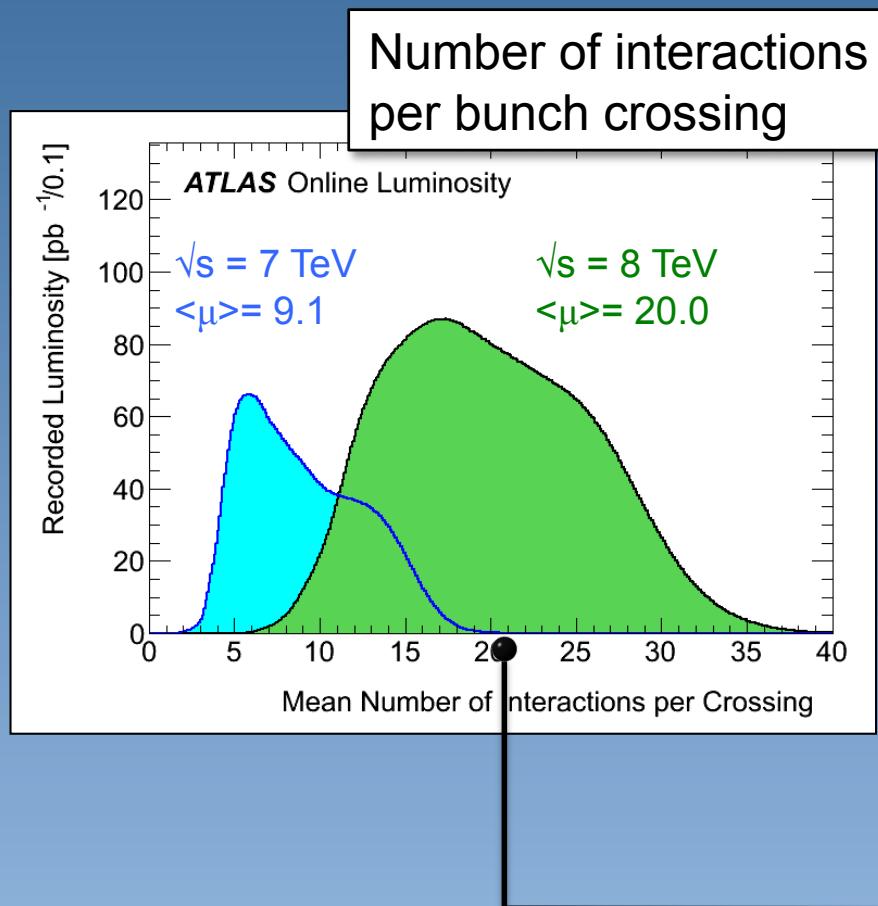
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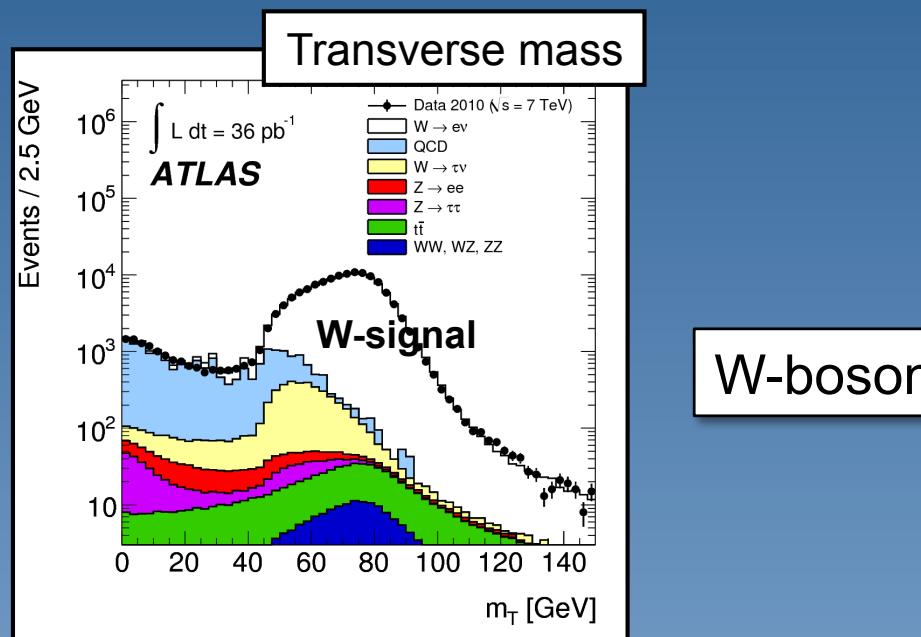
$Z \rightarrow \mu^+\mu^-$ candidate



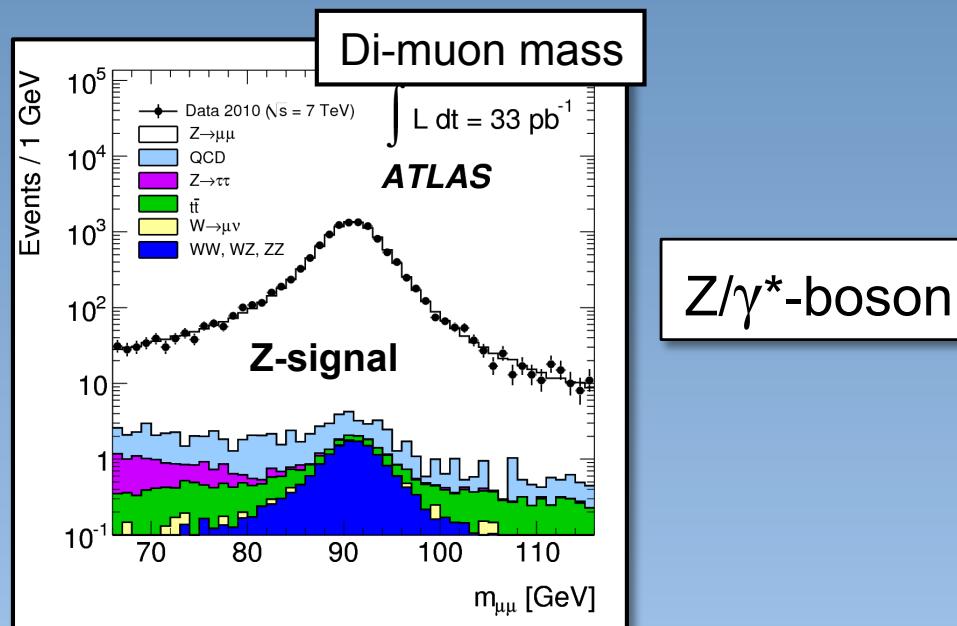
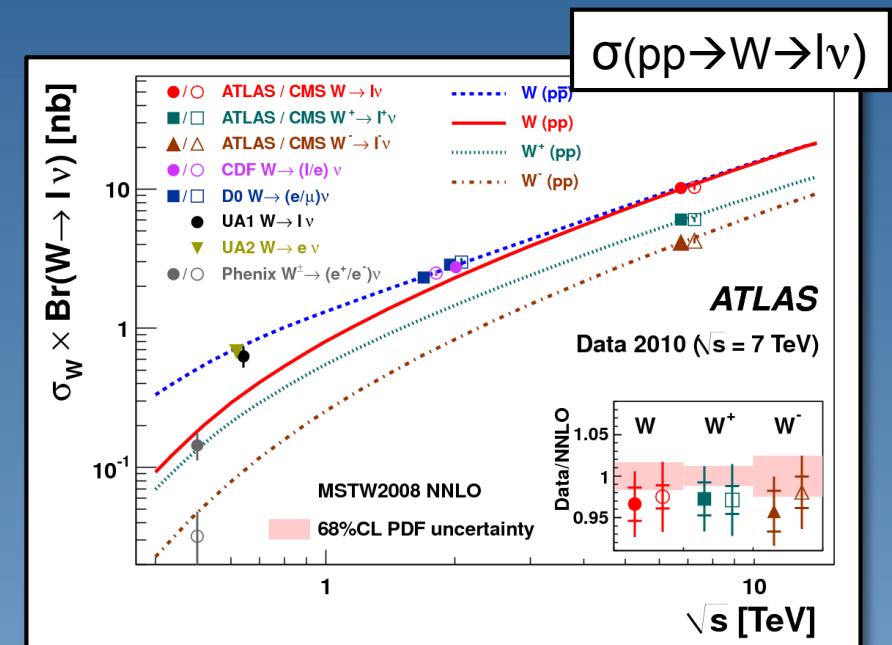
LHC is not a quiet environment



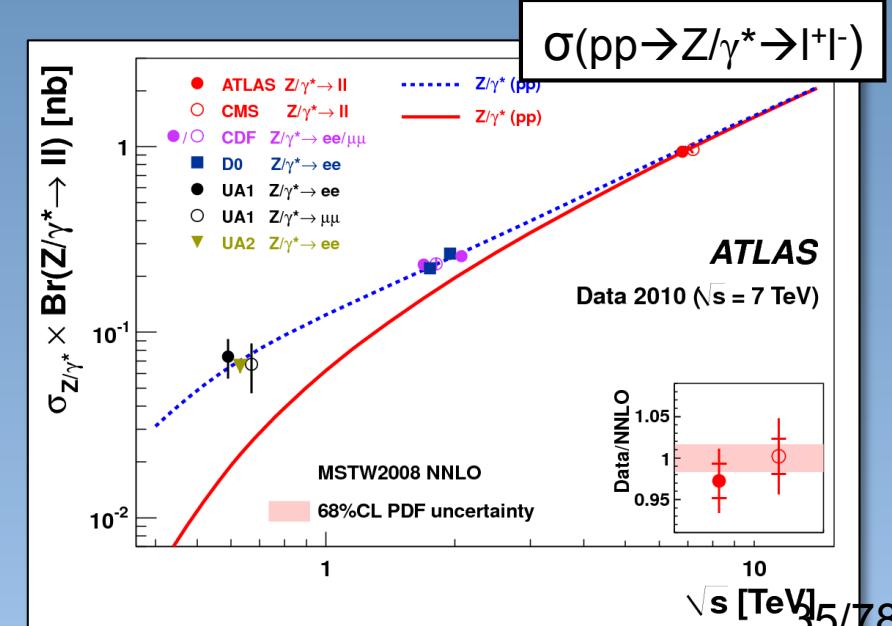
W and Z boson production at the LHC



W-boson

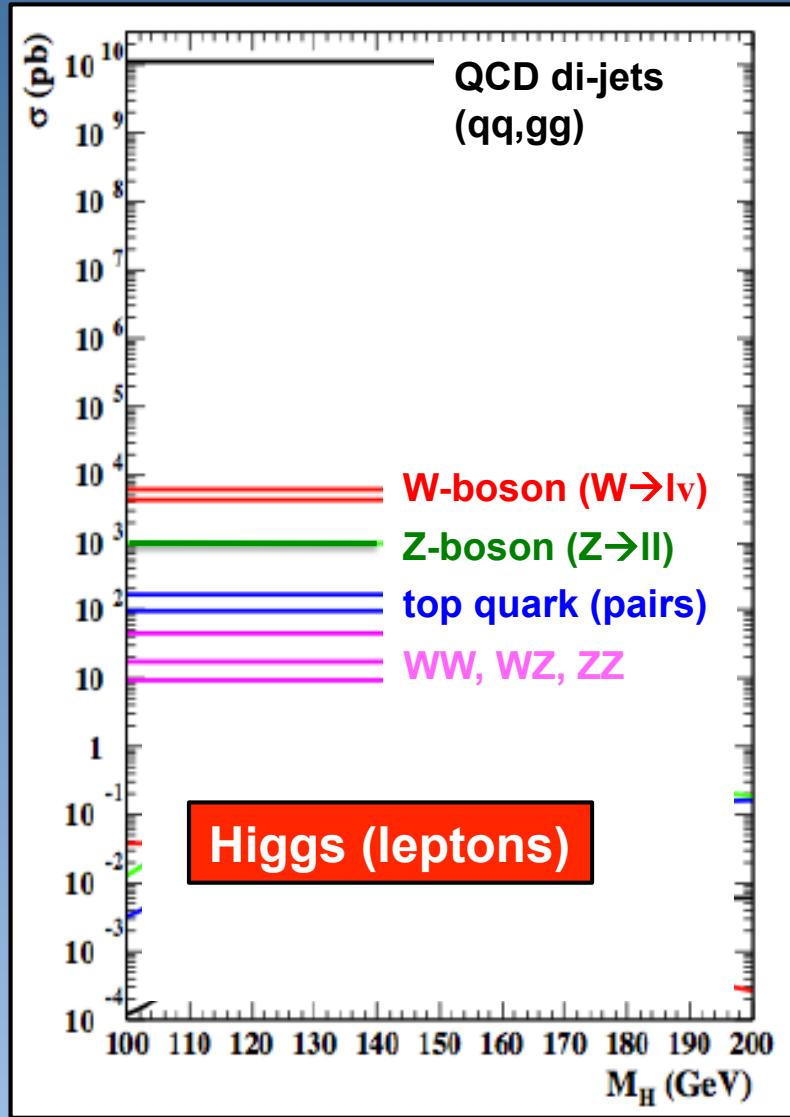


Z/ γ^* -boson



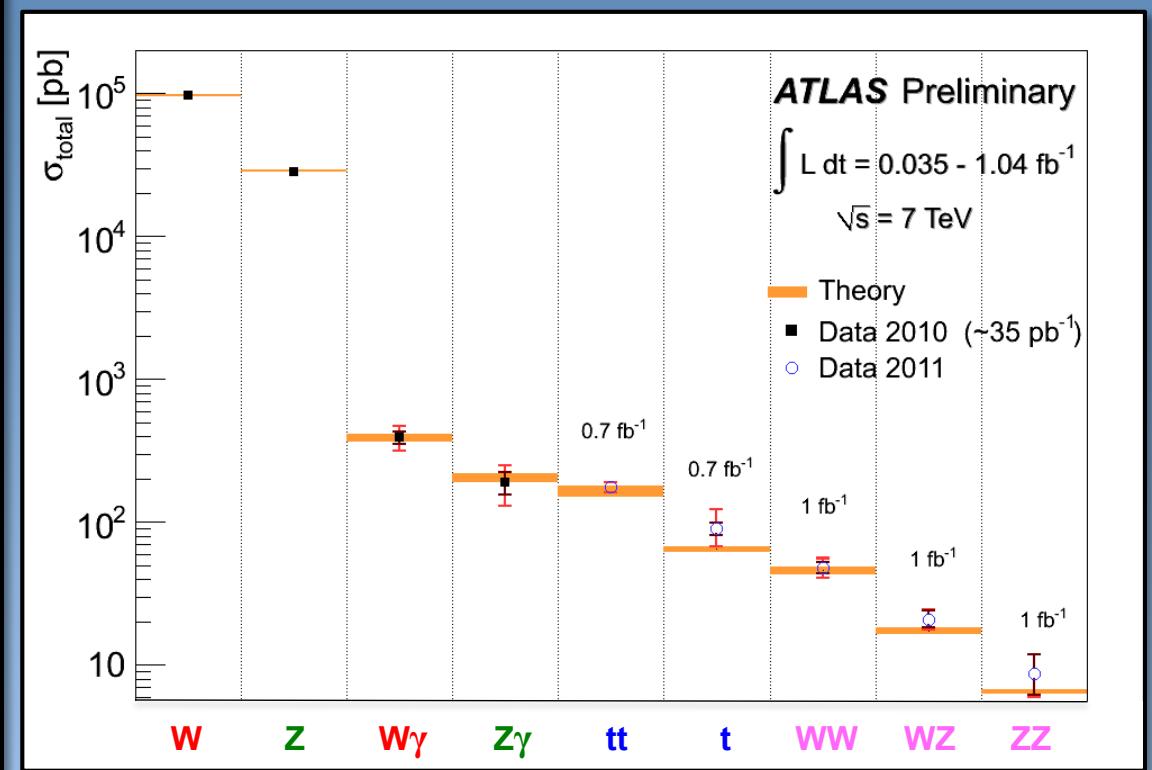
Standard Model processes at the LHC

predicted cross-sections



*Note: many many many jets (quarks/gluons):
→ look for leptons (electrons/ muons)*

ATLAS measurements



The ATLAS collaboration

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⁶, S.M. Farrington^{11B}, P. Farthouet²⁰, D. Fasching^{17Z}, P. Fassina^{11B}, D. Fassoulis¹⁸,
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- J.C. Hill²², N. Hill²², K.H. Hiller⁴¹, S. Hillert²⁰, S.J. Hillier¹⁷, I. Hincliffe¹⁴, E. Hines¹²⁰, M. Hirschbuech¹⁷, J. Hobbs¹⁴⁸, N. Hod¹³³, M.C. Holdenkinson¹³⁹, P. Hodgson¹³⁹, A. Hoekstra¹⁴⁶, J. Hoffmann⁴³, D. Hoffmann⁴³, M. Holheid¹⁴¹, M. Holder¹⁴¹, A. Holmes¹³³, S.O. Holmgren¹⁴⁶, J.L. Holubka⁸⁸, T.R. Homer¹⁷, Y. Homma⁶⁷, D. Horazdovsky¹²⁷, C. Horn¹⁴³, S. Horner⁸⁸, J.Y. Hostachy⁸⁸, T. Hotz²⁰, S. Hou¹, M.A. Houtman¹²⁴, A. Hounamada¹²², J. Howarth⁸², I. Hristova⁴¹, J. Hrvatsic¹¹⁸, J. Huskova¹²⁶, T. Hyra^{ova}⁴, P.J. Hu¹⁷⁸, S.C. Hua¹⁴, G.S. Huan¹, F. Husbant⁸³, F. Huuegging¹²⁶, J. Hwang¹¹⁸, E.M. Hu¹, M.-C. Hu¹, J.-D. Hu¹

- P. Hurst⁵⁷, M. Hurwitz¹⁴, M. Ibbotson⁸², I. Ibragimova¹⁰⁸, O. Isenkina¹⁰⁸, Y. Ikesan¹⁰⁸, C. Leggett¹⁴, M. Lehmann²⁰, G. Lehmann Miotto²⁰, M. Lehmann²⁰, D. Lelouch¹⁷¹, J. Lelouch⁷⁸, M. Leitchouk³⁴, V. Lendermann¹⁰⁸, P. Lenz¹³⁶, K. Lenz¹⁴³, C. Lenz⁹³, J. L. Lenz¹⁶⁹, J. L.

- xx. *Swanson*¹⁰⁷, *R. L. Shultz*¹⁰⁸, *A. Limosum*¹⁰⁹, *M. Liniper*¹⁰⁸, *S.C. Jakobsen*¹¹⁰, *S. Jakobsen*¹¹¹, *H. L. Lipsey*¹¹², *E. Lipese*¹²⁰, *L. Lipinsky*¹²⁰, *A. Lipinsky*¹²⁰, *T.M. Lipinsky*¹²⁰, *A. Jarlakog*⁷⁰, *L. Jeanty*⁸⁷, *H. Liu*⁸⁷, *J.B. Liu*⁸⁷, *M. P. Nemethy*¹⁰⁸, *A.A. Nemethy*¹⁰⁸

- W. Ji^{1*}, Y. Jiang^{1**}, M. E. Lobodzinska^{4†}, P. I. L.G. Johansen¹³, M. Joha A. Loginov^{17b}, C.W.I. R.E. Long⁷¹, L. Lopes¹⁸, R.M. Neves^{10b}, P. Nevski²⁴, P.R. Niedercorn^{11b}, J. Nielsen^{13†}, T. N.

- M.S. Kayl¹⁰, V.A. Kazan¹¹, G.D. Kekelidze⁶⁸, M. Kell¹², K. Kondo¹³, S. Kuroda¹⁴, J. MacLennan McGuire¹⁵, T. Maeno²⁴, P. Mättig¹⁶, C.A. McGrath¹⁰⁴, Y. Matsuura¹⁷, H. Ohshita¹⁴⁰, T.K. Oh¹⁸, A.G. Olchевский⁶⁹, M. Ochiai¹⁹, J. Olazowski³³, C. Orr²⁰, H. Sandak²¹, S. Sandvoss²², H. Sato²³,

- B.P. Kervenov¹, S. Kerst¹, A. Khanov^{1,2}, D. Kharche¹, N. Khovanskiy³, V. Kho¹, A. Maiorova⁴, S. Majer⁵, F. Malek^{6,7}, U. Minkov⁸, A. Manabe^{8,9}, L. Mandel¹⁰, F. Orellana⁴⁹, Y. Oren⁵⁰, Osculati^{50a,50b}, R. Ould-Saada¹¹, A. Ould-Saada¹¹, H. Santos⁴, O. Saasik⁶, L. Sawyer¹²

- P.C. Kim¹⁴², S.H. Kim¹⁶⁰, A. Mann⁵⁴, P.M. Mani¹, N. Osturk⁷, A. Pacheco¹, D. Pallin³³, A. Palma¹, S. Pandit²⁴, D. Panteleev¹, J. Pascual¹, J. Schausberger¹, A.G. Schaumann¹, M. Schioppa¹

- A. Klimentov²⁴, R. Klinge
 T. Kluge⁷³, P. Kluit¹⁰⁵, S.
 M. Kobel⁴³, B. Koblitz²⁰,
 S. Marti-Garcia⁶⁷, A.
 T.A. Martin¹⁷, B. Mar-
 F. Marzano^{132a}, A. M.
 S. Pashinian¹, D. Pava-
 A. Paramonov⁸, S.J. P.
 E. Pasqualucci^{132a}, A.
 J.W. Schulte^{132a}, M.
 M. Schmitz^{132a}, C.
 Schroeck^{132a}, J.W. Schu-

- J. Kohn,¹ L. Kofman,¹ S.J. Maxneide,¹ E.N. A. L. Kononov⁴⁸, R. Konop¹, S.P. Mc Kee⁸⁷, A. Mc G. Kordas¹⁸⁴, V. Koreshev¹, O. Kostylev¹, N. Kuznetsov¹, S.J. Maxneide,¹ E.N. A. L. Kononov⁴⁸, R. Konop¹, S.P. Mc Kee⁸⁷, A. Mc G. Kordas¹⁸⁴, V. Koreshev¹, O. Kostylev¹, N. Kuznetsov¹, J.A. McFayden¹³⁹, H. M. Perrino^{72a}, P. Perrot¹, J. Petersen²⁹, T.C. Peters¹, D. Petschelt⁴¹, M. Petruzzelli¹, H. Severini¹, M. Shapiro¹, A. Shlomo¹

- O. Kortner^a, S. Kortner^a, A. Koutsman^{10b}, R. Kowa^a, V.A. Kramarenko⁹⁷, G. K. T.R. McMahon¹⁰, T.J. M. Medinnis⁴¹, R. Med. J. Meinhardt⁴⁸, B. Me E. Piccaro⁷⁸, M. Piccin M. Pinamonti^{164a, 164c}, M. Plumondon¹⁶⁹, W. O. A. Smit^{12a}, P. Sicho^{12a}, D. Silverst D. G.

- F. Krejci¹²⁷, J. Kretzschmar
 J. Kroeseberg²⁰, J. Krstic¹¹
 S. Kuehn⁴⁵, A. Kugel^{58c}
 A. Mengarelli^{15a,15b}, S.
 C. Meroni^{50a}, F.S. Mer-
 J. Meyer¹⁷³, J. Meyer²
 L. Poggiali^{11a}, T. Pogh
 V. Polychronakos²⁴, D.
 G. A. Pohle¹², D. P.
 P. Sinervo¹⁶, T.B. Sjurs
 K. Sliwa¹⁶

- M. Kuna⁸³, N. Kundu¹¹⁸, W. Kuykendall¹³⁸, M. Ku-
li, Lashko⁸⁰, J. Labby⁴, P. Miele²⁰, S. Migas⁷³, R.J. Miller⁸⁸, W.J. Mi-
ke¹⁰⁰, M. Minami¹⁰⁷, J.A. Mi-
ller¹⁰⁰, S. Mirelman¹⁰⁰, S.P. Popelis¹²⁷, I.N. Pota-
pov¹²⁷, P. Pralavorio⁸³, S. Pra-
mela¹⁰⁰, D. Ralston¹⁰⁰, D.
Rothberg¹⁰⁰, C. Ross¹⁰⁰, B.C. Smith¹¹¹,
J. Snow¹¹¹, M. Solar¹²

- E. Ladygin⁶⁸, R. Lafaye¹,
W. Lampf¹, E. Lancon¹³⁶
M. Minami^{1, 2, 3}, A. M. Mirabet¹,
G. Mirabelli^{132a}, L. Miura¹³³,
V.A. Mitsou¹⁶⁷, S. Mizutani¹³³,
M. Mizutani¹³³, T. Mizutani¹³³,
P.M. Richardson¹, D. Przybilla¹³,
X. Prudent⁴³, H. Pryzgoda¹³,
Y. Pylypchenko¹¹⁷, J. Qian¹³⁰,
J. Sondergaard¹, S. Spagnoli¹,
M. Strocen¹³

- A.J. Lankford¹, P. Lam¹, S. Moed¹, V. Moeller¹, M. Raas^{10*}, V. Radescu¹¹, E. Staneck¹², A.V. Larionov^{12*}, A. Larionova¹², A.M. Moisseev^{12*}, R. S. Rajagopal²⁴, S. R. P.N. Redoff¹³, F. Rausch¹⁴, I. S. Sapegin¹⁵, P. Starov¹⁶, I. S. Staneck¹², W. Lavrijen¹⁴, P. Laycock¹⁷, S. Montesano^{8a, 8b}, F. M. Spadaccini¹⁸

- E. Le Menedec^{13b}, M. Lee^{13c}, J.S.H. Lee^{13d}, S.C. Lee^{13e}
 A. Moraes^{13f}, A. Morai^{13g}, P. Morettini^{13h}, M. Mo¹³ⁱ, J.D. Morris^{13j}, H.G. M^{13k}, B. Resende^{13l}, P. Rezn^{13m}

- E.J.W. Moyse⁸⁴, M. M. A. Muijs¹⁰⁵, A. Muir¹⁶, A G Myaskov¹²⁸ M. P. S. Rieke⁸¹, M. Rijpstra¹, G. Rivoltella^{80a,80b}, F. IEMF Robinson⁷⁷, M. R. P. Strizene¹, I. Stumer², C. Suhr¹⁰⁶

- Y. Nakahama^{11a}, K. N. S. Rodier⁸⁰, D. Rodriguez¹, K. Suruliz¹,
N.R. Nation²¹, T. Natu¹, I. Sykora¹,
A. Ni¹, -119a, 119b, G. N. A. Romanouk⁹⁸, V.M. A. Taga¹¹,

- A. Negri¹, G. Neri¹, M. Rose², G.A. Rosent³, L.P. Rossi^{2,4}, L. Rossi⁵, C.R. Royon^{1,6}, A. Rozza⁷, A. Talyshyan⁸, N. Tamouh⁹, P. Tusa^{1,2,6}

- G. Rudolph⁴², F. Rühr,
K. Runge⁴³, O. Runolf,
P. Ruzicka^{12a}, Y.F. Ry-

- A.F. Saavedra^{10a}, I. Sa
G. Salamanna^{10a}, A. S
R.S. Reuter
S. Thoma
E. Thomae
G. Wieden
C. Wieden

- C.J.W.P.
T. Todorov
K. Tolleson

- C. Topfel¹
D.R. Tove
S. Trincavelli

- # aboration

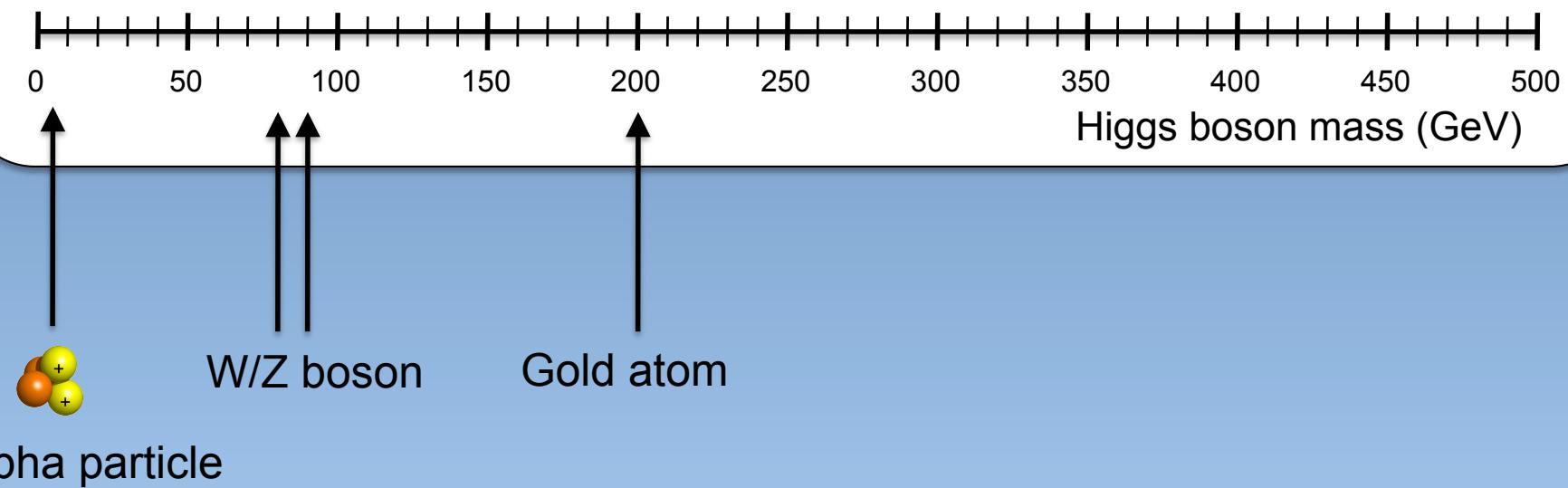
Trigger:
From 40 MHz (LHC) to a few hundred Hz



Bottleneck: disk space

Selecting interesting physics from an e-nor-mous background

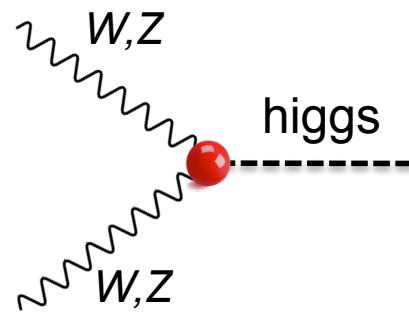
The mass of the Higgs boson



How can you produce a Higgs particle ?

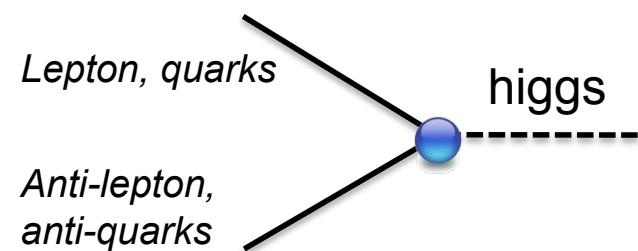
Gauge bosons

*Massive gauge boson ?
... then it coupling to the Higgs!*



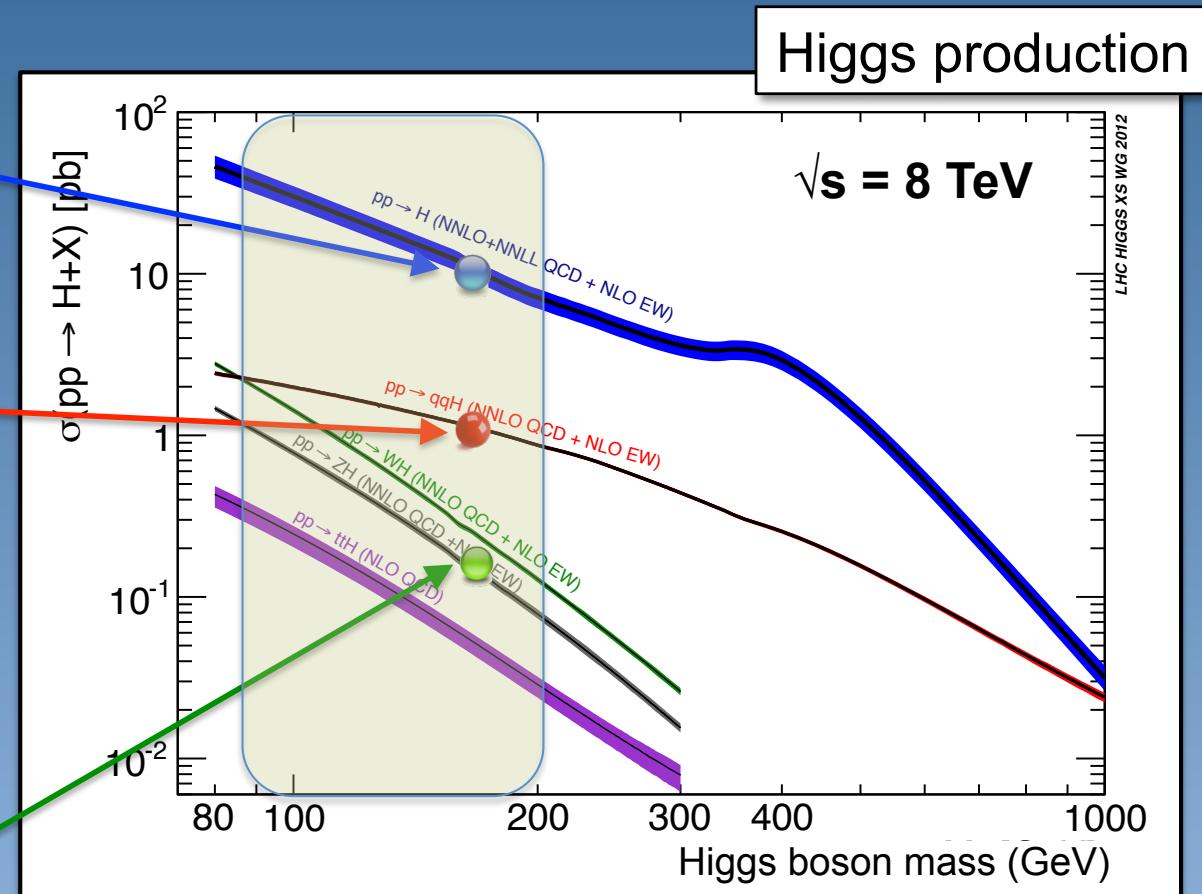
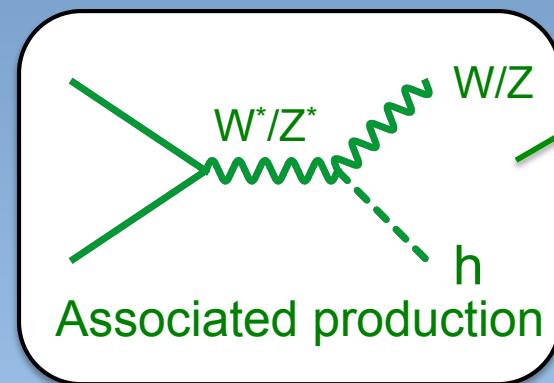
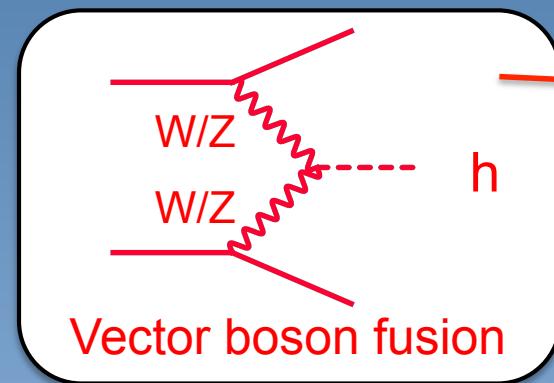
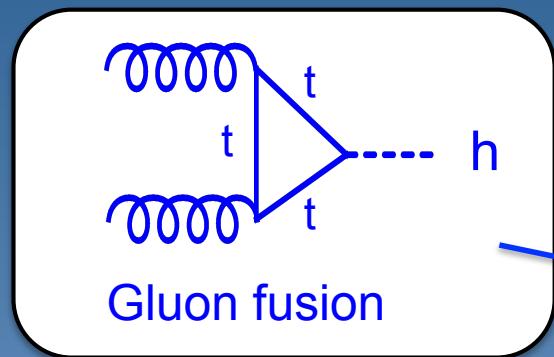
Fermions

*Massief fermion ?
... then it couples to the Higgs!*



Higgs boson production

Production of the Higgs boson



How many Higgs bosons have been produced at the LHC (until July 4th) ?



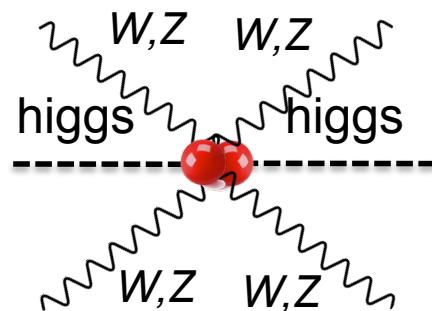
If the Higgs boson exists:

$$\begin{aligned} m_h = 125 \text{ GeV: } & 212.000 \\ m_h = 200 \text{ GeV: } & 77.000 \end{aligned}$$

How does the Higgs boson decay ?

Gauge bosons

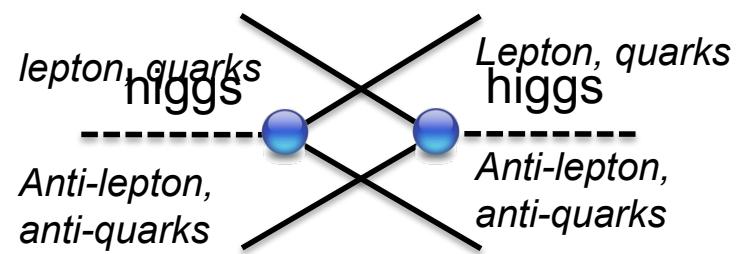
*Massive gauge boson ?
... then it coupling to the Higgs!*



$$\Gamma(h \rightarrow VV^*) \propto m_V^4$$

Fermions

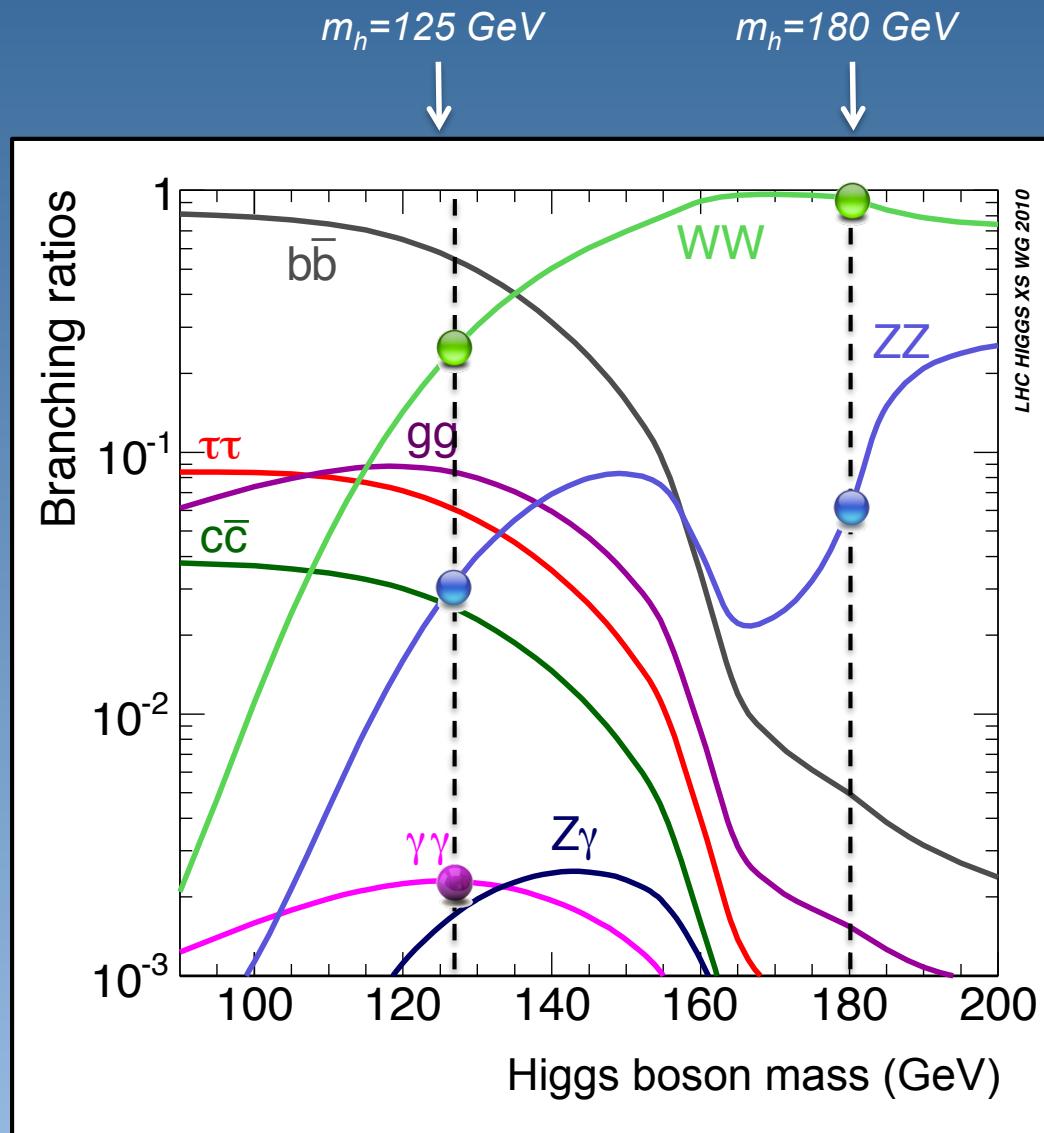
*Massief fermion ?
... then it couples to the Higgs!*



$$\Gamma(h \rightarrow f\bar{f}) \propto m_f^2$$

Higgs boson decay

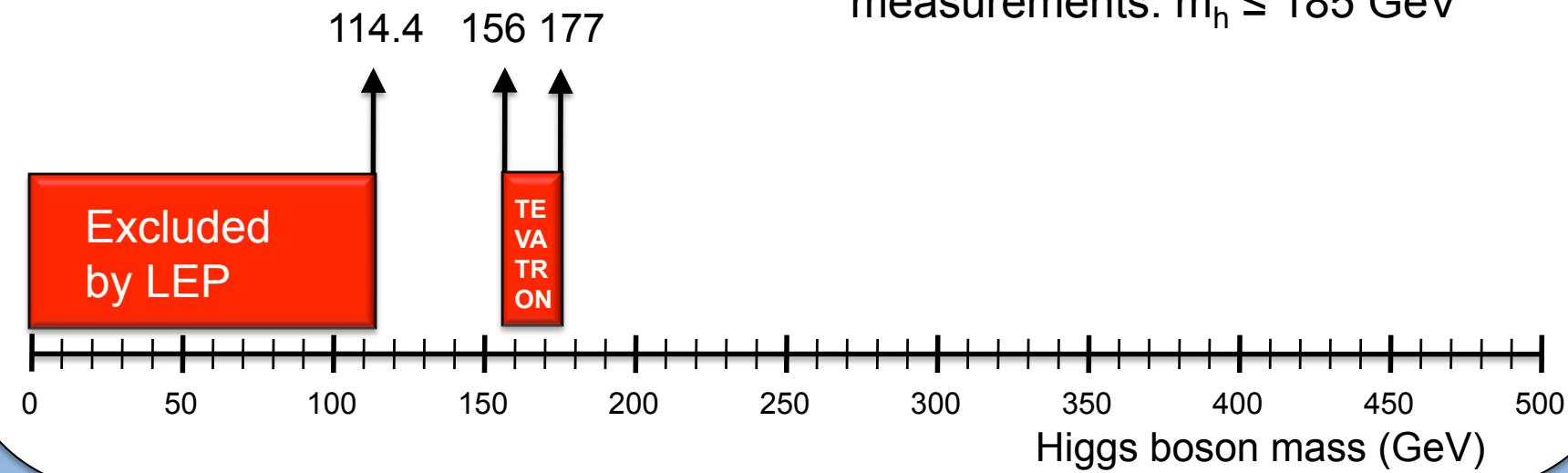
Decay of the Higgs boson



Subsequent gauge
boson decay

$W \rightarrow qq$	67%
$\ell\nu$	33%
$Z \rightarrow qq$	70%
$\nu\nu$	20%
WW	10%

Status beginning of last year



Indirect:

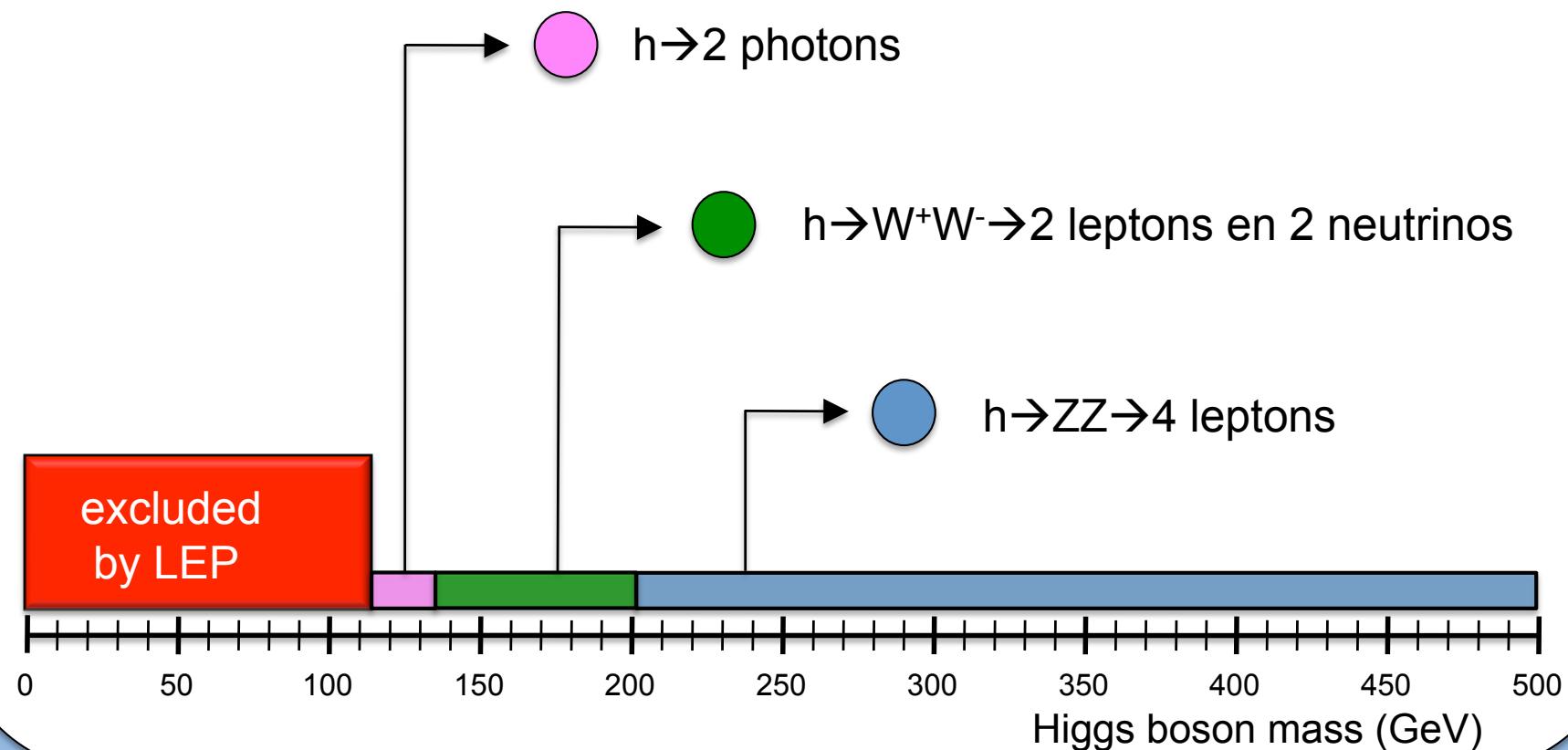
Theory: $m_h \leq 600$ GeV

Electroweak precision-
measurements: $m_h \leq 185$ GeV

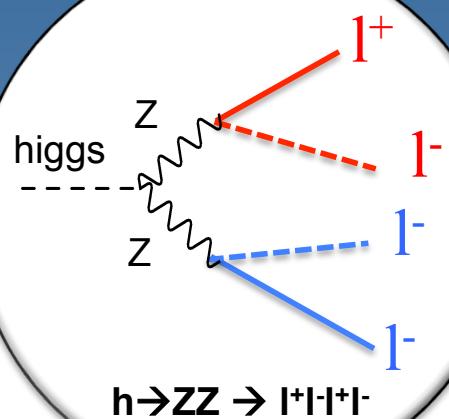
Experiments at the LHC Sensitive for mass range: 114-1000 GeV 45/78

The clearest fingerprint of the Higgs boson

For every mass here are multiple channels: these are most sensitive ones



High massa's: $m_H > 200$ GeV



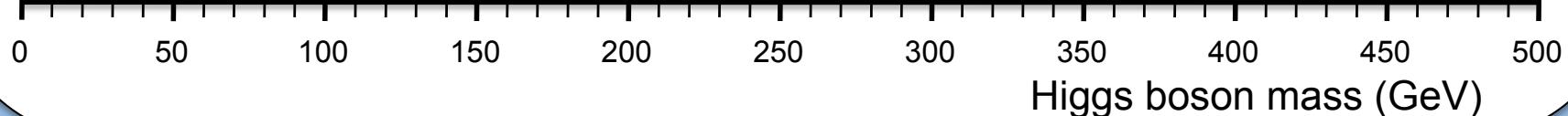
The golden channel
- 4 leptons (e/ μ)
- very rare: Higgs or ZZ

Fingerprint: 4-lepton mass



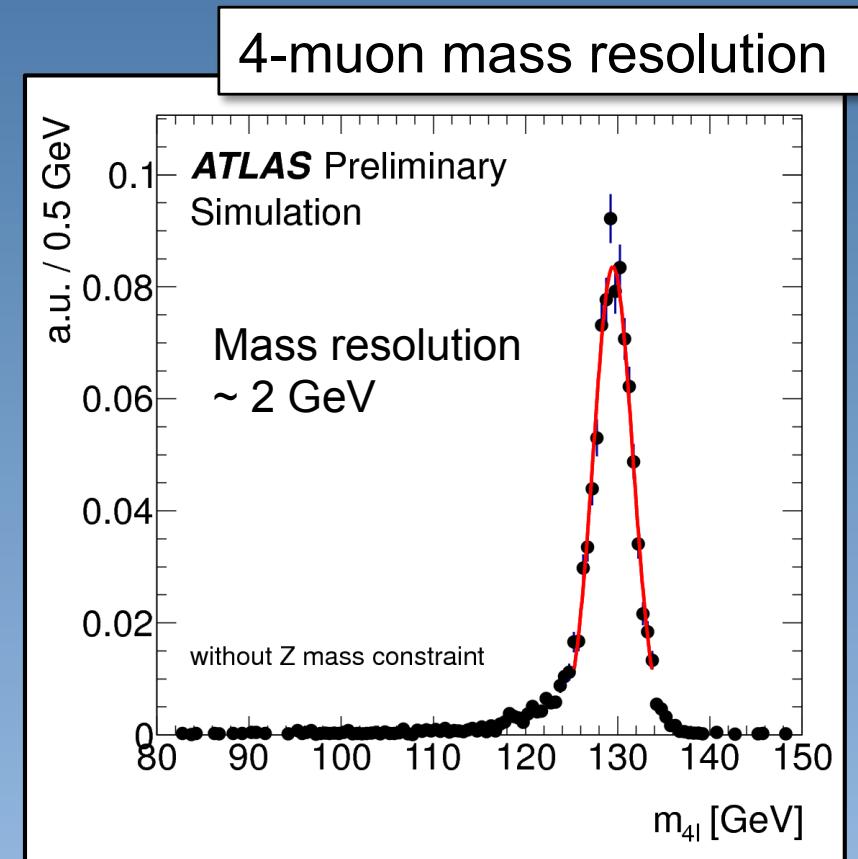
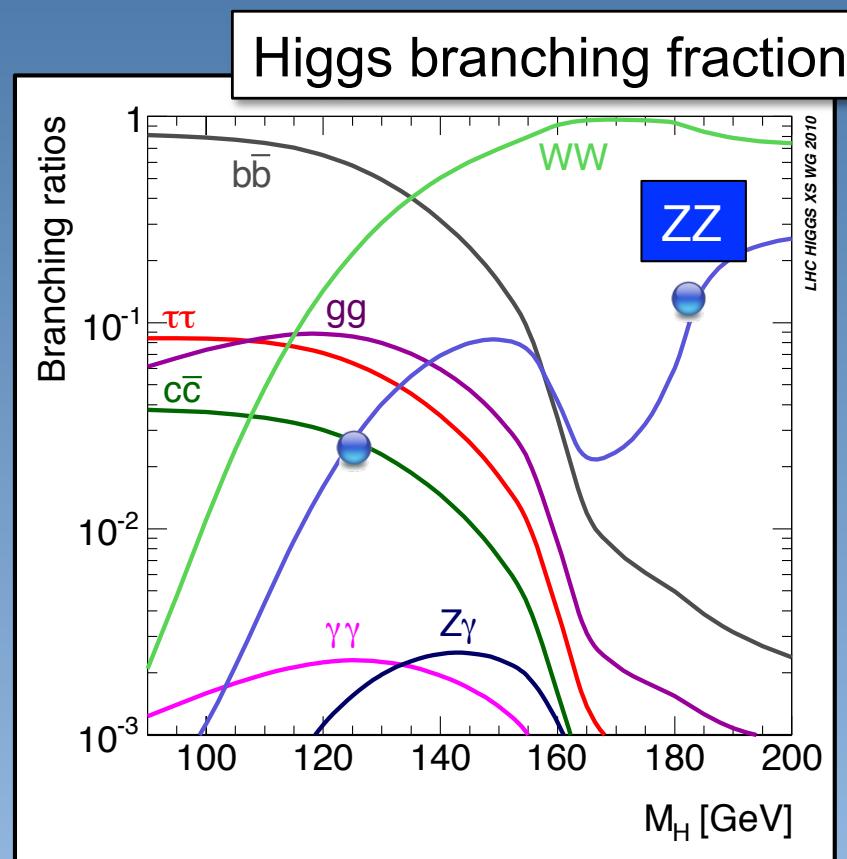
$h \rightarrow ZZ \rightarrow 4$ leptons

Excluded
by LEP

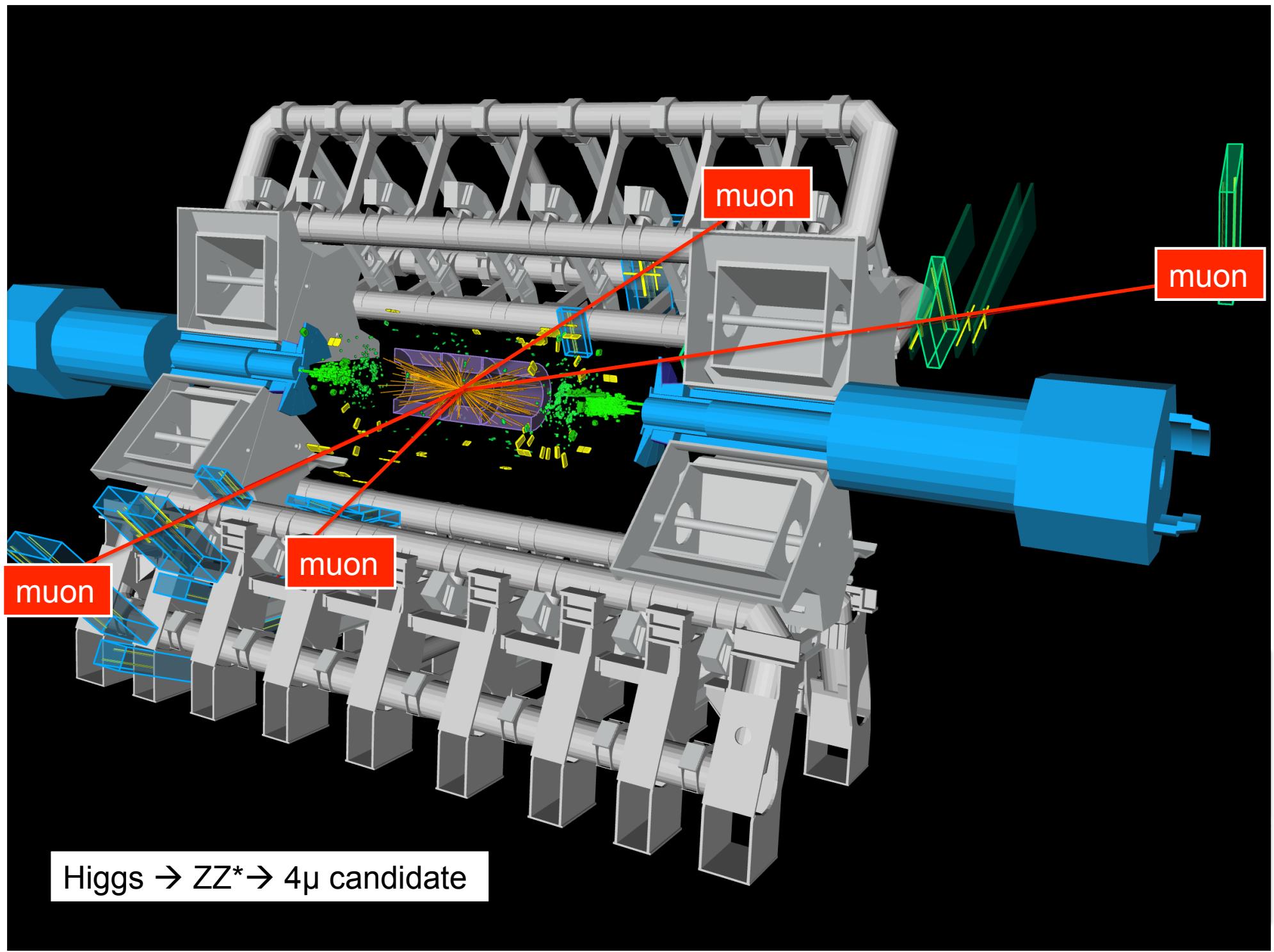


Higgs $\rightarrow ZZ \rightarrow 4$ leptons

Small number of beautiful collisions



Note: if $m_h = 125$ GeV: only 1 in 8,000 Higgs bosons results in a 4 lepton (lepton = electron or muon) final state



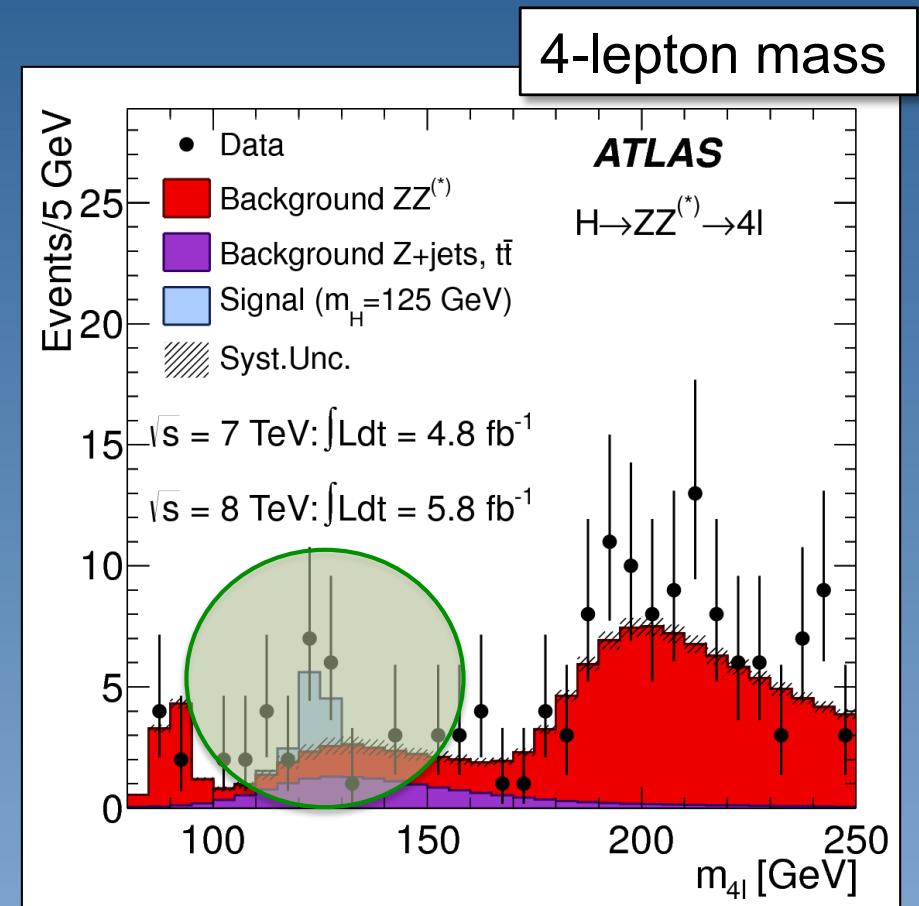
Higgs $\rightarrow ZZ \rightarrow 4$ leptons: small sample of beautiful events

Example: 125 GeV Higgs

1) Theory: $H \rightarrow ZZ \rightarrow 4l$ ($l=e, \mu$):

$$\begin{aligned} \Gamma(h \rightarrow ZZ) &= 2.64 \% \\ \Gamma(Z \rightarrow ll) &= 3.37 \% \end{aligned} \quad \left. \right\} 1.25 \cdot 10^{-4}$$

2) ATLAS detector: reconstruct 4μ
detector acceptance, Atlas muon
reconstruction and selection criteria



$m_{4l} = 125 \pm 5 \text{ GeV}$

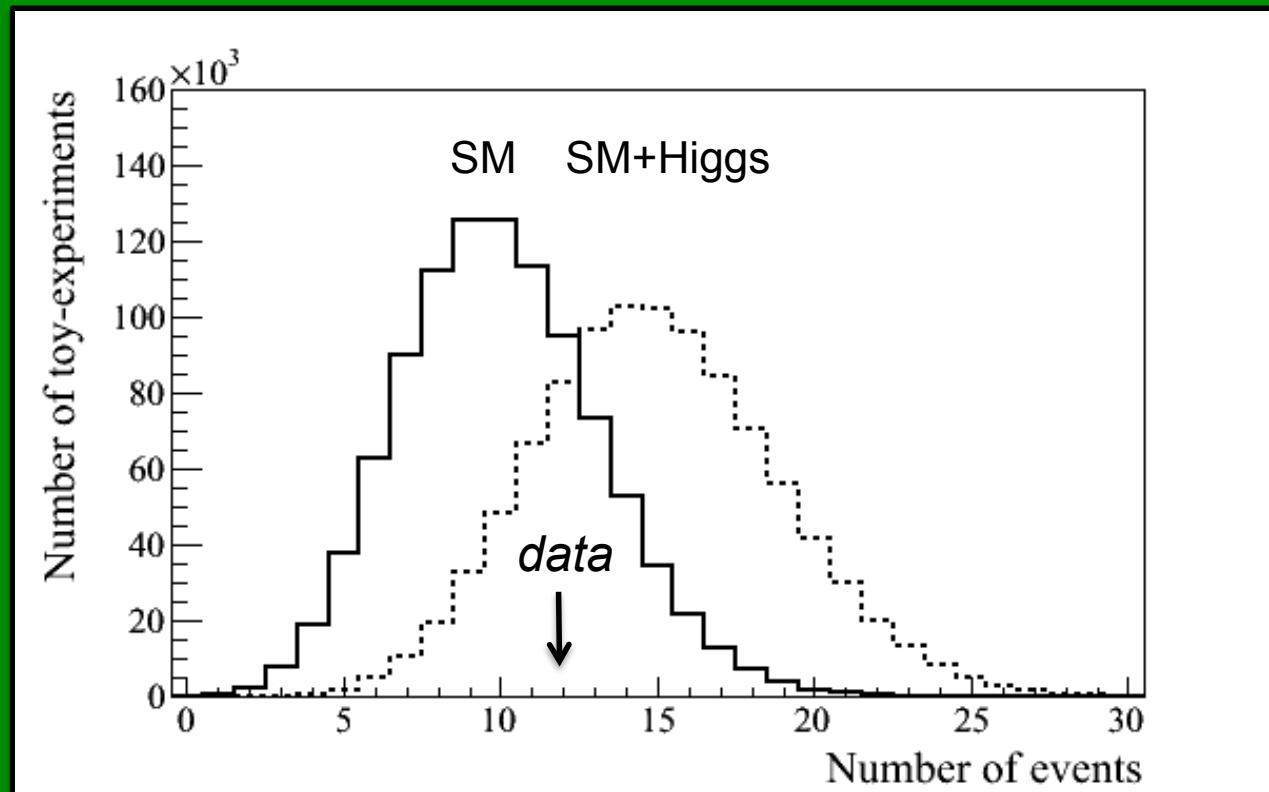
Expectation: SM (+Higgs)
Observation: Data

4.8 (10.1)
13

EXAMPLE: Quantifying compatibility with SM hypothesis

Measurement

SM	10
Higgs	5
Data	12



For a given number of events, the p-value is:

- 1) Expected excess (p-value)
- 2) Observed excess (p-value)

*"The probability to observe as many events as this number (or even more) under the assumption that the Higgs does **not** exists"*

→ Can be translated into a (gaussian) significance

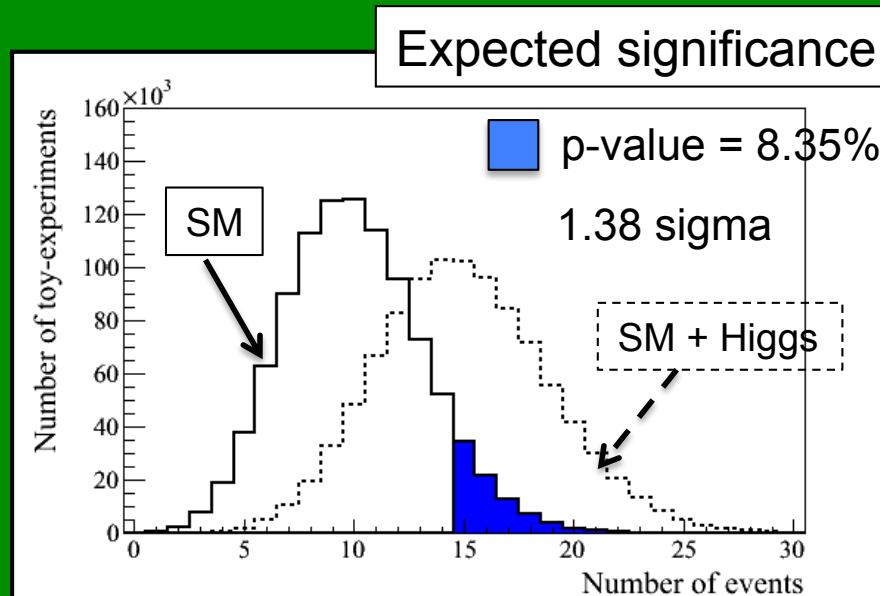
Quantifying an excess: p-value

Measurement

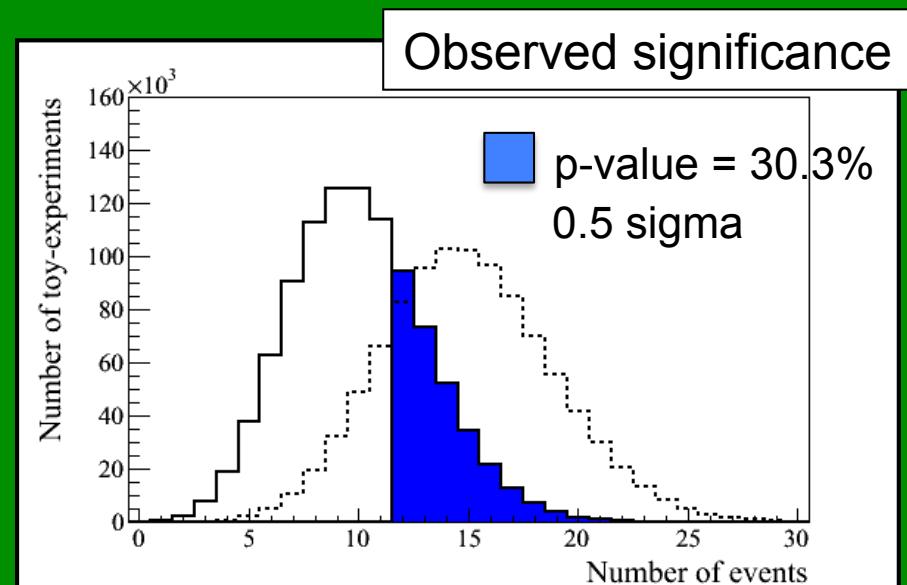
SM	10
Higgs	5
Data	12

p-value $< 2 \cdot 10^{-7}$ → discovery!

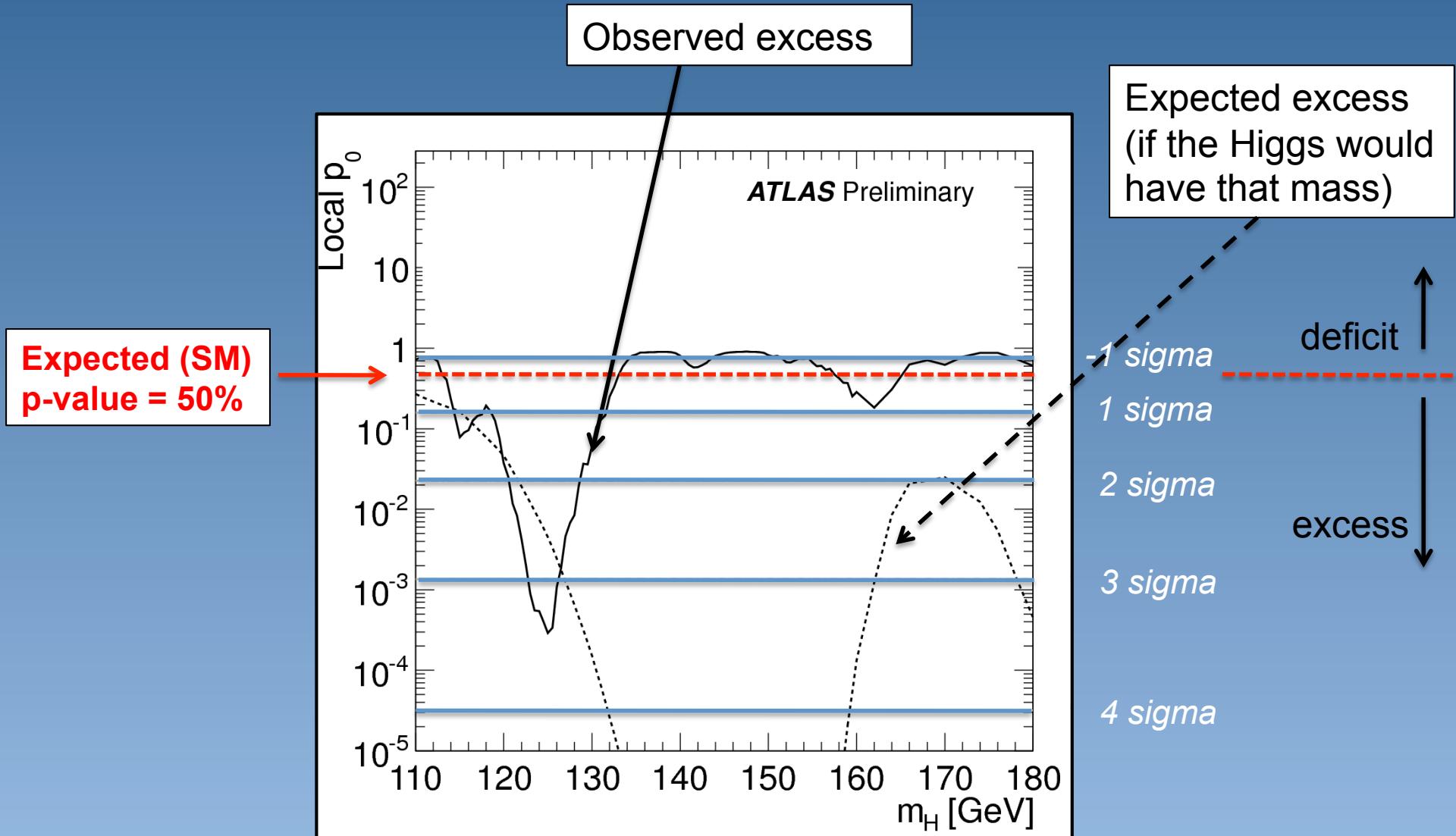
1) What is the expected significance ?



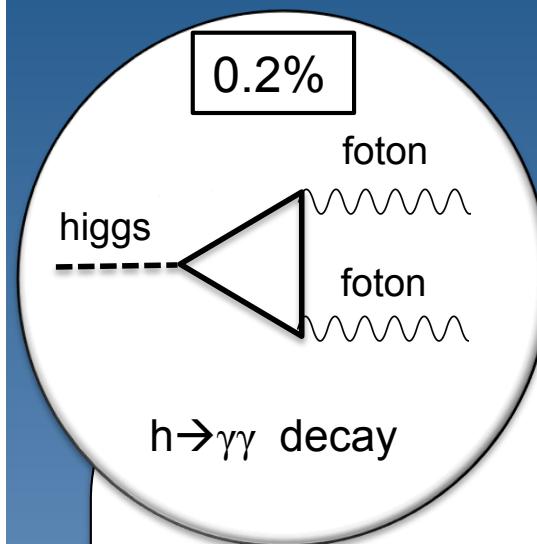
2) What is the observed significance ?



Quantifying the excess:



The low-mass region

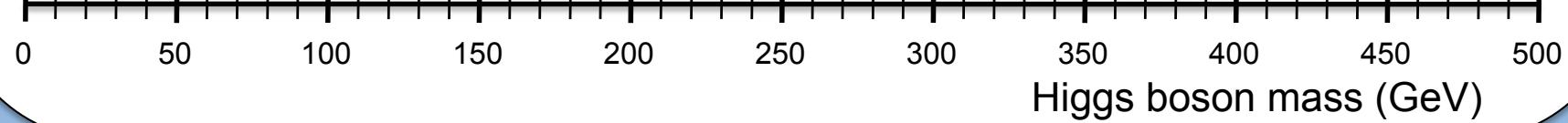


$h \rightarrow 2 \text{ photons}$

- 2 energetic isolated photons
- detector sensitive for 'fake' signals

Fingerprint: 2-photon mass

excluded
by LEP



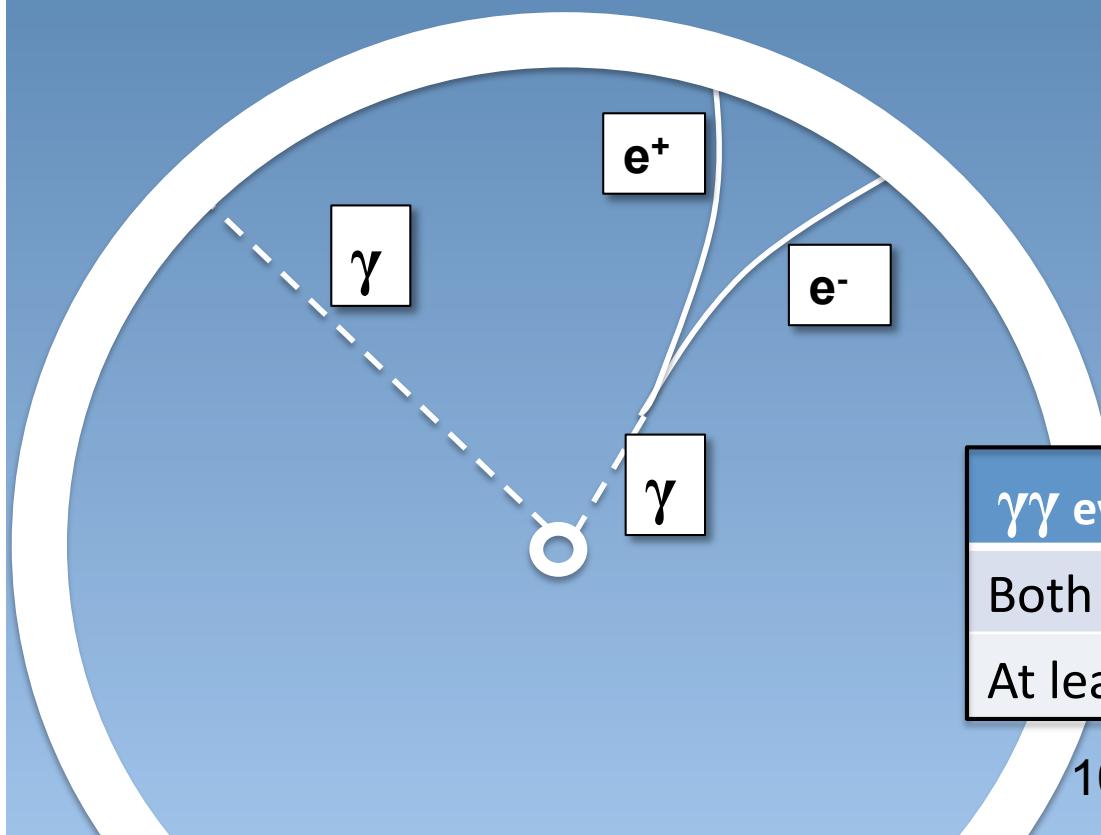
Photon identification and event selection

Selection: 2 isolated photons: $P_T > 40$ (30 GeV)

40% selection efficiency



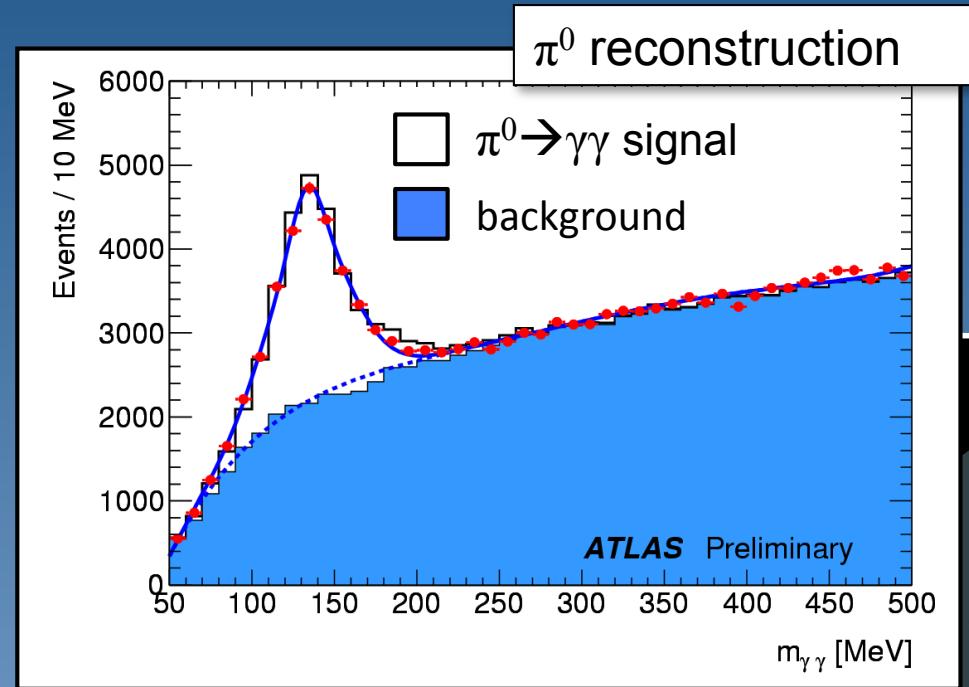
“Elk nadeel heet z’n voordeel”



$\gamma\gamma$ event classification	Fraction
Both unconverted	41%
At least 1 converted	59%

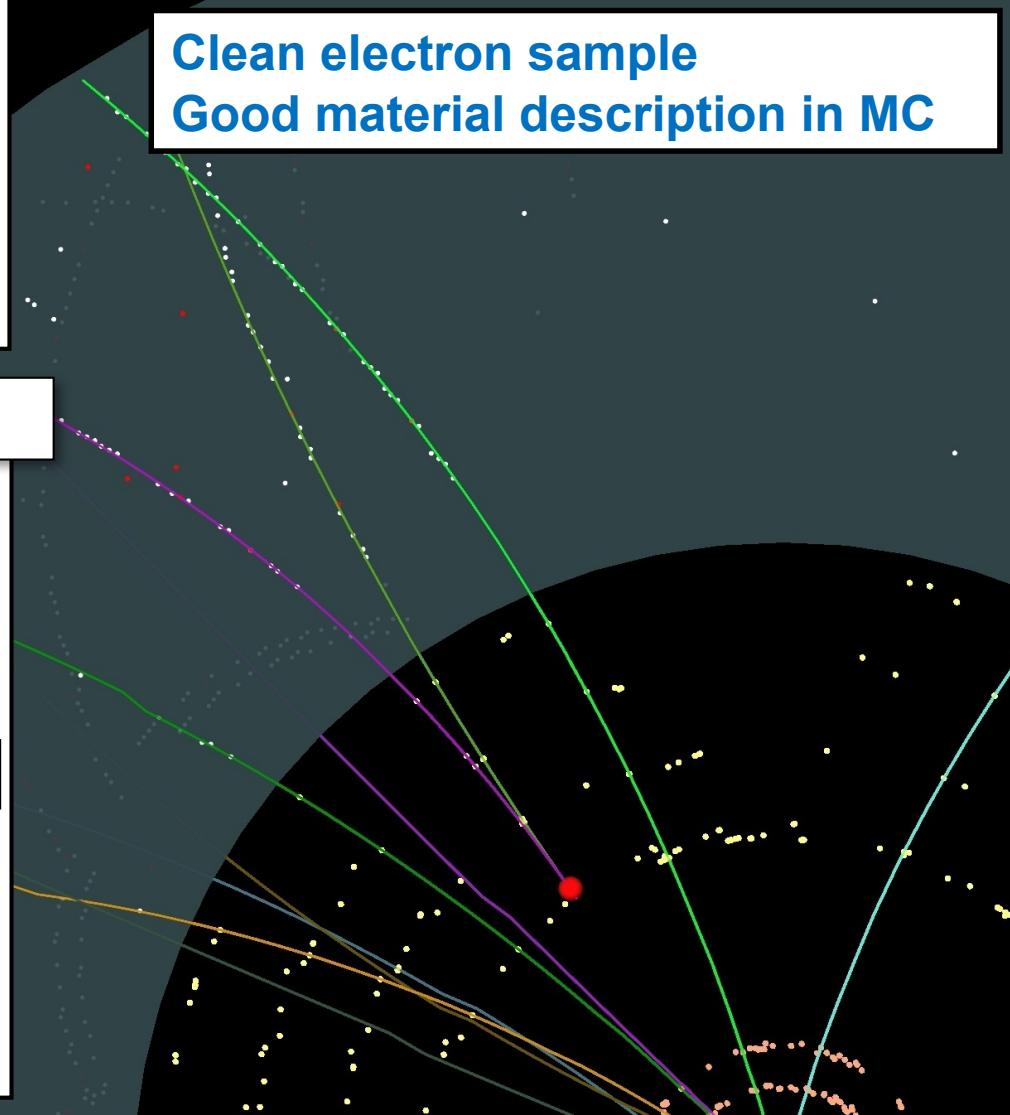
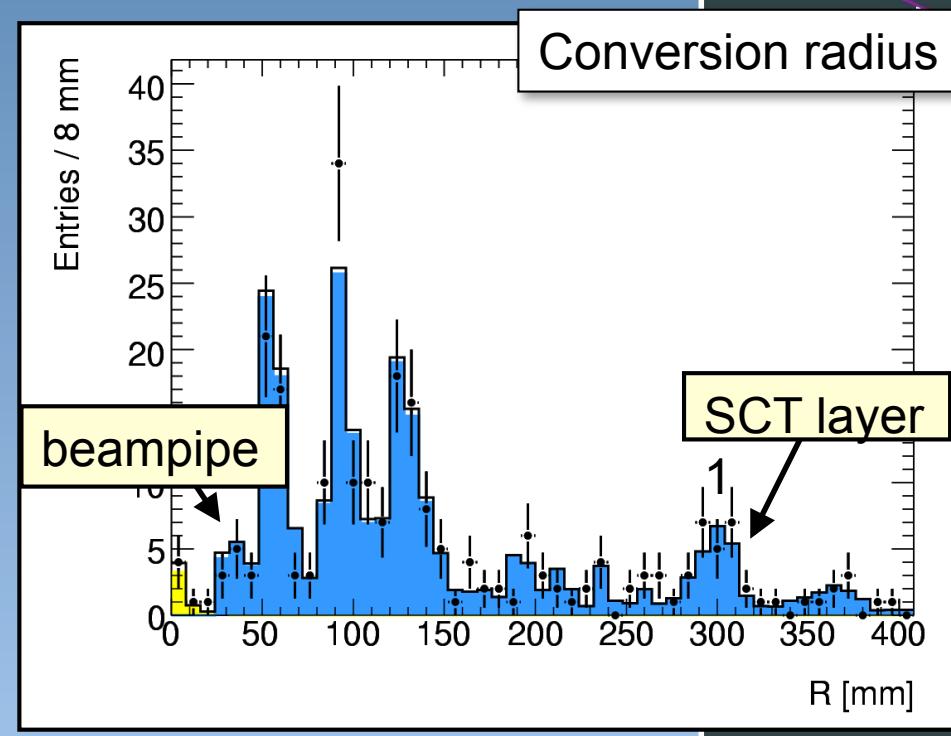
10 separate topology-categories

55/78



Photon conversions

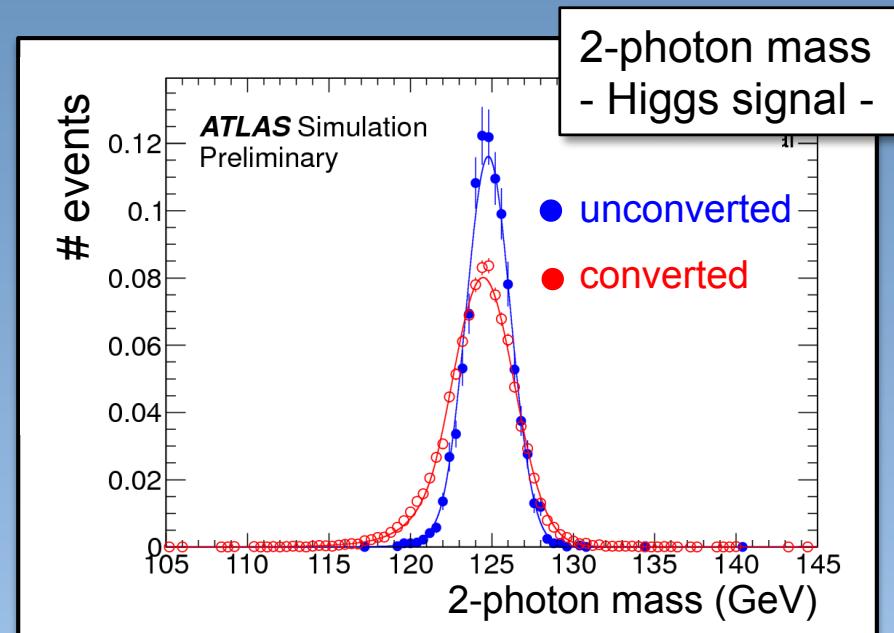
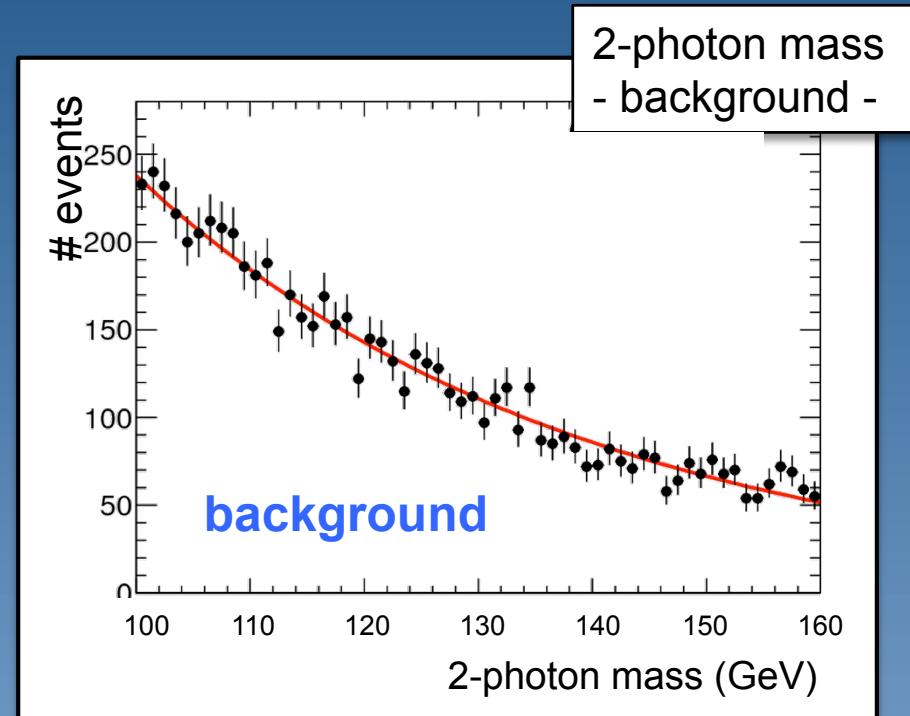
Clean electron sample
Good material description in MC



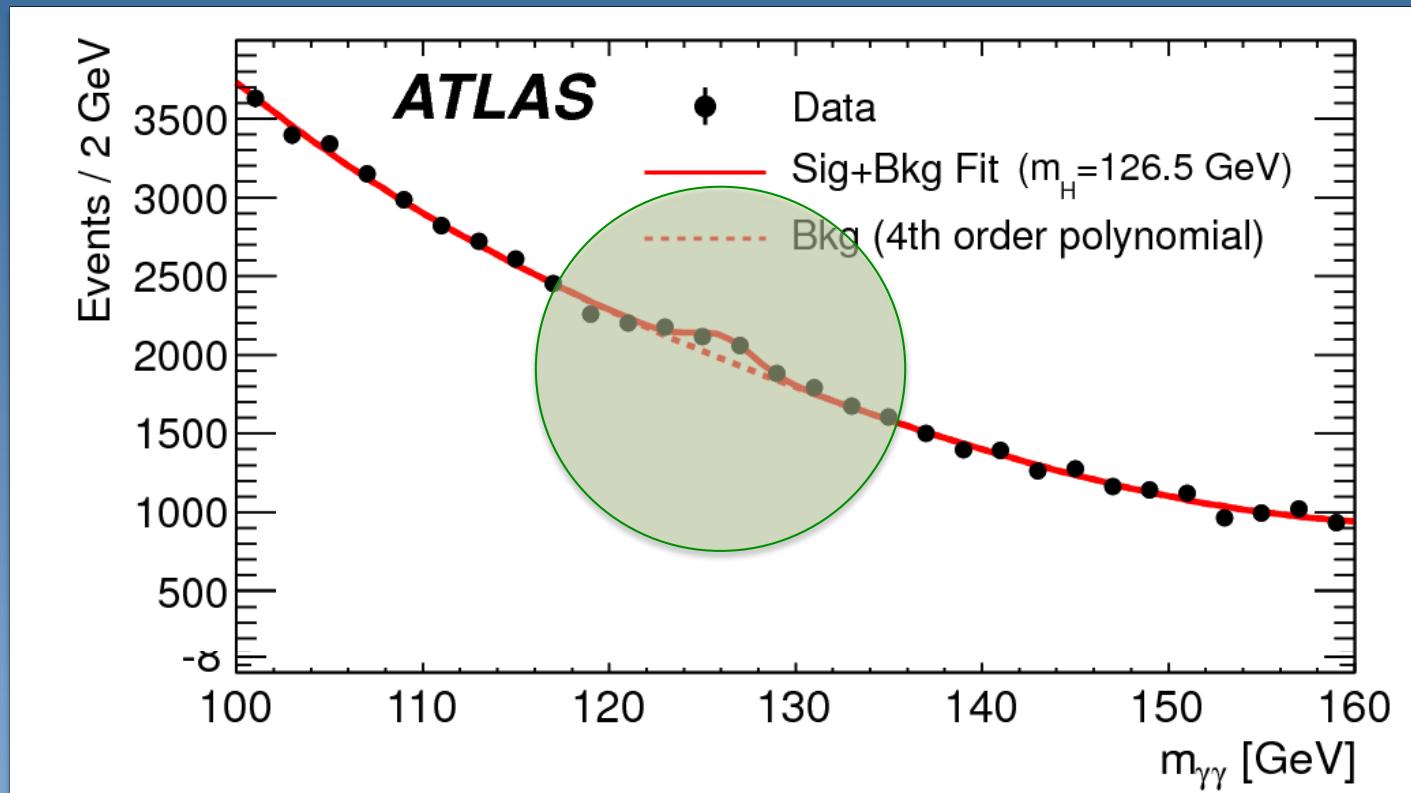
Di-photon events

Other origin
59,000 events

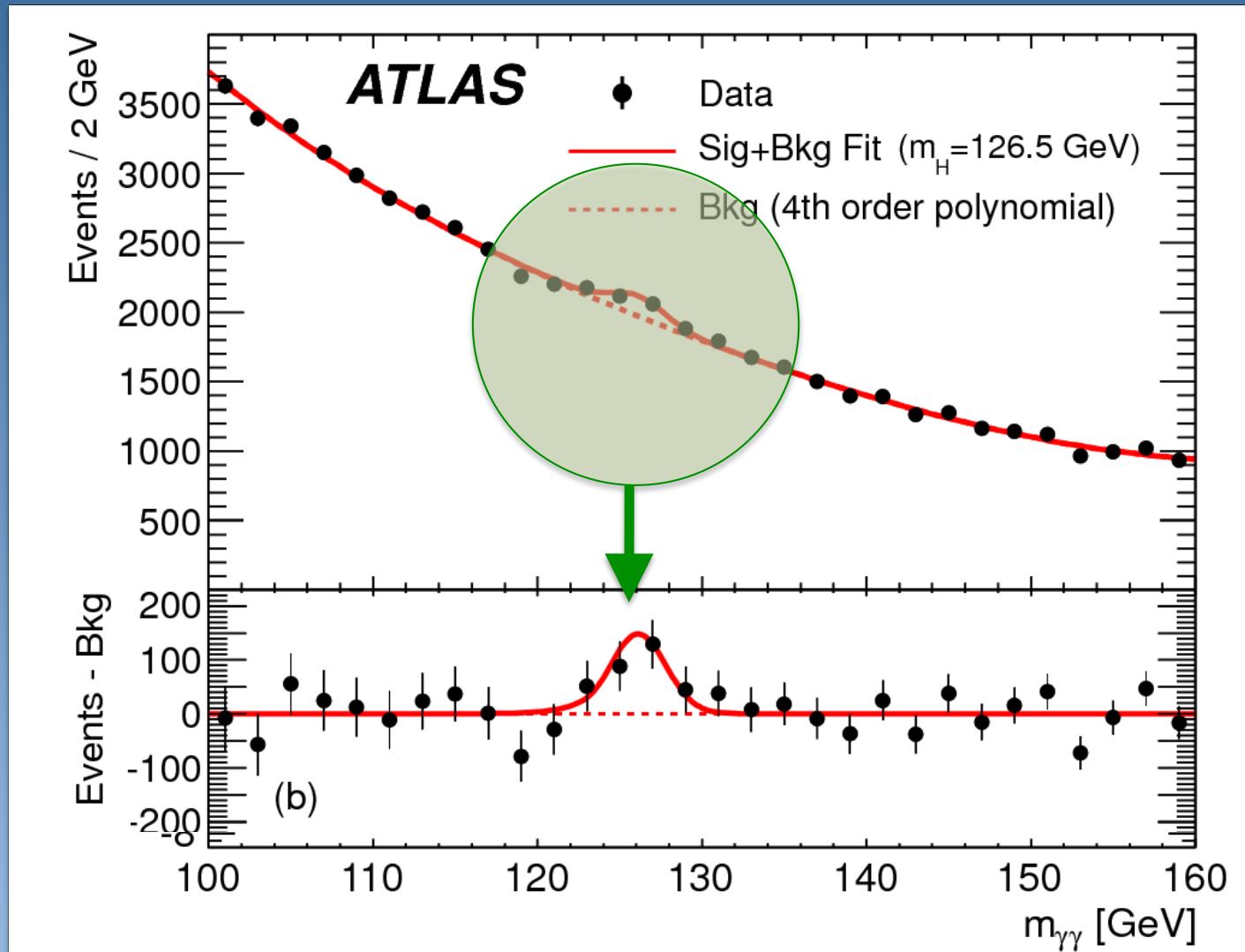
Higgs
190 events



The ATLAS data



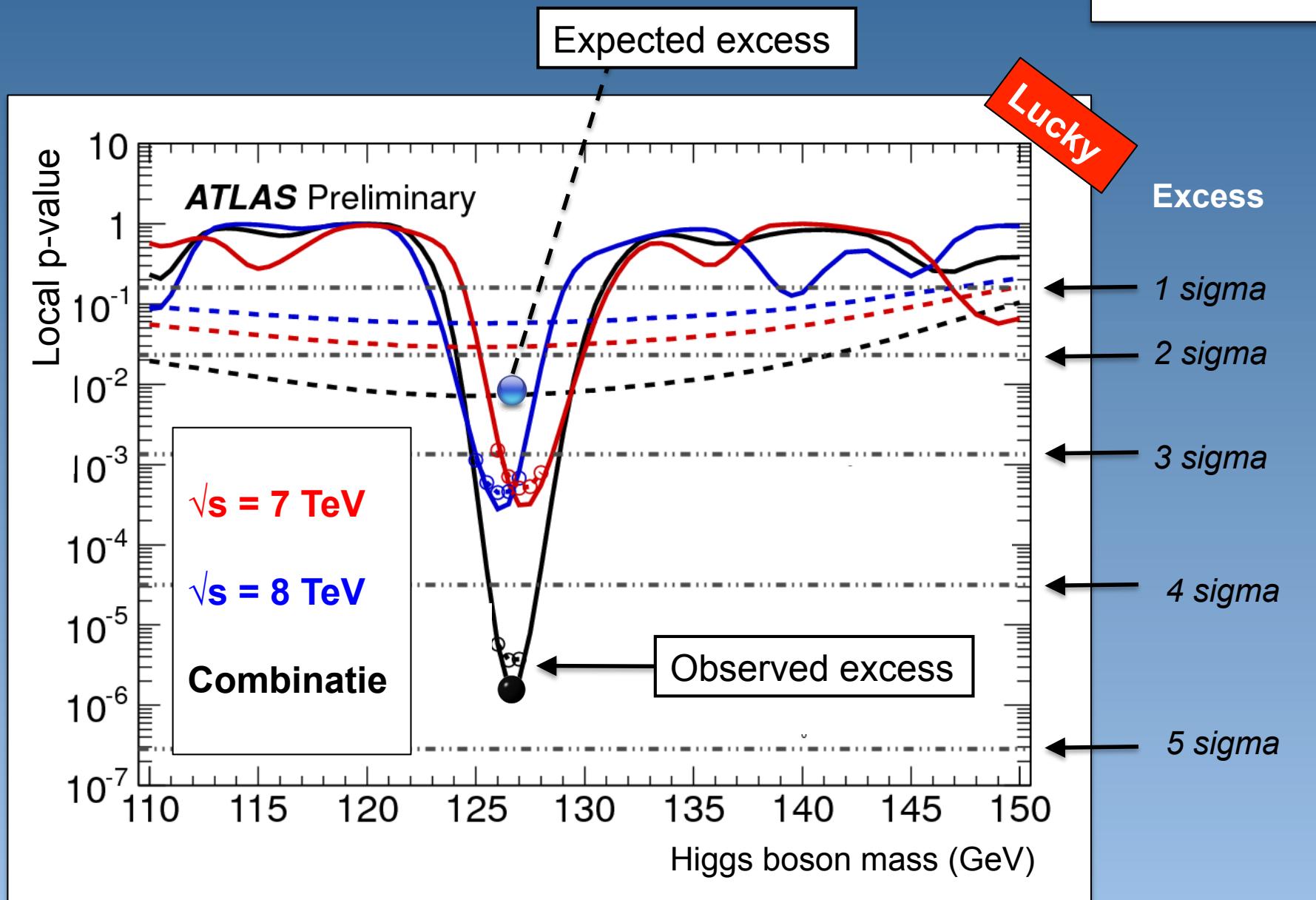
The ATLAS data



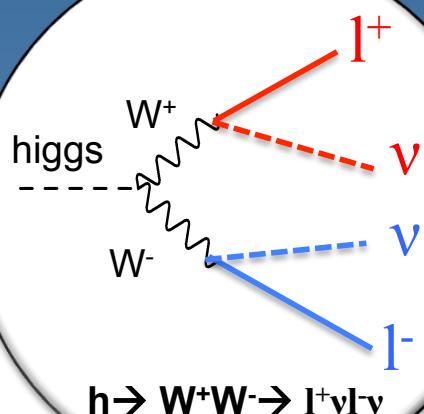
Quantify the significance ?

The excess in a single channel: $h \rightarrow \gamma\gamma$

expected
observed



Mid-range: $130 < m_H < 200$ GeV

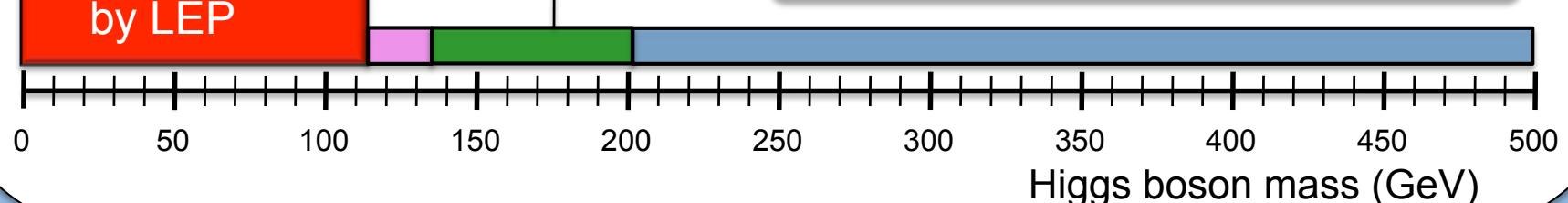


excluded
by LEP

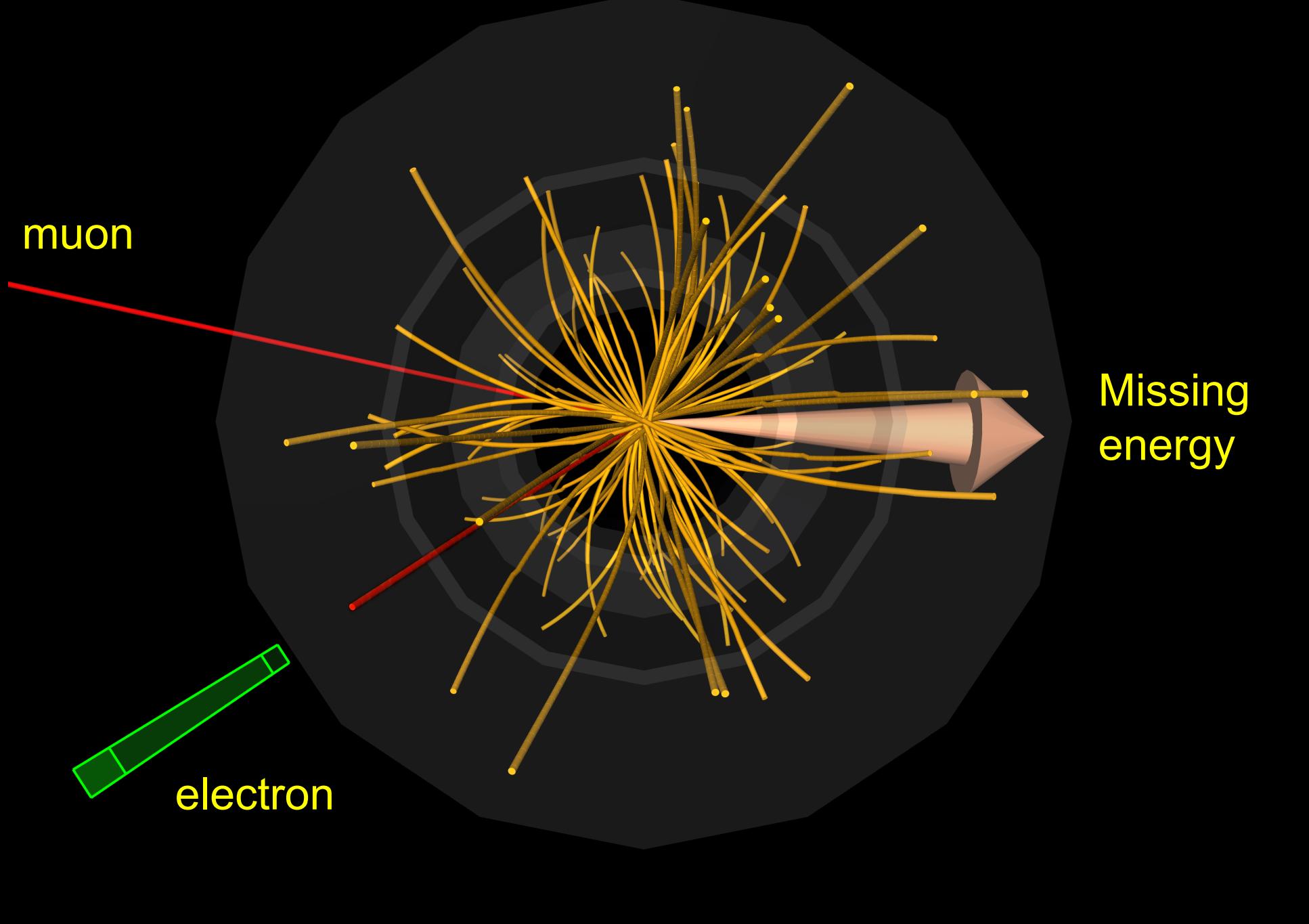
$h \rightarrow W^+ W^- \rightarrow 2 \text{ leptons en } 2 \text{ neutrinos}$

- 'Common' final state
- Difficult to isolate

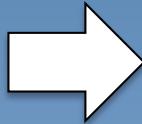
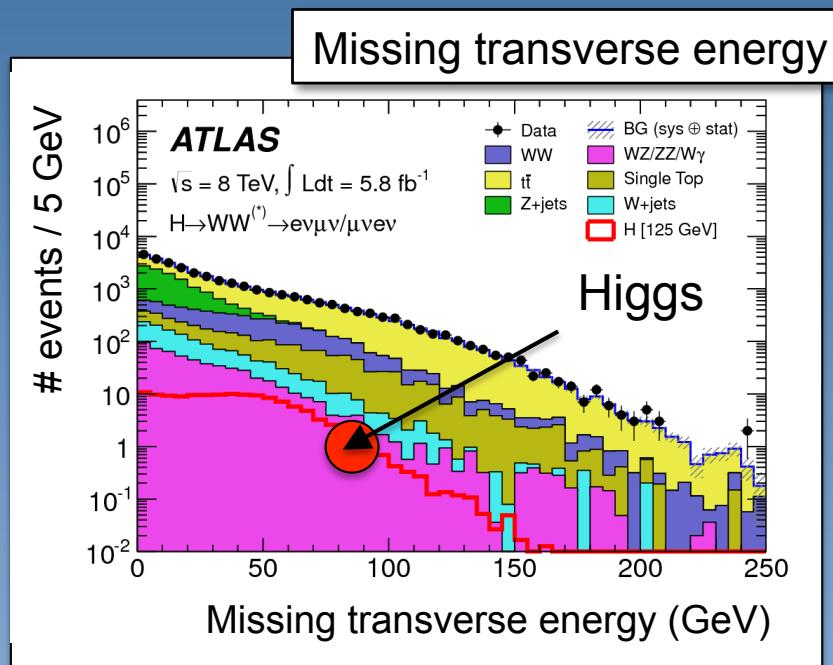
Fingerprint: none (no peak)



Final state: 2 leptons and 2 neutrino's + jets



The possible Higgs signal is well hidden



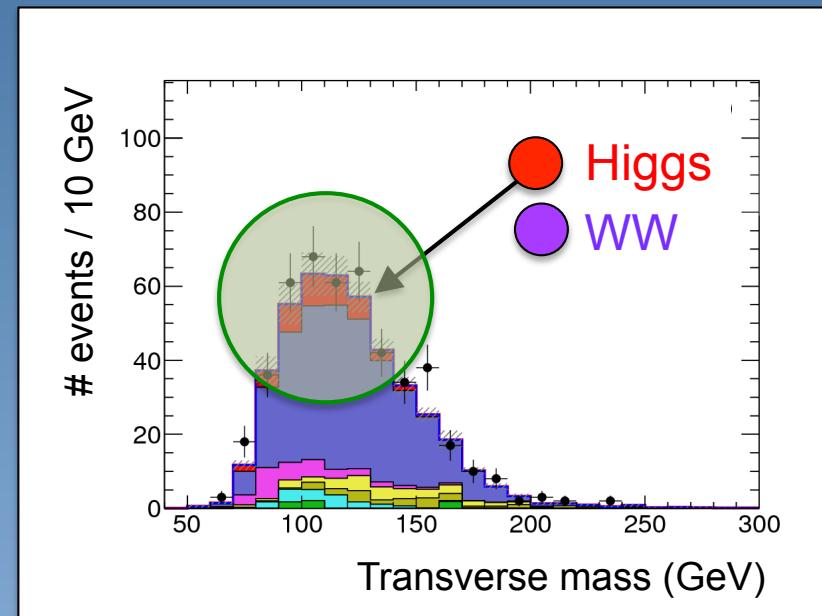
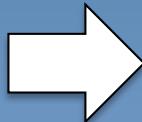
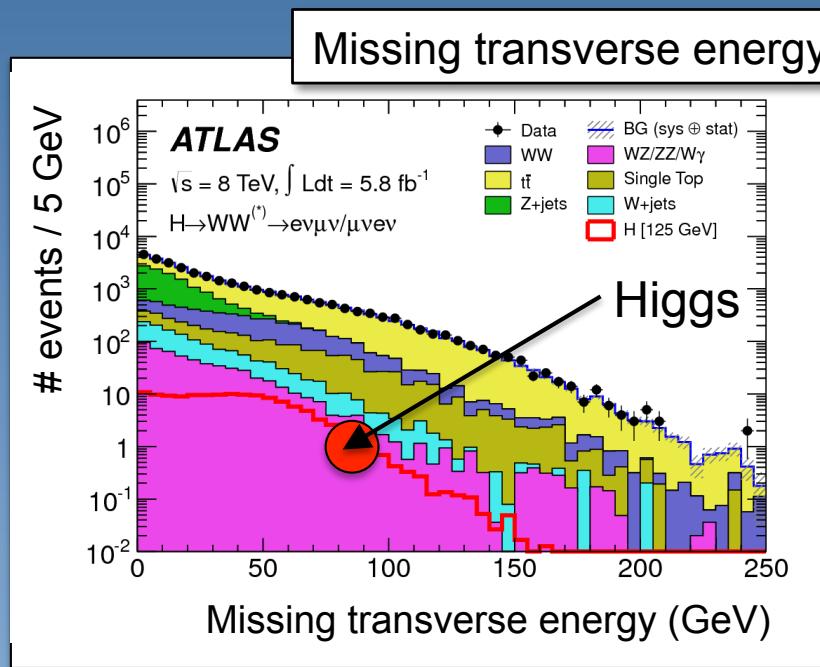
Delicate multi-dimensional selection

signal < 1/1,000 of background

Looking for diamonds



The possible Higgs signal is well hidden

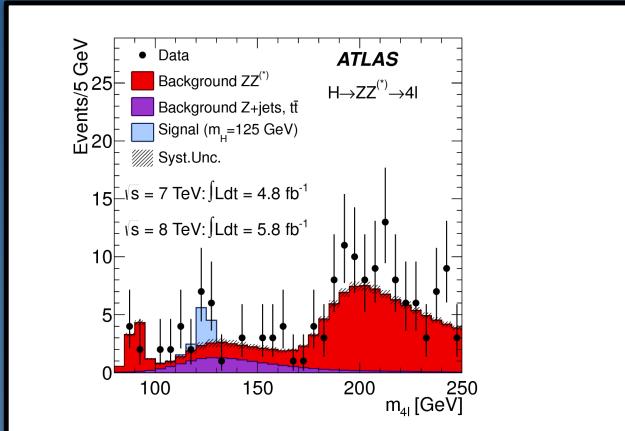


S/B $\sim 1/1,000$

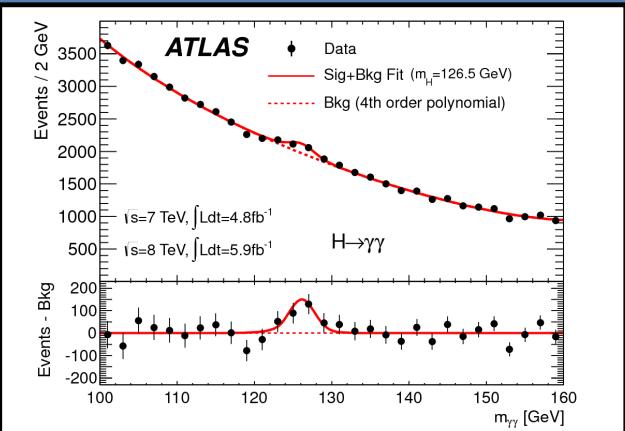
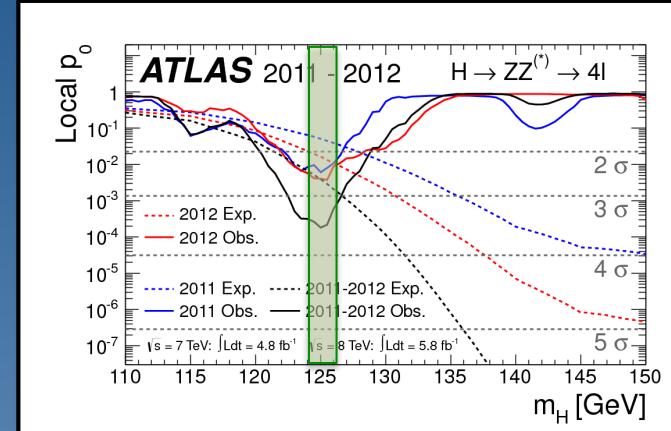
S/B $\sim 1/10$

Combine the evidence
 $h \rightarrow \gamma\gamma$, $h \rightarrow ZZ^*$ and $h \rightarrow WW^*$

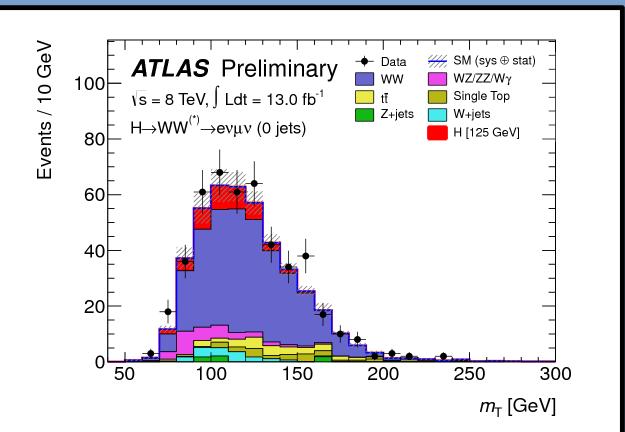
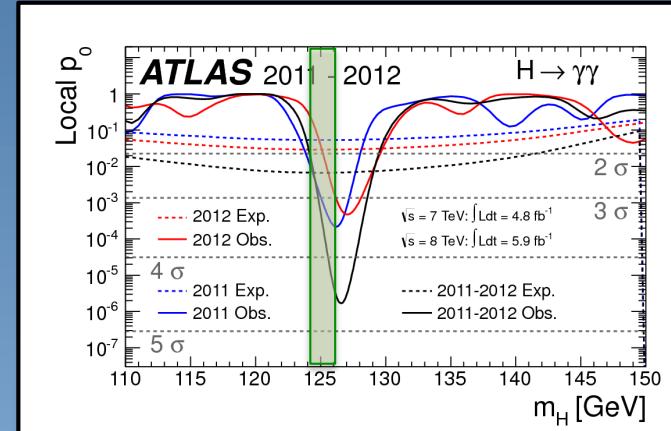




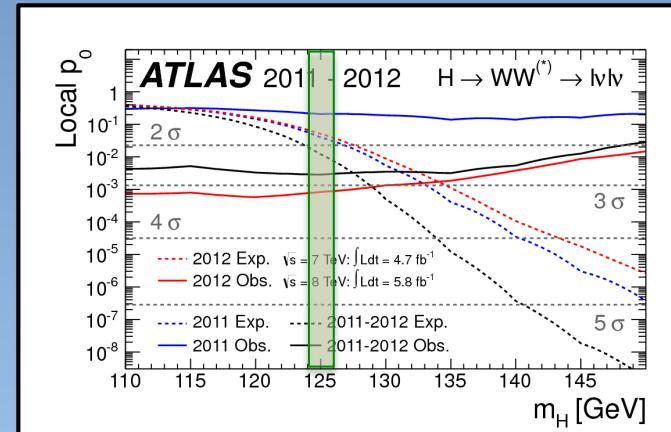
$H \rightarrow ZZ$



$H \rightarrow \gamma\gamma$

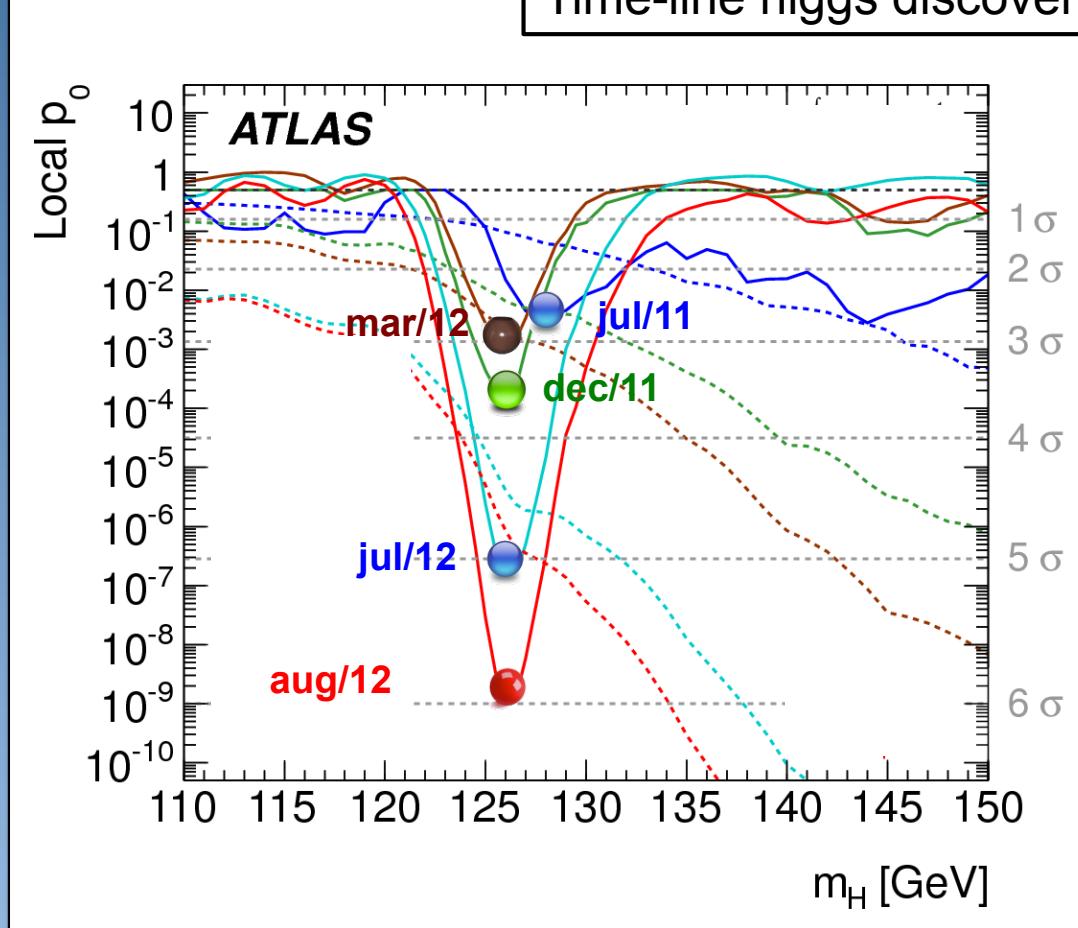


$H \rightarrow WW$

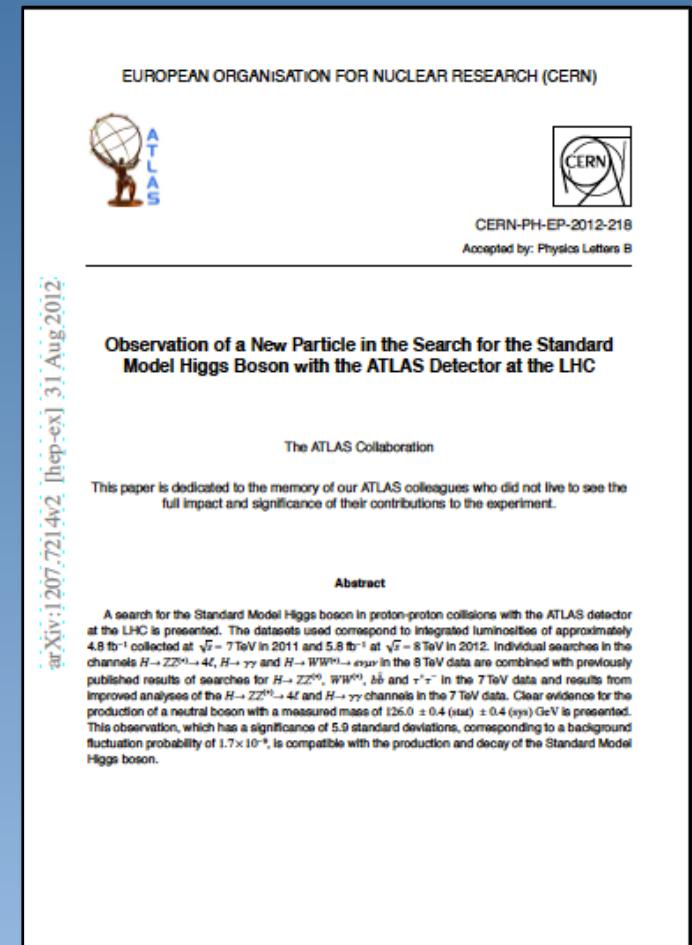


A discovery in slow-motion

Time-line higgs discovery



$$m_h = 126.0 \pm 0.4 \text{ (stat)} \pm 0.4 \text{ (syst)} \text{ GeV}$$



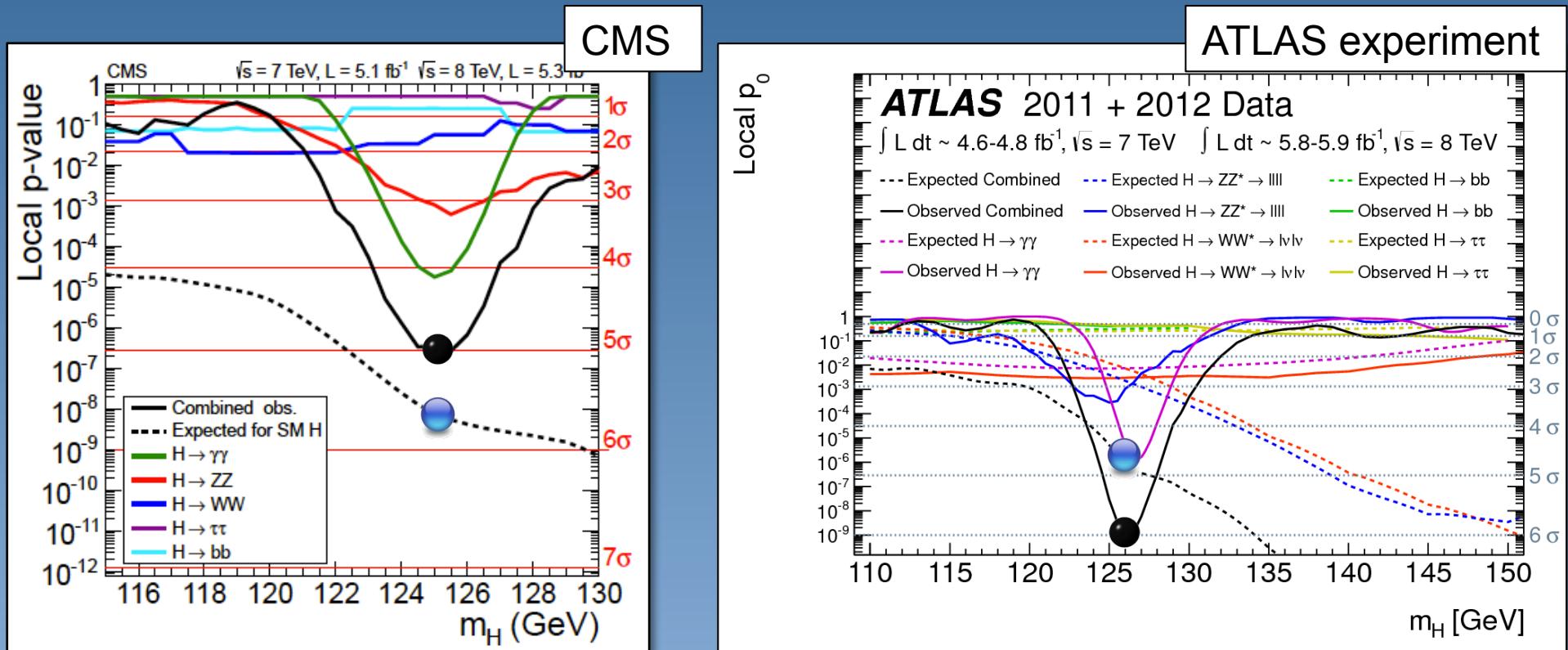
'Are you sure Ivo ?'

'Yes Diederik, we're pretty sure'

'Independent cross-check ?' ' Yep!'



Comparison results from ATLAS and CMS experiments



Conclusion (in July):

- expected
- observed

- 2 experiments observe new particle
- CMS more sensitive than ATLAS
... but ATLAS was lucky
- ATLAS had highest observed significance

Presentation CMS en ATLAS experiment: Higgs boson discovery



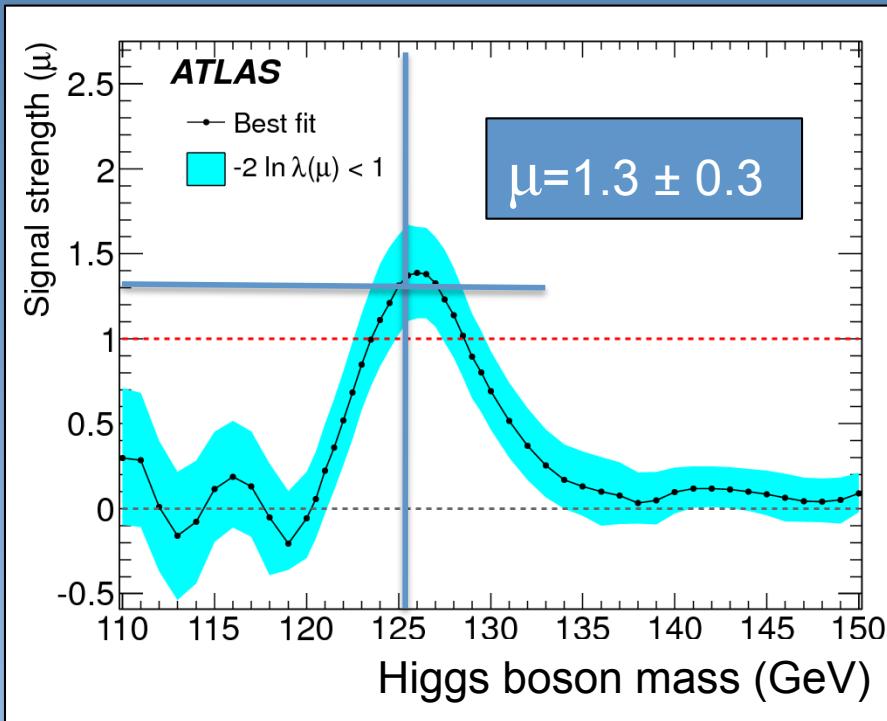
Is it the SM Higgs boson ?

SM: all properties are well predicted ... as a function of it's mass

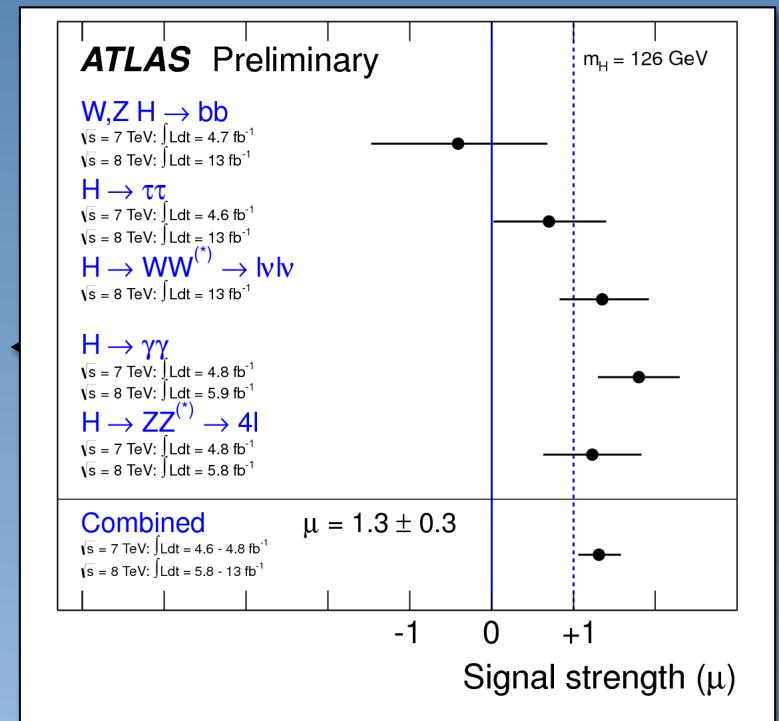
SuSy: 5 Higgs bosons, lightest one is like the SM Higgs boson,
... but slightly different

Looking at the Excess: signal strength (μ) compared to SM

Overall production cross-section scaling



Coupling to gauge bosons and fermions

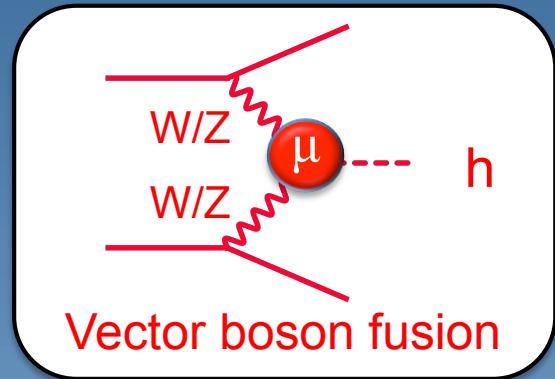


Properties: [couplings]

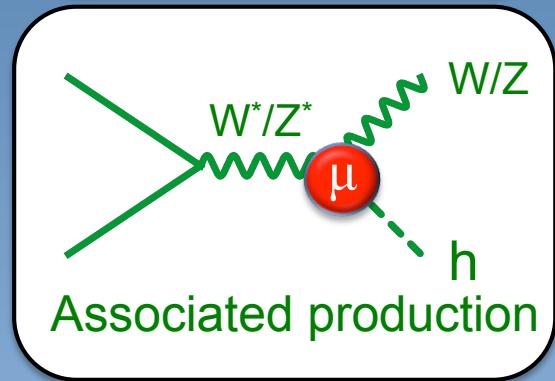
coupling to gauge bosons



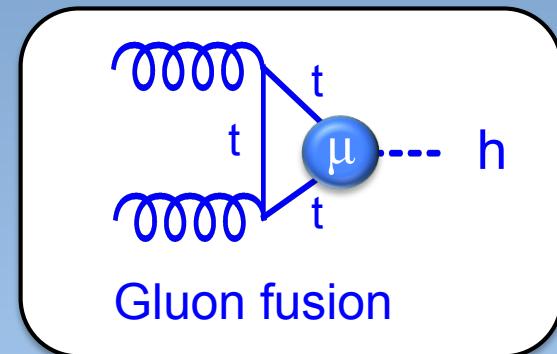
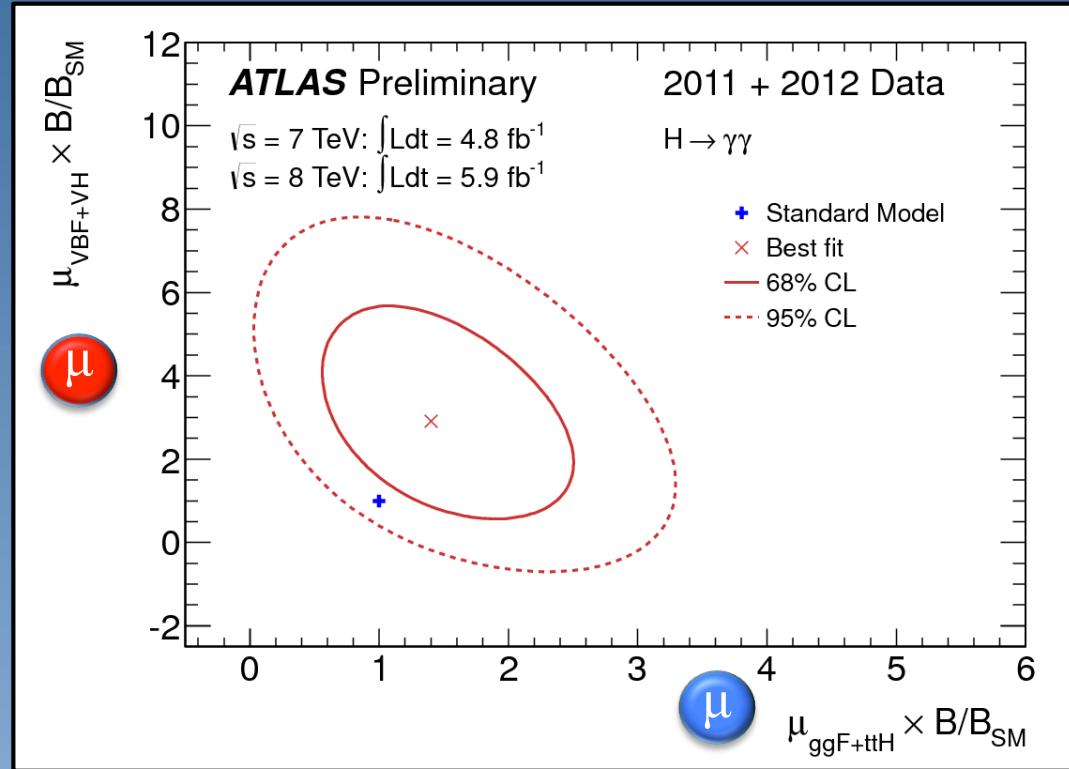
and fermions (Yukawa)



Vector boson fusion



Associated production



Gluon fusion

Properties: [spin/CP]

Prediction: the Higgs boson is a scalar particle

Spin 1 ? No!

- produced through gluon-gluon fusion (2x spin-1)
- decays to 2 photons (2x spin-1)

Spin-0 ? Spin-2

Spin/CP parameters: look at decay angles in $gg \rightarrow H \rightarrow ZZ \rightarrow 4 \text{ leptons}$

Not many plots/numbers public yet

Some ‘tiny’ issues in a world with a Higgs field:

- Radiative corrections (from top) to Higgs mass are huge $\propto \Lambda^2$. Extreme fine-tuning to get light Higgs.
Note: one of the motivations for super-symmetry
- Higgs does not give any hint to particle masses or dark matter
 $\rho_{Higgs}^{vac}(\Omega_\Lambda)$
- Contribution to factor 10^{54} off

Knowledge on mass Higgs boson

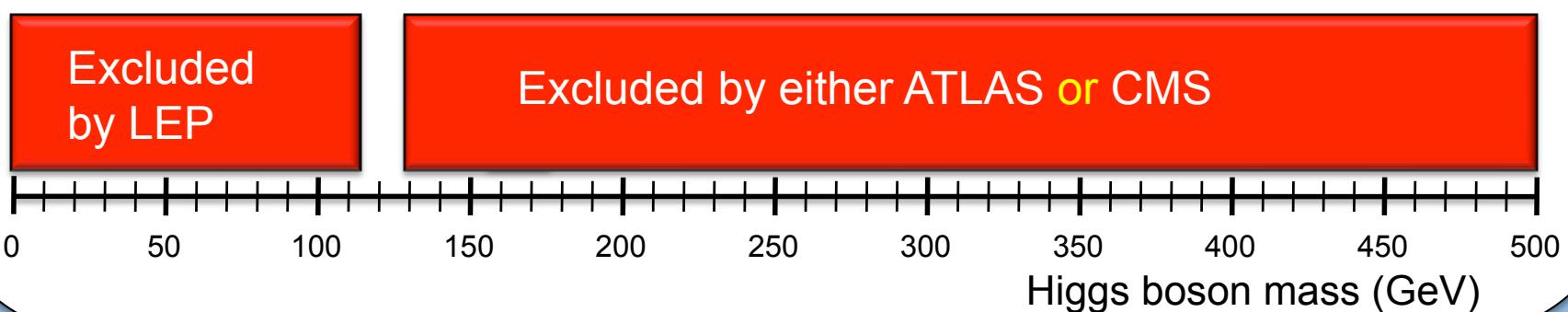
It's here



Indirect:

Theory: $m_h \leq 600$ GeV

Electroweak precision
measurement: $m_h \leq 185$ GeV



Summary



- 1) Large Hadron Collider and experiments
are working fan-tas-tic!
- 2) Search for the Higgs boson in 2011+2012
Discovery !!
- 3) Standard Model has serious shortcomings
Looking forward to new physics (SuSy, extra dimensions, ...)

All ATLAS public results: papers, numbers and plots
<https://twiki.cern.ch/twiki/bin/view/AtlasPublic>

backup

REFERENCES

ATLAS experiment

All ATLAS public results: publications, numbers and plots

<https://twiki.cern.ch/twiki/bin/view/AtlasPublic>

Higgs paper: “*Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC*”, Phys. Lett. B716 (2012) or <http://arxiv.org/abs/1207.7214>

CMS experiment

All CMS public results: publications, numbers and plots

<https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResults>

Higgs paper: “*Observation of a new boson at a mass of 125 GeV with the CMS experiment at the LHC*”, Phys. Lett. B716 (2012) or <http://arxiv.org/abs/1207.7235>