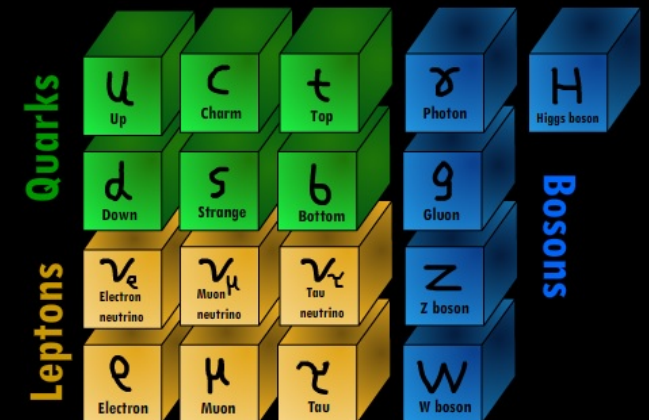




PHY3004: Nuclear and Particle Physics
Marcel Merk, Jacco de Vries, ...



The Standard Model

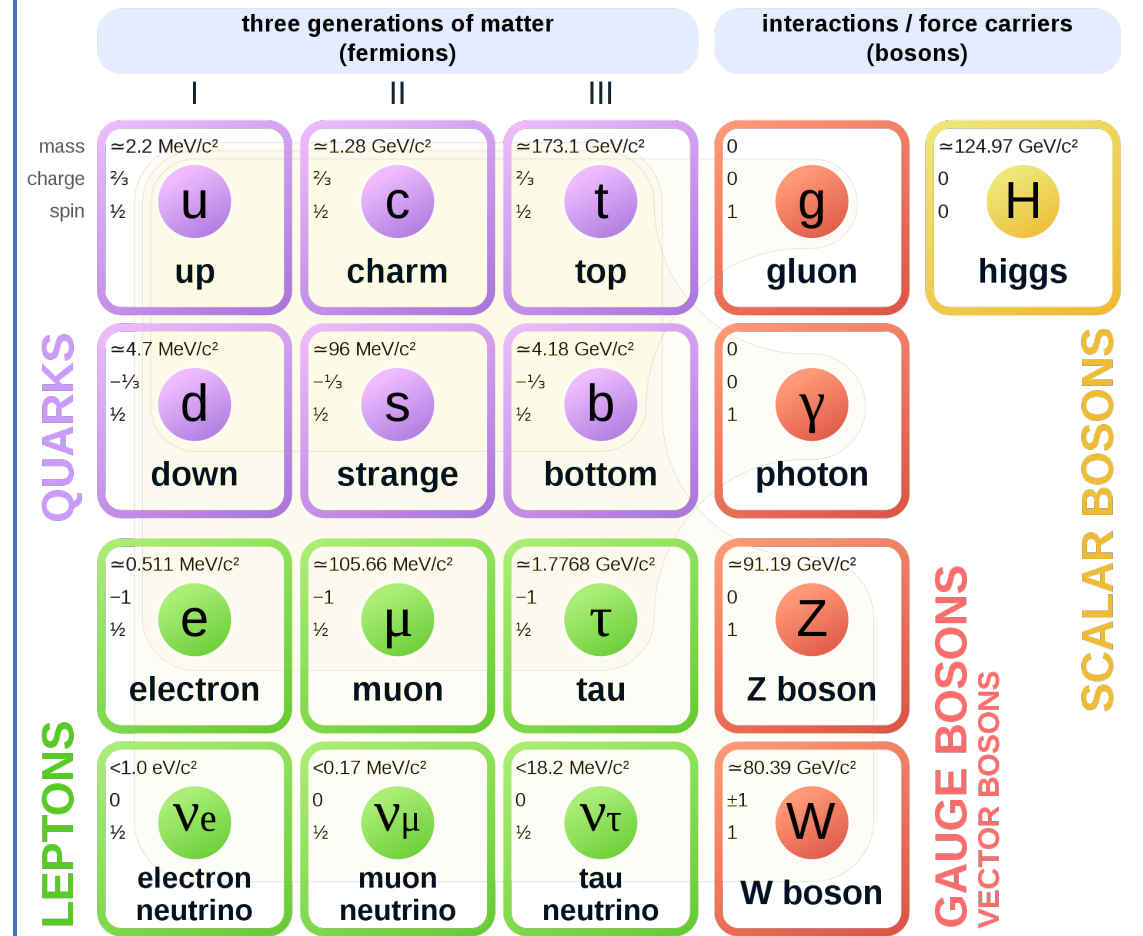


Elementary particles

Classification of particles

- **Lepton**: fundamental particle
- **Hadron**: consist of **quarks**
 - **Meson**: 1 quark + 1 antiquark (π^+ , B_s^0 , ...)
 - **Baryon**: 3 quarks (p , n , Λ , ...)
 - **Anti-baryon**: 3 anti-quarks
- **Fermion**: particle with half-integer spin.
 - Antisymmetric wave function: obeys Pauli-exclusion principle and Pauli-Dirac statistics
 - All fundamental quarks and leptons are spin- $\frac{1}{2}$
 - Baryons ($S=\frac{1}{2}$, $\frac{3}{2}$)
- **Boson**: particle with integer spin
 - Symmetric wave function: Bose-Einstein statistics
 - **Mesons**: ($S=0$, 1), **Higgs** ($S=0$)
 - **Force carriers**: γ , W , Z , g ($S=1$); graviton ($S=2$)

Standard Model of Elementary Particles

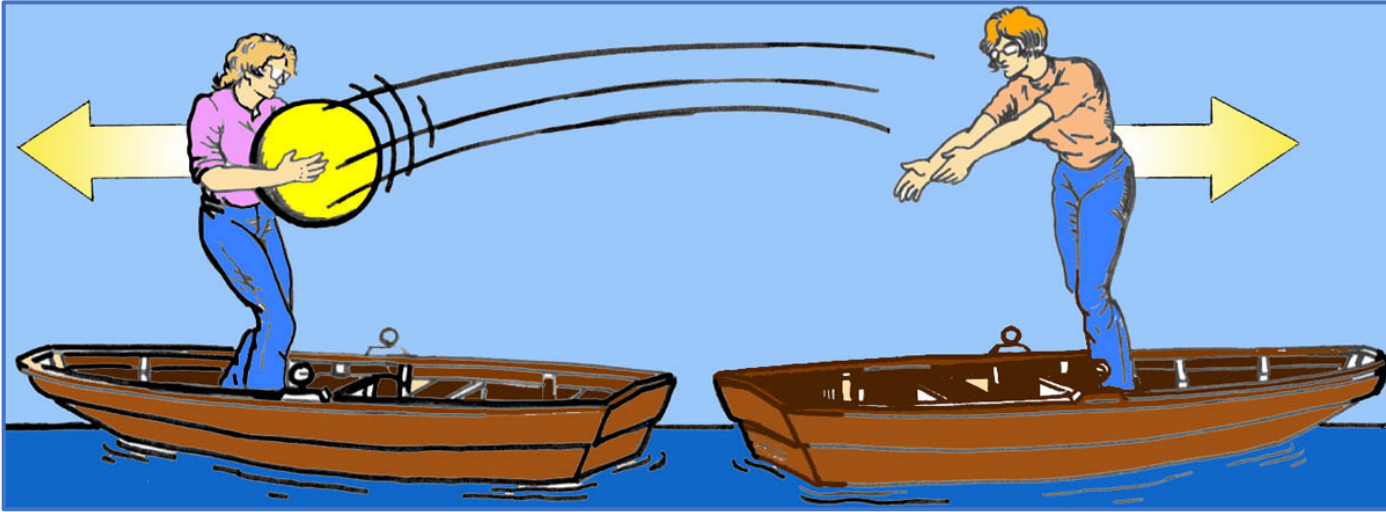


- Electromagnetism
- Weak Interaction
- Strong Interaction
- No Gravitation

Lecture 2 : Forces

[illegible]

Attractive and Repulsive forces and the quantum exchange



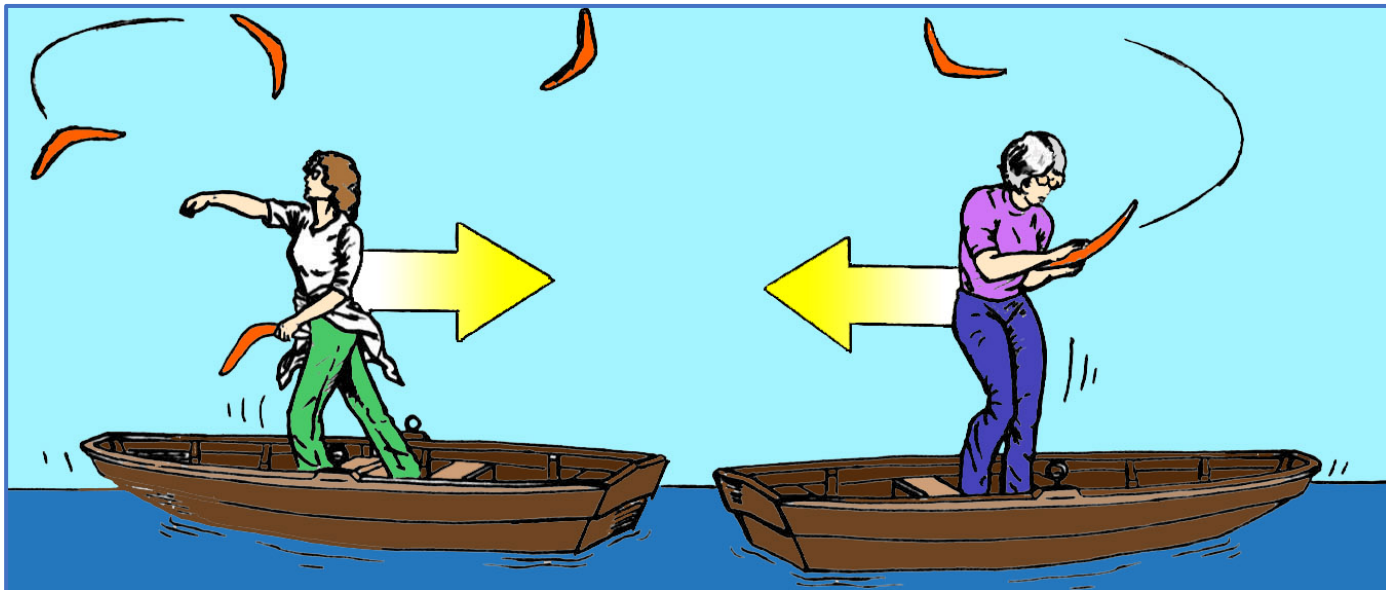
There is no
“action at a distance”

EM: photon

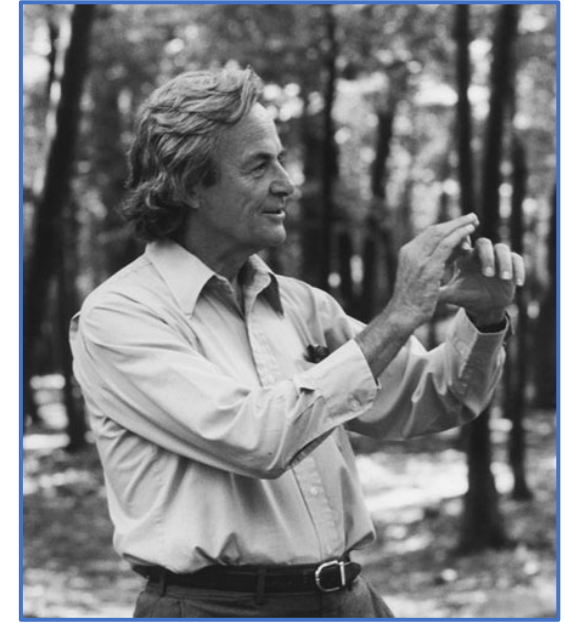
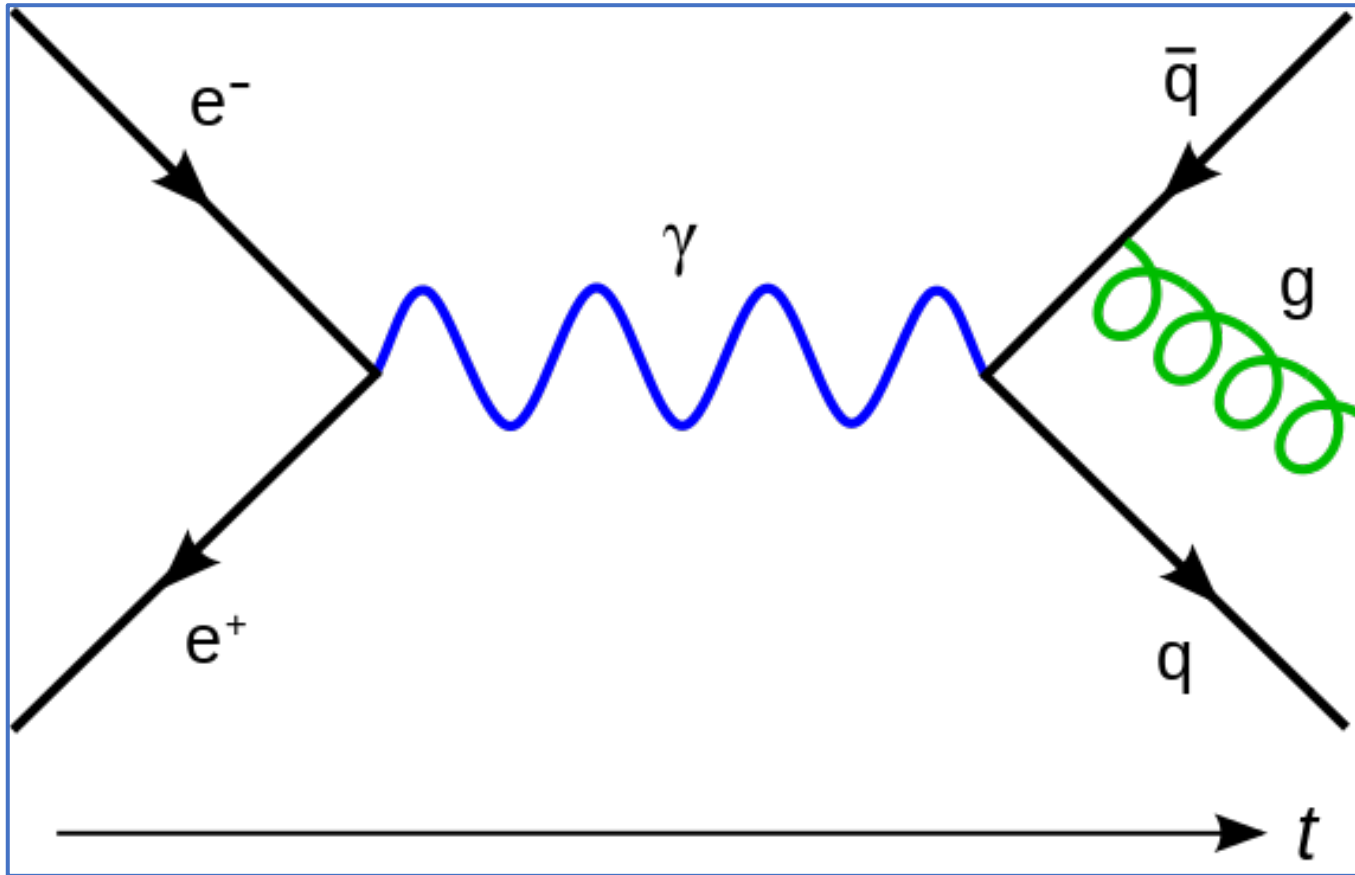
Weak: W, Z bosons

Strong: gluons

Gravitation: graviton(?)



Example of a quantum process



Feynman diagram

Note:

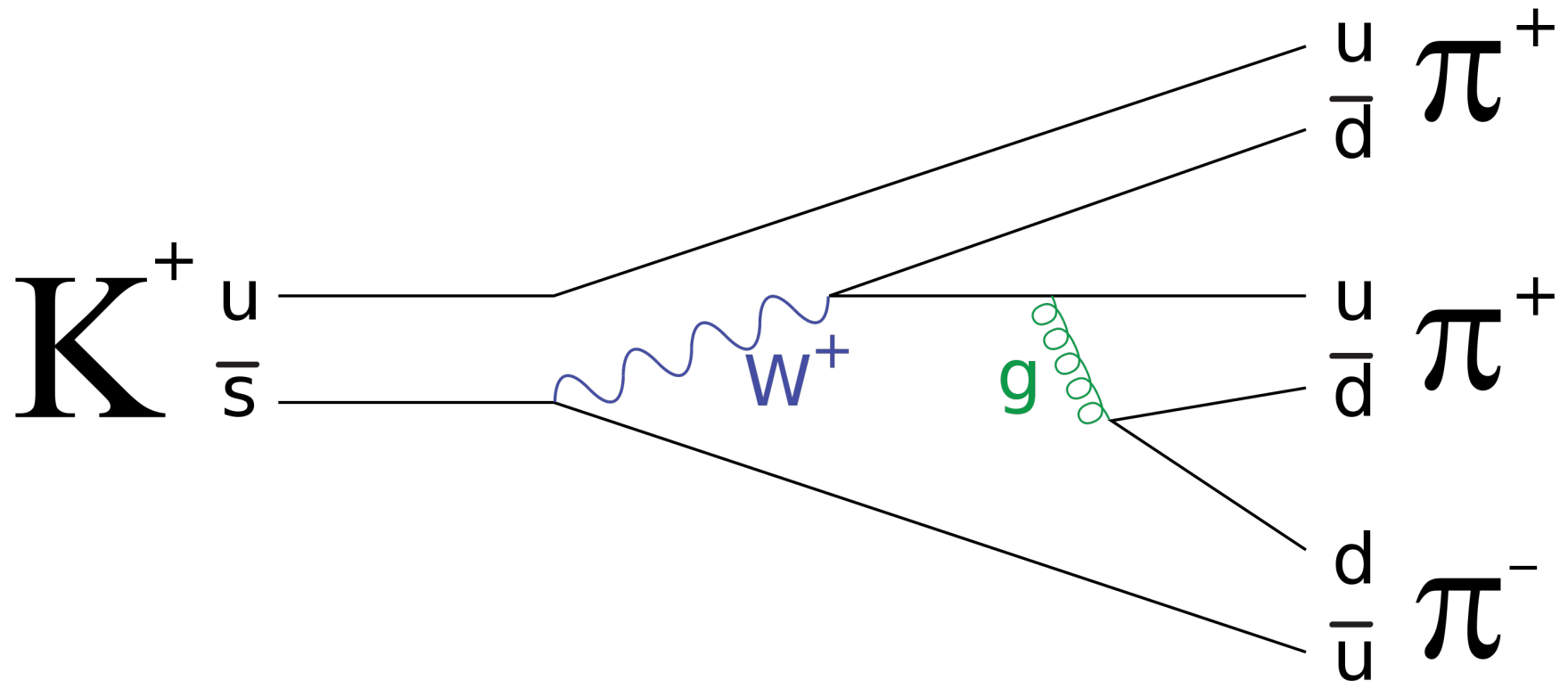
This is a **graphic** representation of a **calculation** that represents a quantum event. We use it to **talk** about it.

Do **not** take it **literally** as what happens

- Electron-positron annihilation
- Produce a “virtual” photon
- Photon produces a quark + antiquark
- Quark can radiate a gluon

Another example

- Kaon decay with *weak* interaction mediated by a W -boson
- Quark anti-quark produced by the *strong* interaction mediated by a gluon

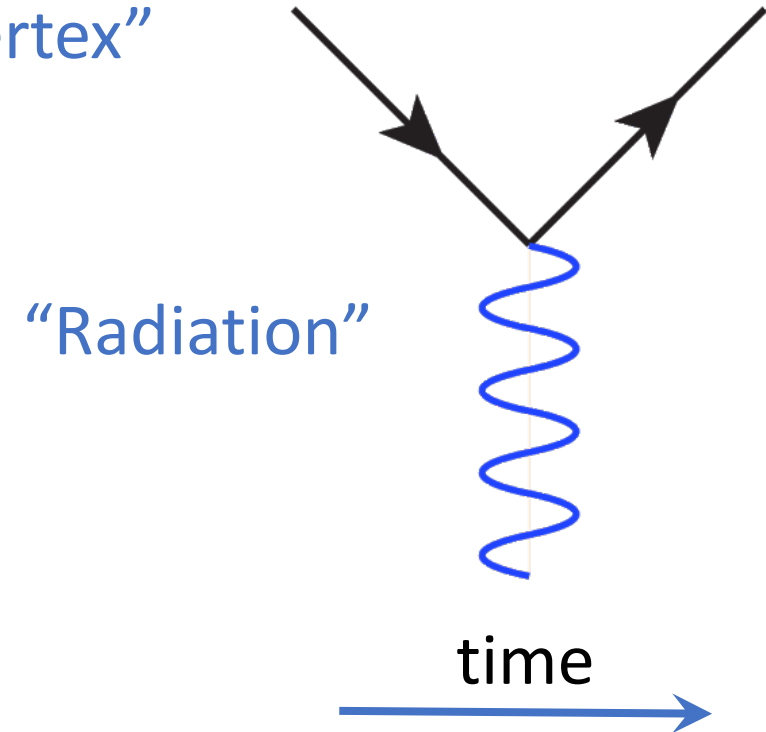
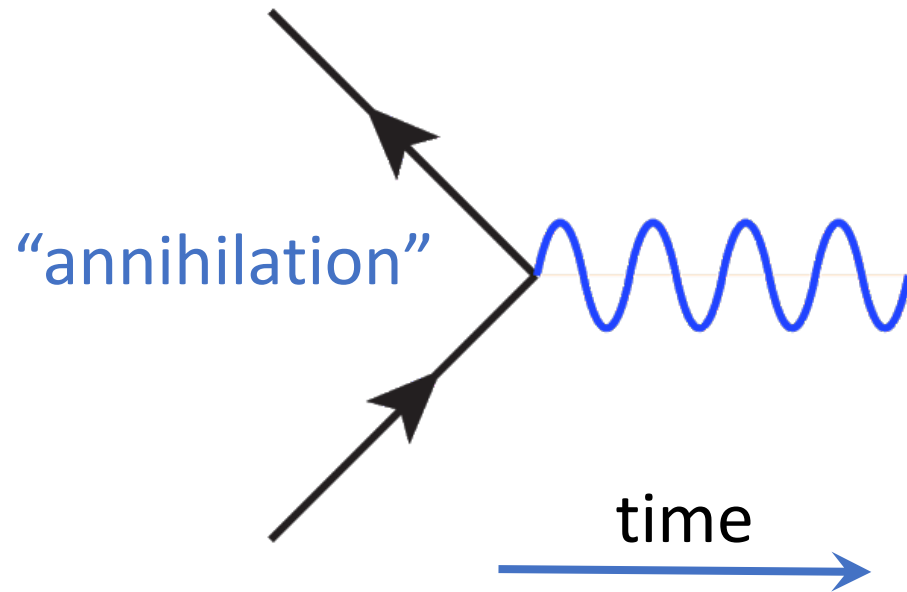


Part 1

The Electromagnetic Interaction Quantum Electrodynamics (QED)

Quantum Electrodynamics (QED)

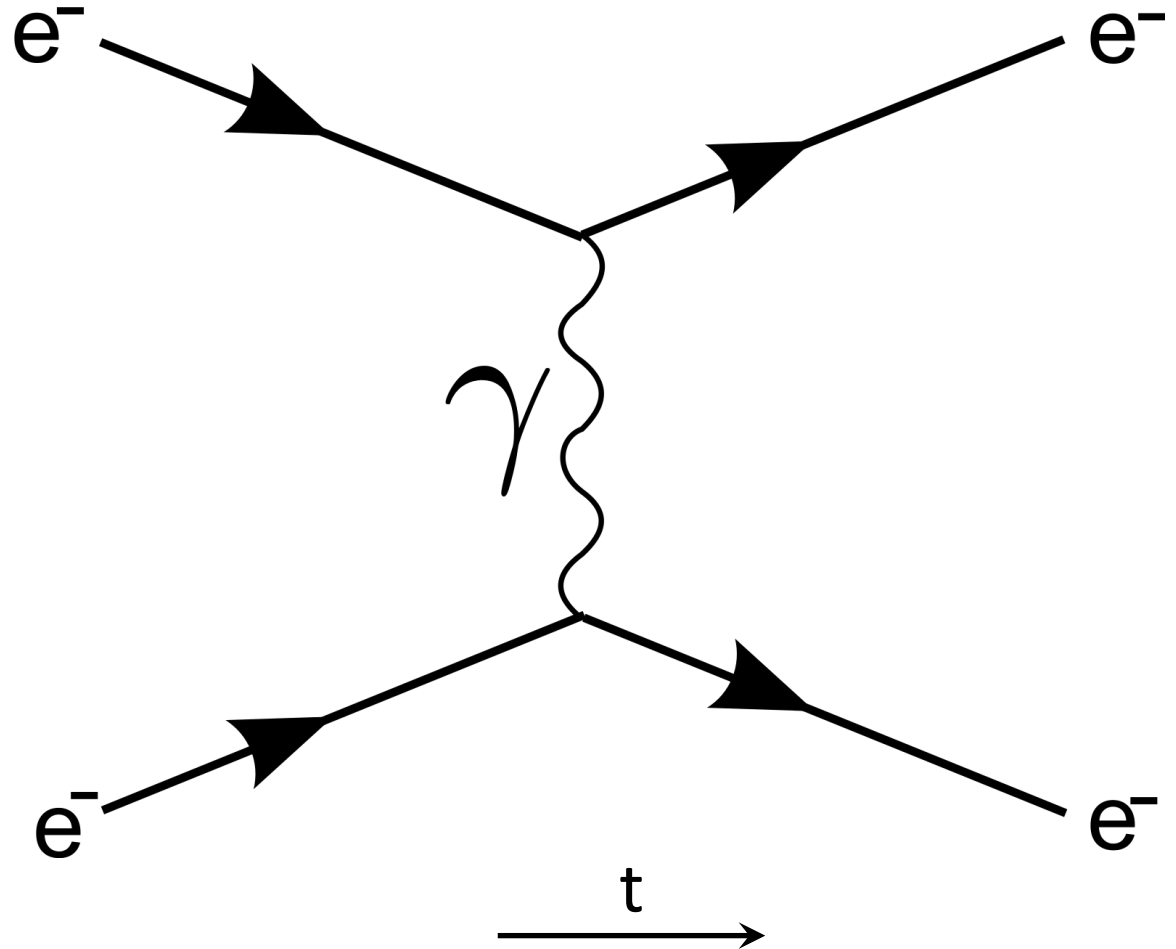
- There is only one fundamental quantum interaction process: ie one “vertex”



- “strength” of the “vertex coupling” is equal to the elementary charge e
- Note: this vertex can only be part of a process since energy and momentum cannot be conserved at the same time in a “2-to-1” process
 - This is an off-shell or virtual photon

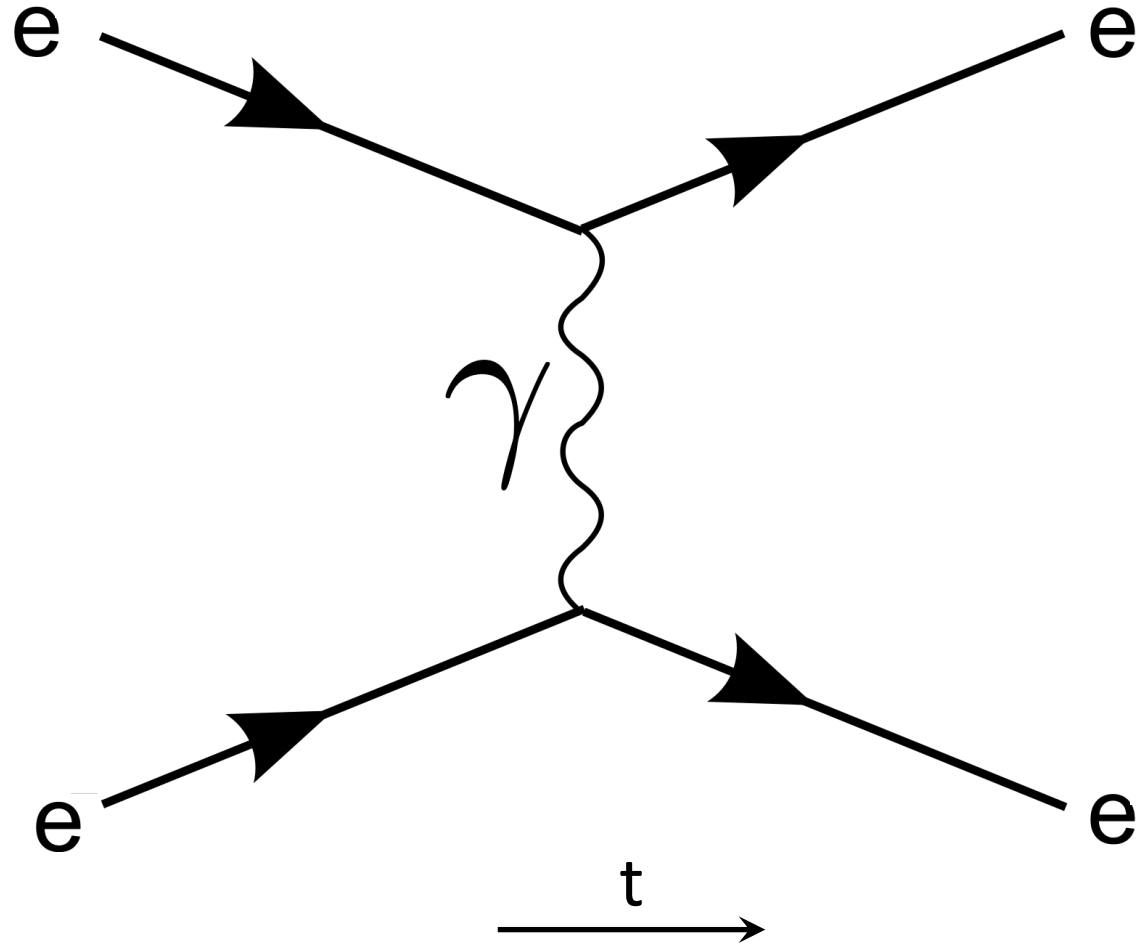
A Real QED Process: Möller scattering

- “Coulomb” repulsion of colliding electrons
- Convention: the direction of an arrow w.r.t. time determines whether it is a colliding particle or antiparticle
 - Should leave out the charge to be unambiguous



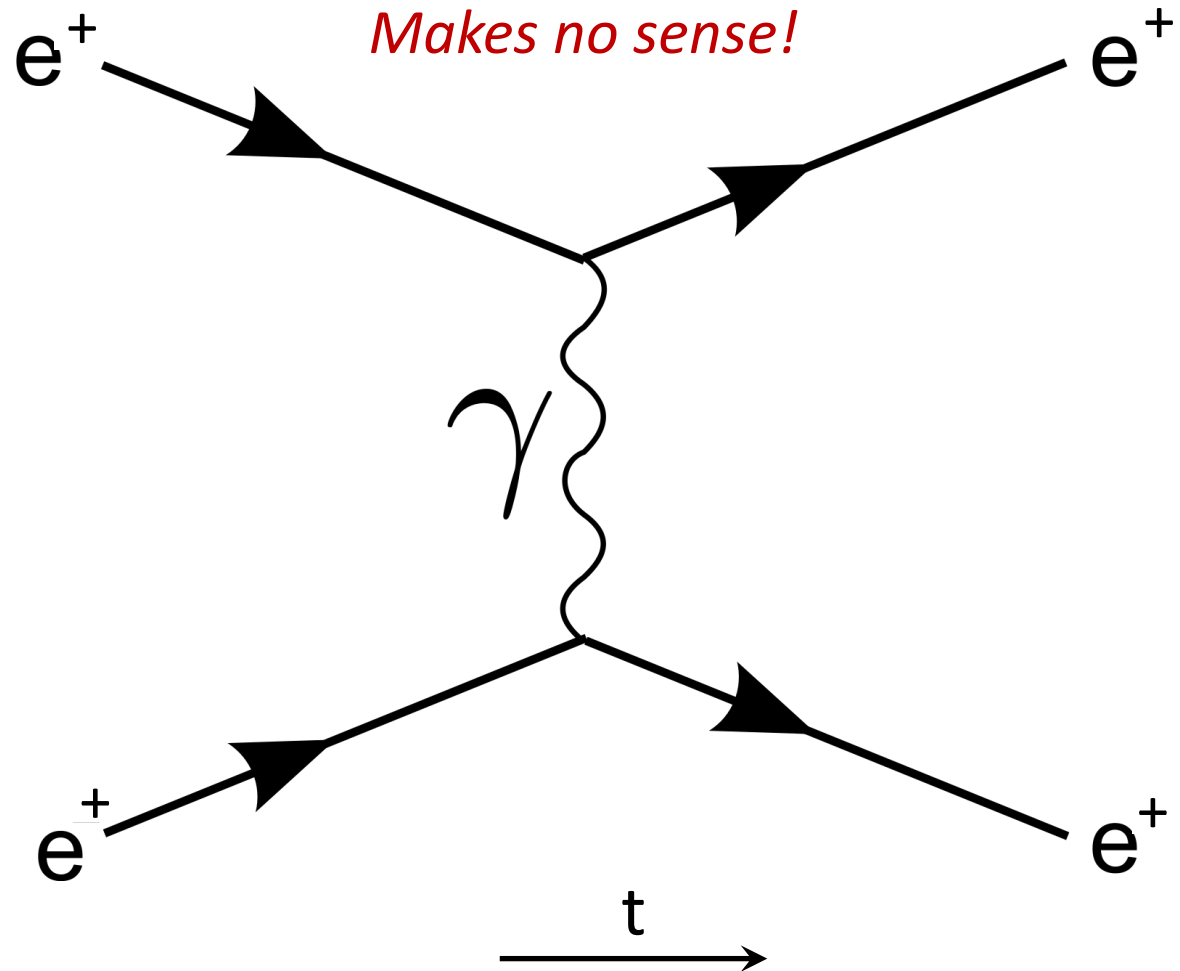
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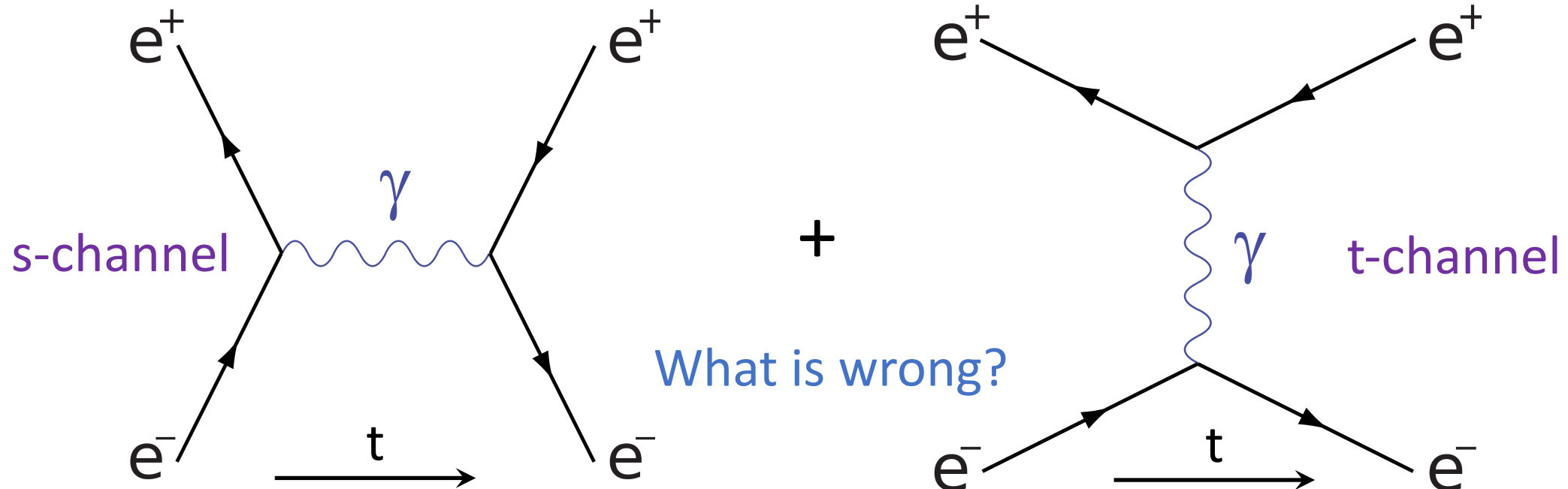
A Real QED Process: Möller scattering

- “Coulomb” repulsion of colliding electrons
- Convention: the direction of an arrow w.r.t. time determines whether it is a colliding particle or antiparticle
 - Should leave out the charge to be unambiguous
- Feynman diagrams are always drawn with *particles*
 - *Anti-particles* are represented with an arrow pointing against the direction of time.



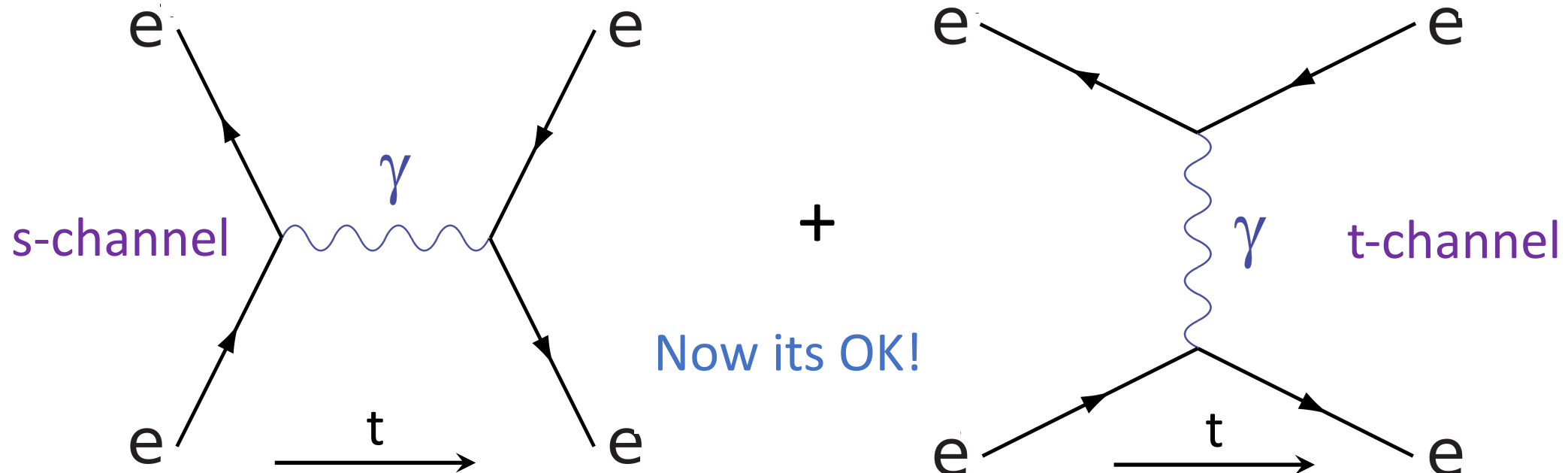
A Real Process: Bhabha Scattering

- Scattering of electron and positron, two quantum processes (“Feynman diagrams”) in one real process:
 - Spacelike “s-channel” exchange
 - e^- and e^+ annihilate into a photon, which converts back to a e^- and e^+ pair
 - Timelike “t-channel” exchange
 - e^- and e^+ scatter in each others EM field



A Real Process: Bhabha Scattering

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Feynman Diagram

- Calculation using Feynman rules:

$$\mathcal{M} = \frac{e^2}{q^2}$$

- Fine structure constant:

$$\alpha = \frac{e^2}{4\pi}$$

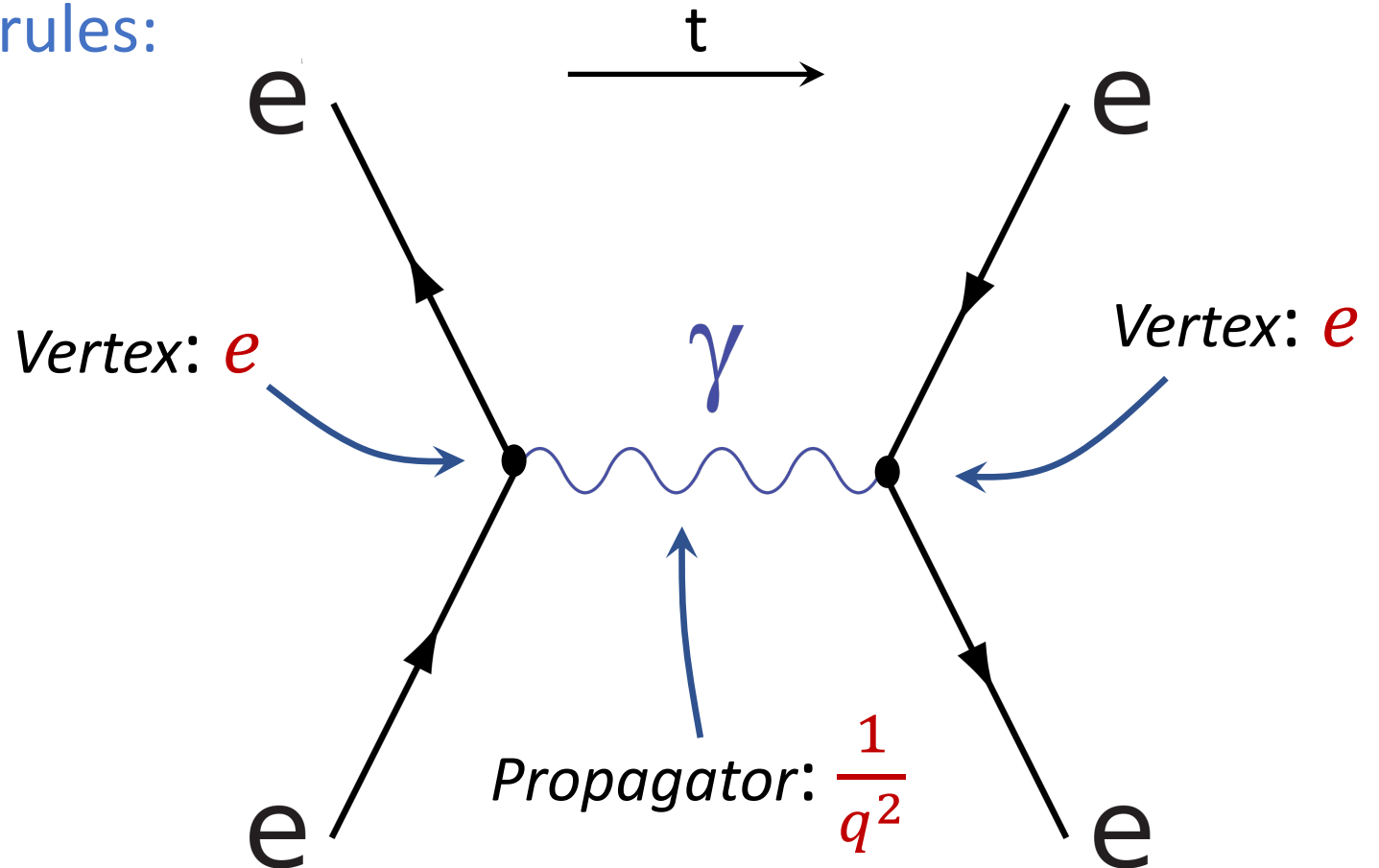
$$\text{Hence: } \mathcal{M} \propto \frac{\alpha}{q^2}$$

q^2 is “virtuality” of the photon

- Strength of the interaction

- Determined by the coupling constant e :

$$\alpha = \frac{e^2}{4\pi} = 1/137 \quad \text{why???}$$

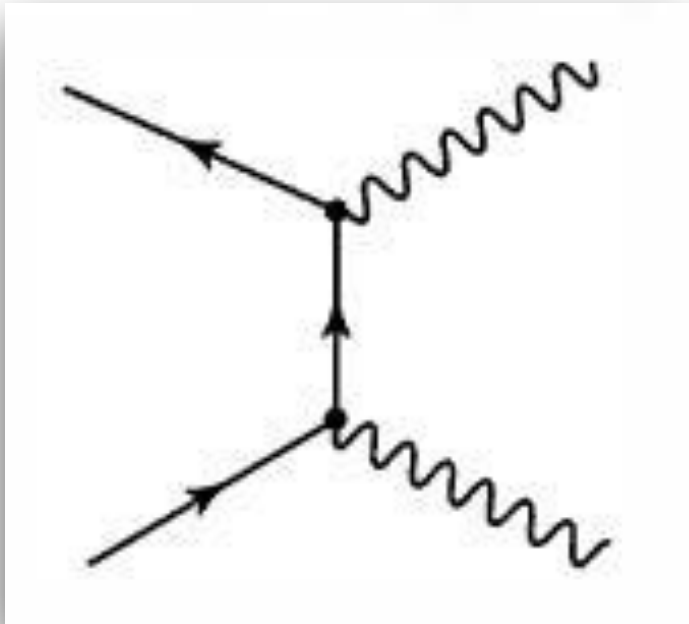


Note: a photon should have $q^2 = 0$!

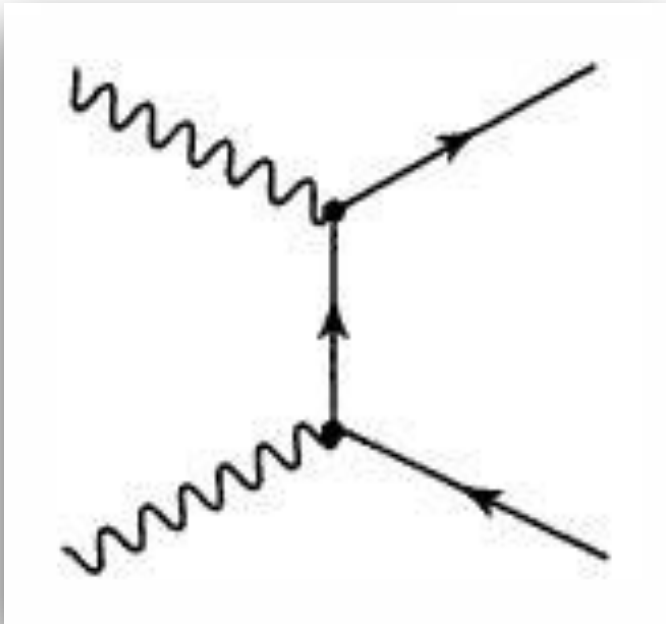
Scattering with photons

Real processes with external photon lines

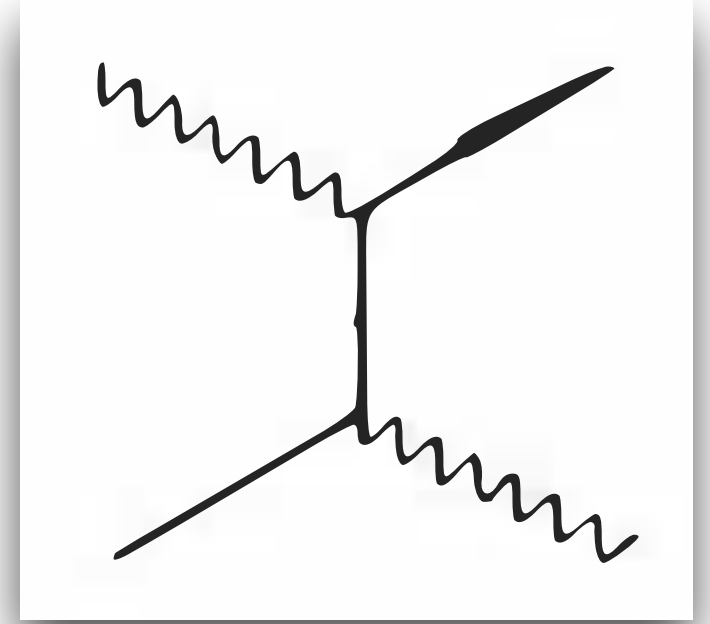
t \longrightarrow



Pair annihilation



Pair production

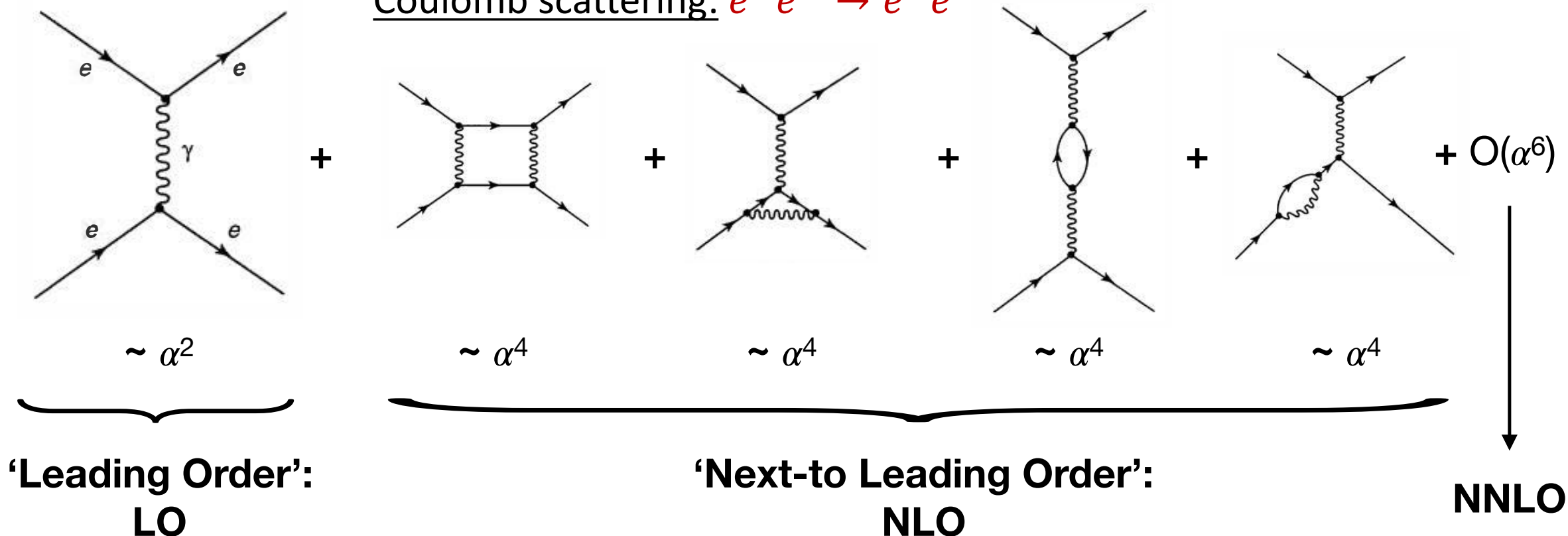


Compton scattering

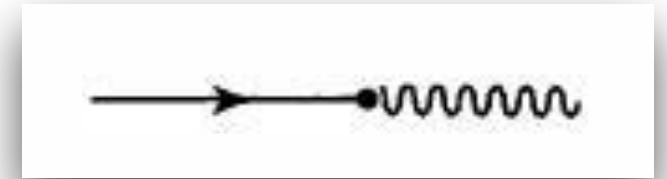
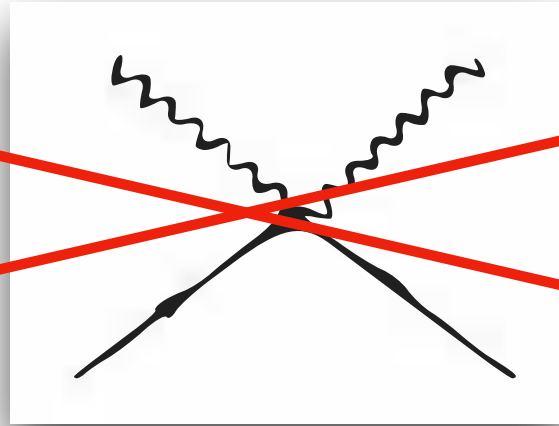
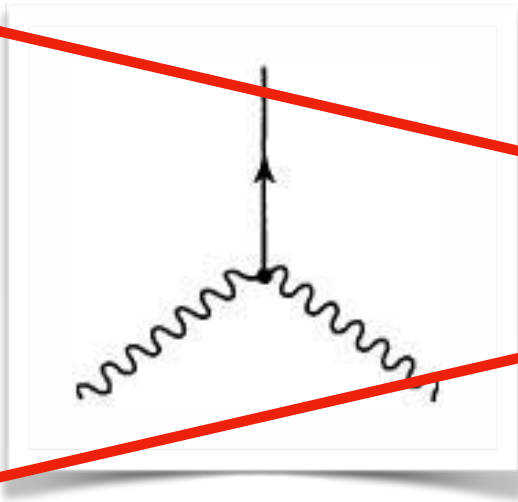
Feynman diagrams we use describe the lowest order “Perturbation Theory”. Calculations with Feynman diagrams are never exact, always an approximation.

- We simply do not know how to do exact calculations in particle physics
- QED is extremely precise because many complicated diagrams have been calculated

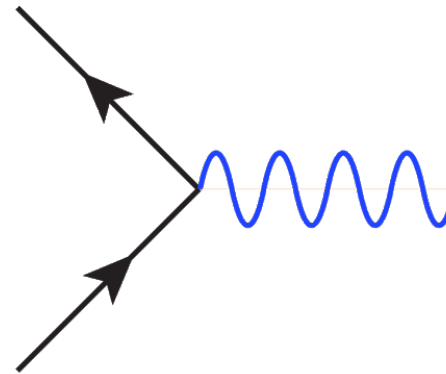
Coulomb scattering: $e^-e^- \rightarrow e^-e^-$



Not all diagrams exist



- You can only use combinations of the fundamental QED vertex:



The background of Feynman Diagrams

- More about scattering and Feynman diagrams in Lecture 4
- For an intuitive approach to Feynman diagrams, see wikipedia: https://en.wikipedia.org/wiki/Quantum_electrodynamics and references therein, or the booklet of Feynman:
 - “The strange theory of light and matter”

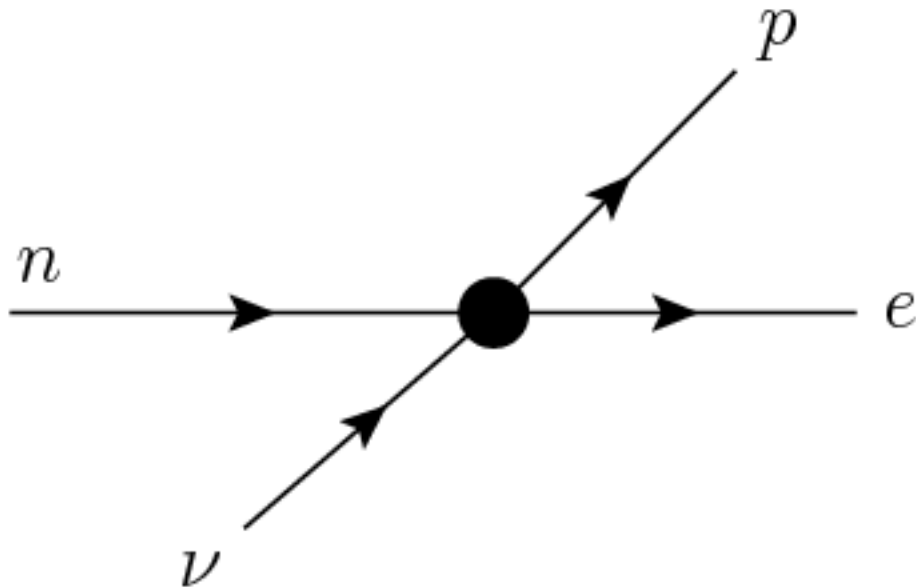
Part 2

The Weak Interaction

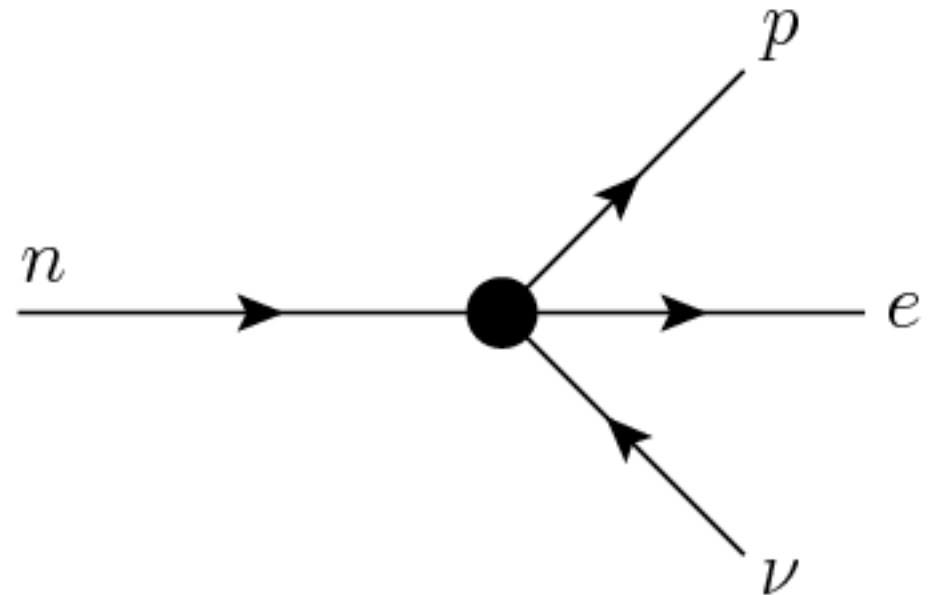
Historical: Fermi 4-point interaction

- Becquerel discovered radioactivity base on the weak force
 - $n \rightarrow p + e + \bar{\nu}$
- The original model for the weak interaction is from Fermi
 - For obvious reasons it is called the 4-point 'contact' interaction (1933).
 - The strength of the interaction was given by Fermi's constant $G_F = 10^{-5}$

$$n + \nu \rightarrow p + e$$

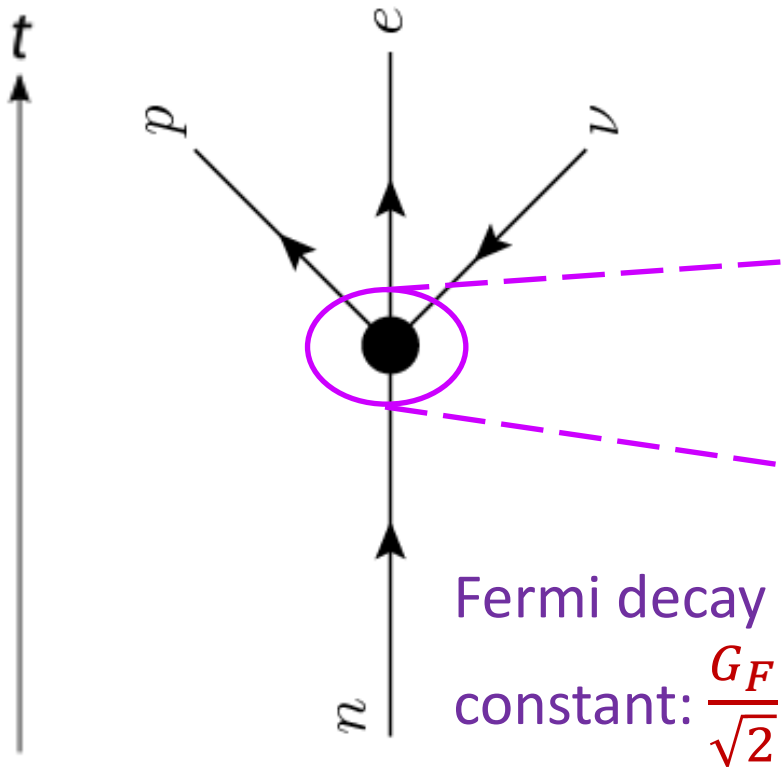


$$n \rightarrow p + e + \bar{\nu}$$

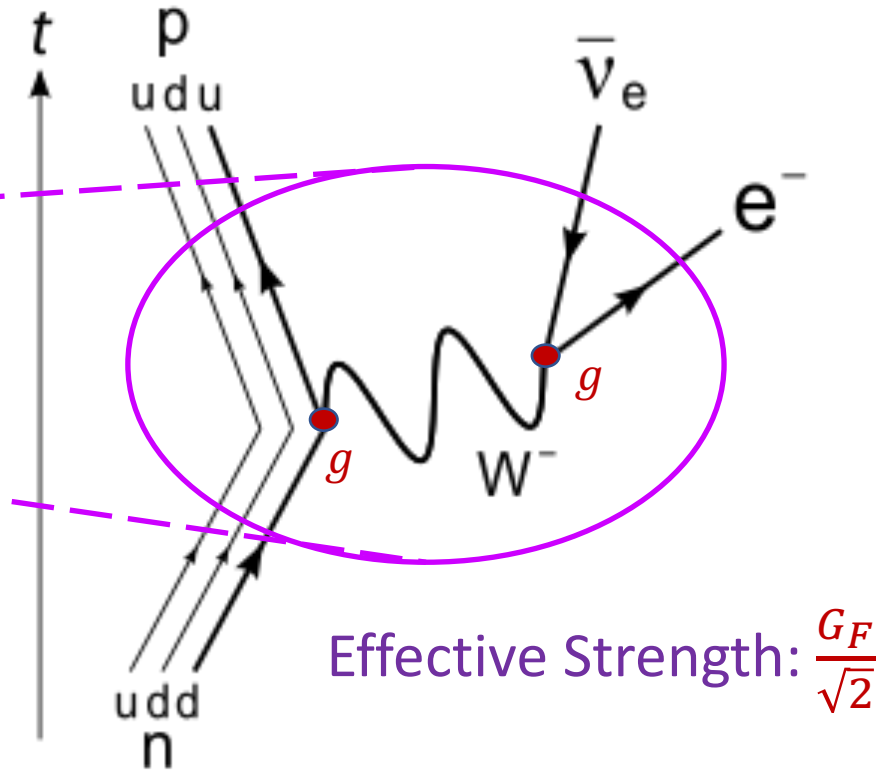


Quantum zoom-in on Fermi 4-point interaction

- Fermi 4-point interaction



- Very short range W exchange

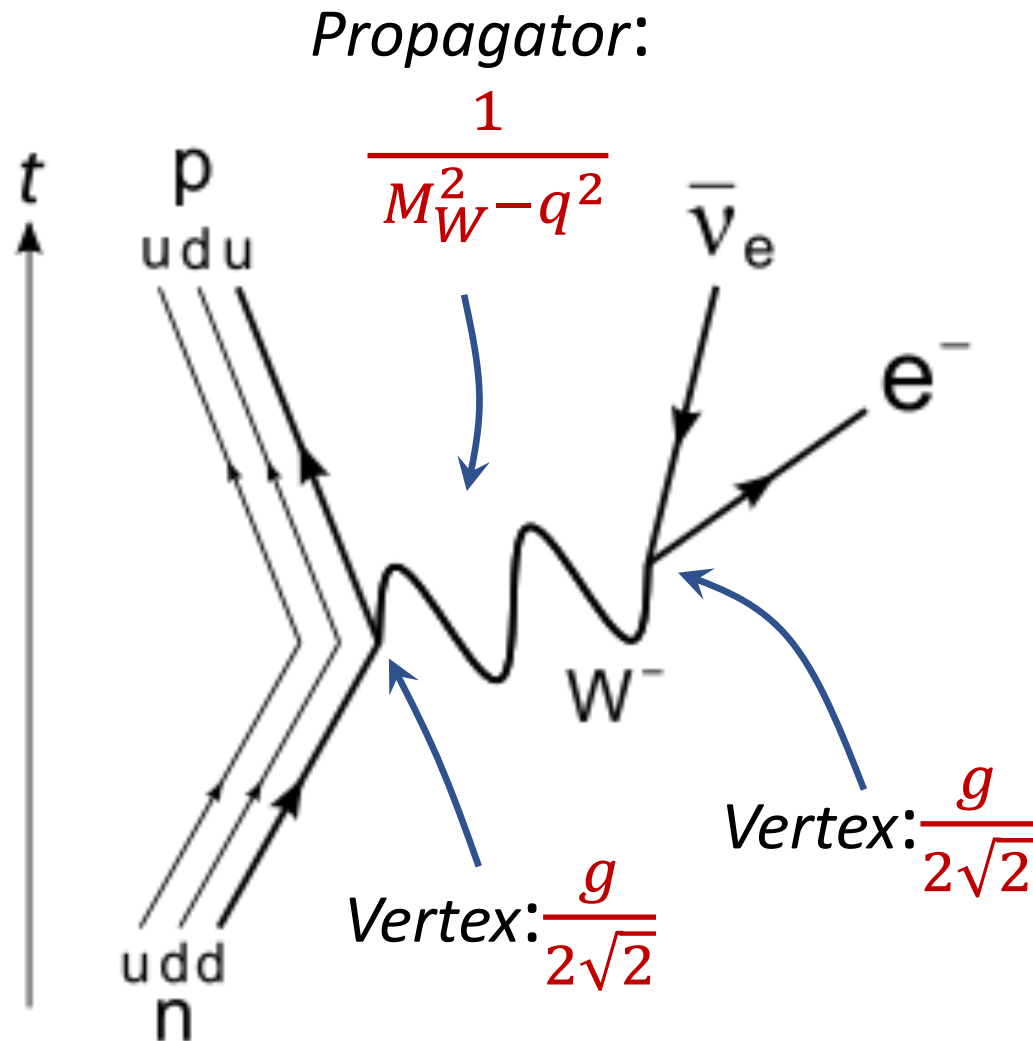


$G_F = \text{Fermi Constant}$

- Similar to electrodynamics the weak interaction has a propagator
 - W spin-1 boson
- At quark level it involves a $d \rightarrow u$ transition

The Weak Interaction

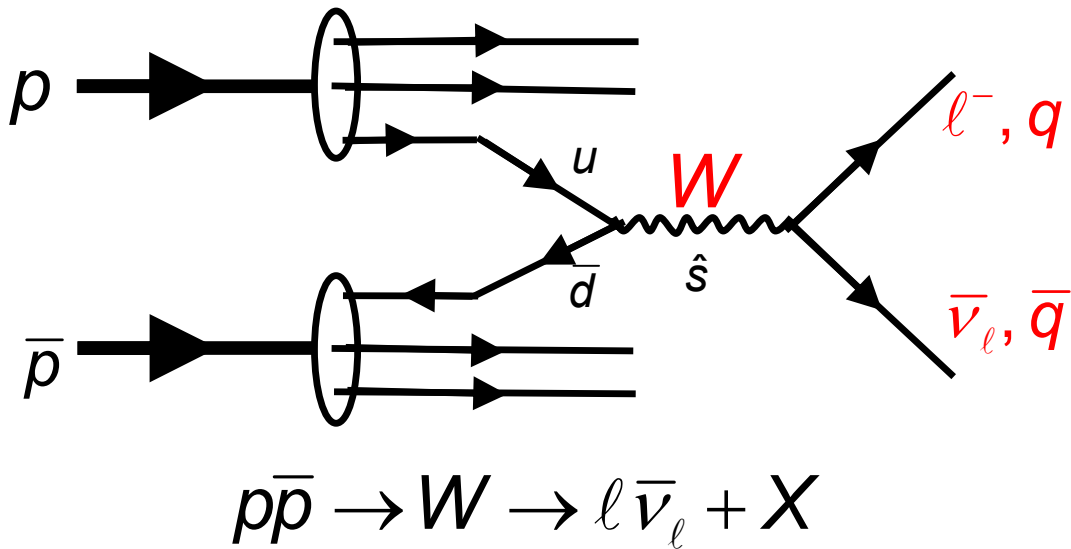
- Feynman rules



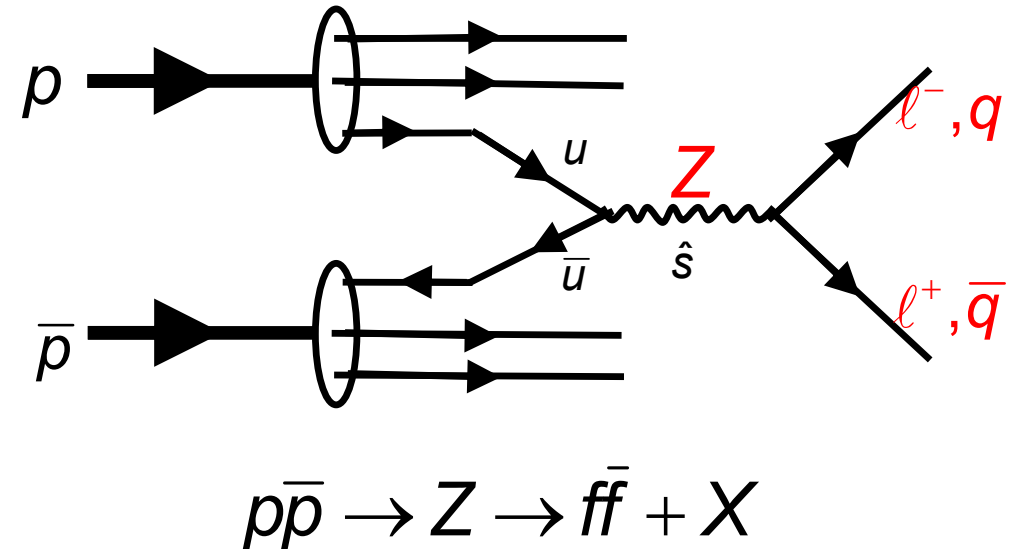
- W -boson has electric charge
 - The process is called a *charged current*
- q^2 is the momentum transfer from the hadronic to leptonic vertex
 - Here due to mass difference $m_n - m_p$
- Effect of propagator is only noticeable when $q^2 \sim M_W^2$
 - If q^2 is large: **resonance**
 - In that case weak force is strong
 - If q^2 is small: then effectively a 4-point Fermi coupling
 - Propagator $\sim 1/M_W^2 \rightarrow$ weak force
 - $G_F = \frac{g^2}{8M_W^2}$; $\alpha_w = \frac{g^2}{2\pi} = \frac{1}{30} > \alpha_{QED}$
- The weak coupling constant is large!
 - Interaction is weak due to large W mass

W and Z bosons

- Glashow, Salam, Weinberg (GSW model, 1968) predict three weak mediators
 - W^+ , W^- , Z - bosons
- Discovery of the W and Z bosons at CERN, 1983, UA1 and UA2 experiments
 - Look at high energy proton-proton collisions
 - $M_W = 80 \text{ GeV}/c^2$, $M_Z = 91 \text{ GeV}/c^2$ as predicted by GSW model



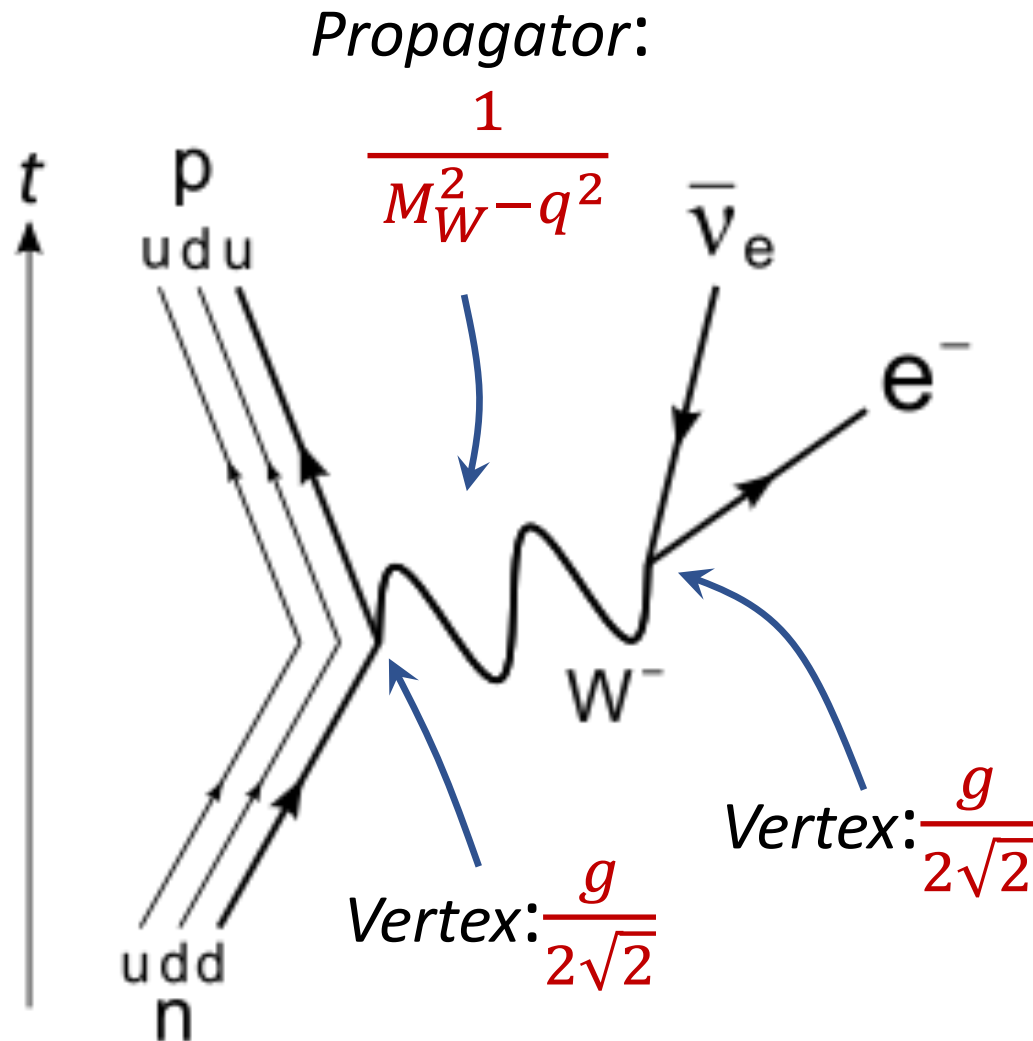
Flavour changing charged currents: $u\bar{d}$



Flavour conserving neutral currents: $u\bar{u}$

Virtuality

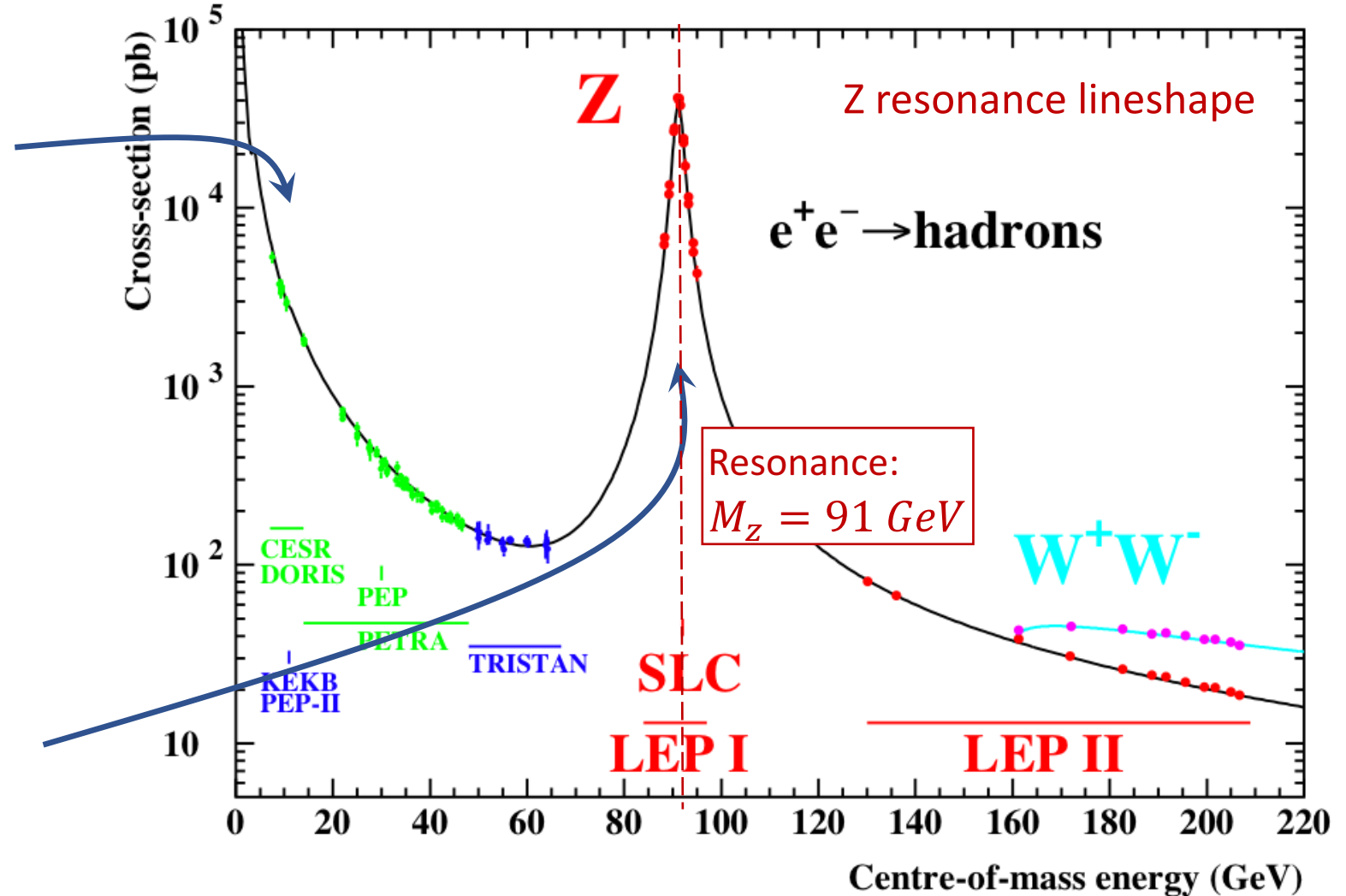
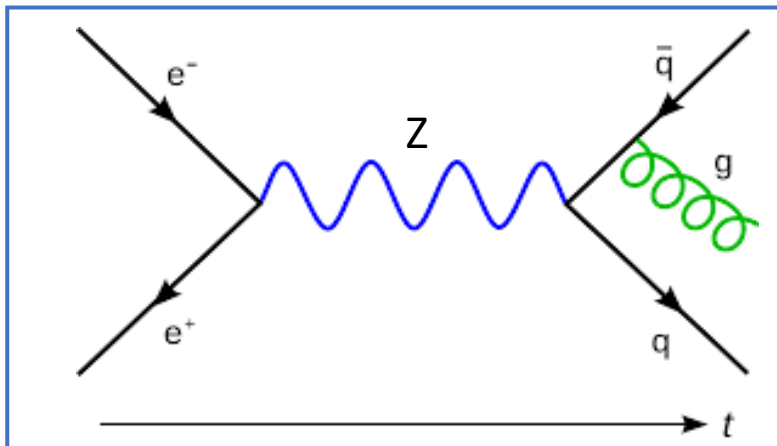
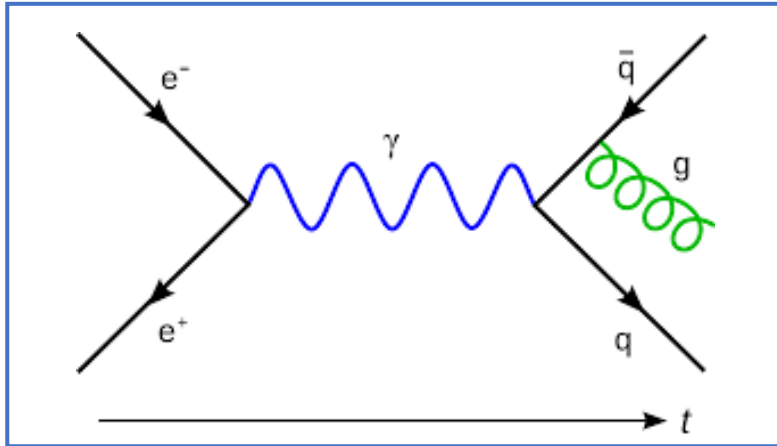
- How can beta decay $n \rightarrow p + e + \bar{\nu}$ work, if the W has a mass of 80 GeV??



- Exchange of “virtual” particles, also called “off-shell”: $E^2 - \vec{p}^2 \neq m^2$
- Heisenberg:
 E is undetermined as $\Delta E \Delta t \geq \hbar/2$ with $\Delta E = \sqrt{\vec{q}^2 + m^2}$
 For small $q \rightarrow$ range force R
 $R \sim c \Delta t \approx \frac{\hbar}{2mc} \sim 1.2 \times 10^{-18} \text{ m}$
- Notice the resonance behaviour for $q^2 \approx M_W^2$

Precision studies with Z bosons: “Electroweak Force”

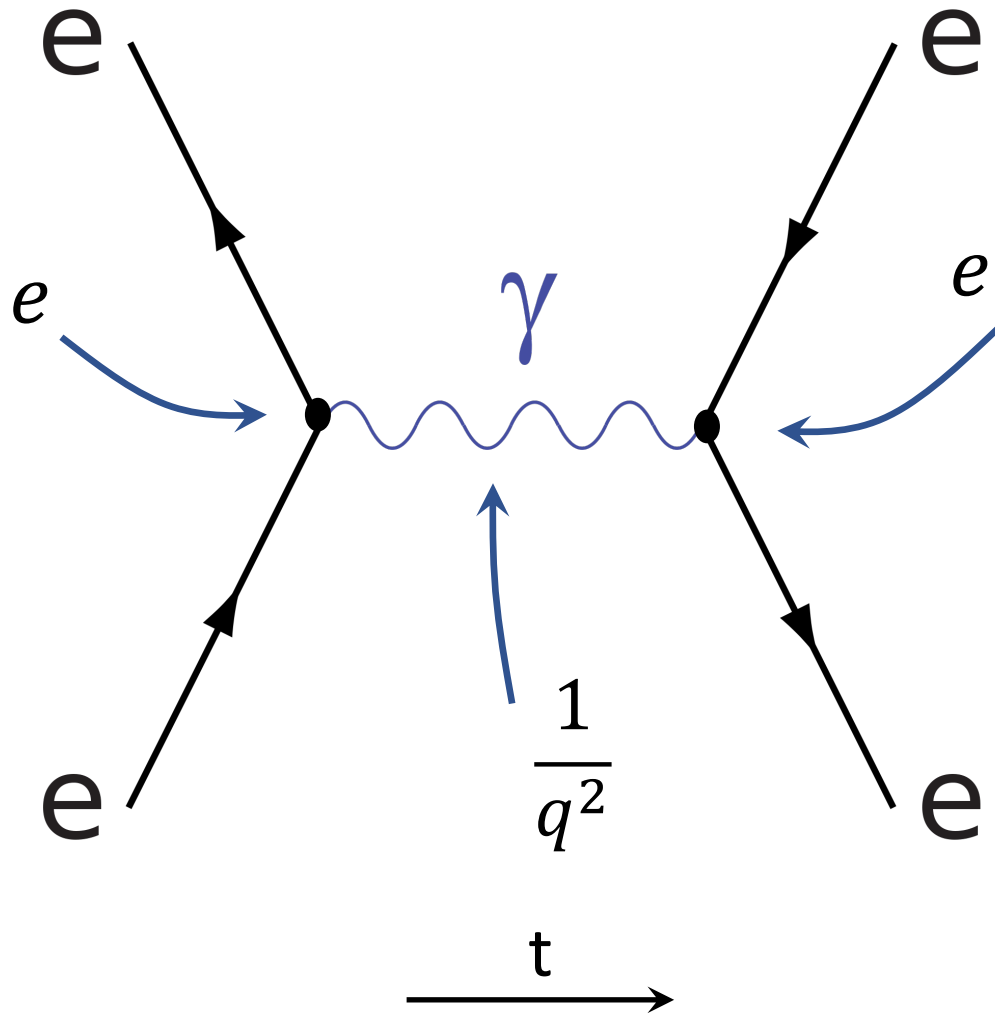
- The LEP collider did many precision studies of the weak interaction with e^+e^- annihilation collisions at collision energy around $91 \text{ GeV}/c^2$.



Similarity of Electromagnetic (γ) and Weak (Z) force

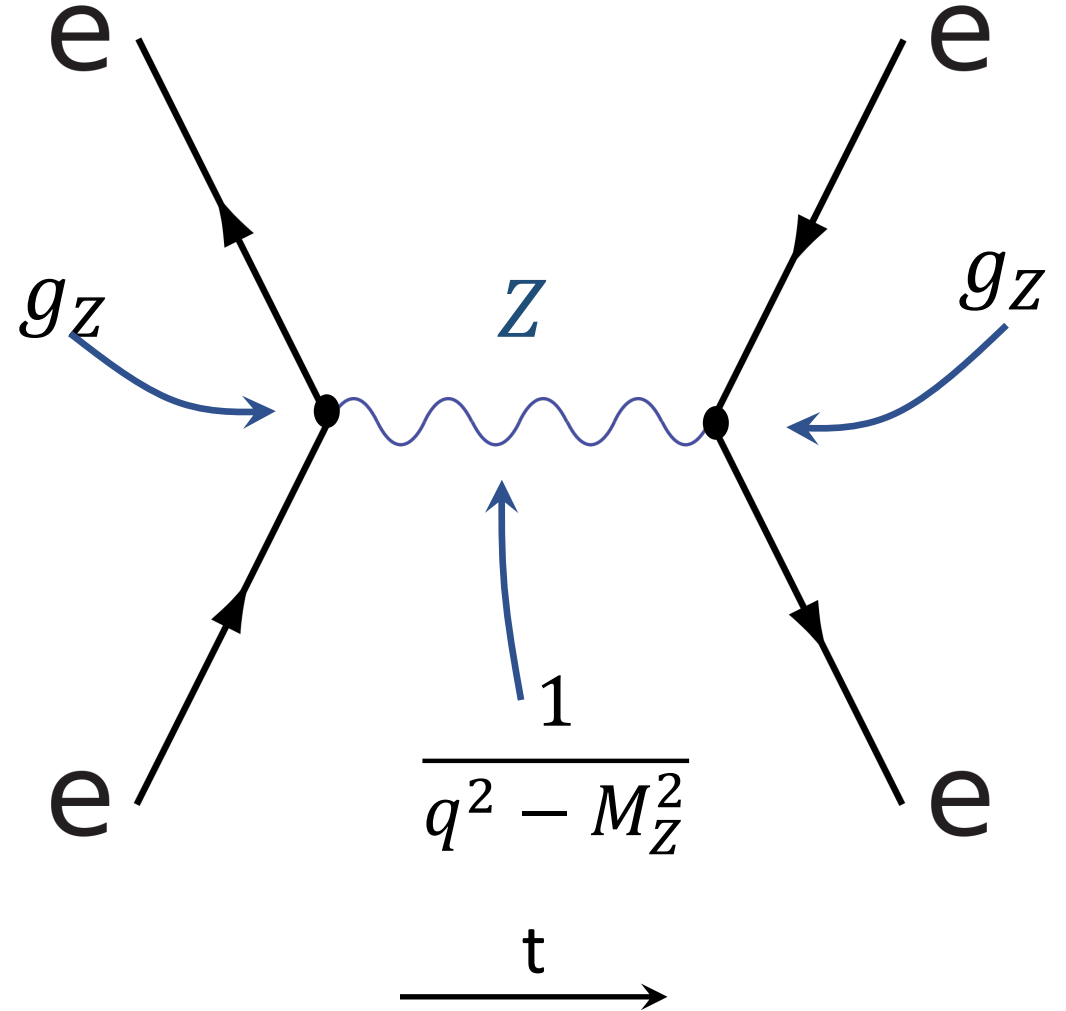
- Electromagnetic:

$$e^+e^- \rightarrow e^+e^-$$



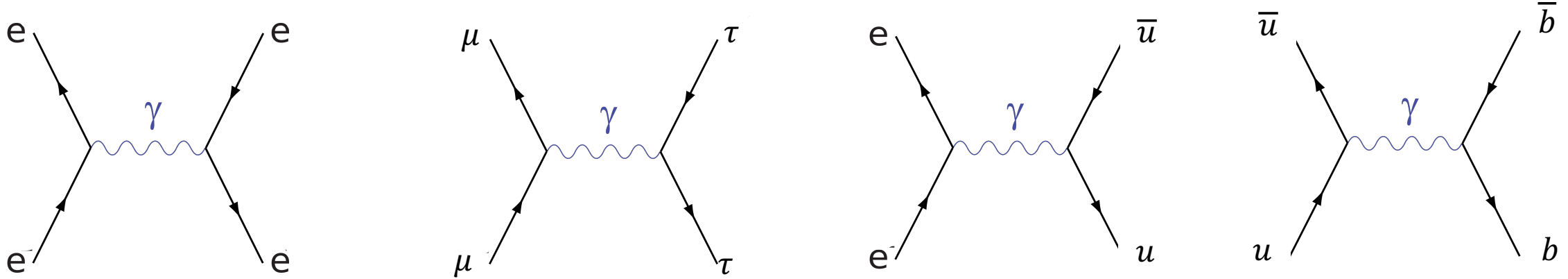
- Weak neutral current:

$$e^+e^- \rightarrow e^+e^-$$

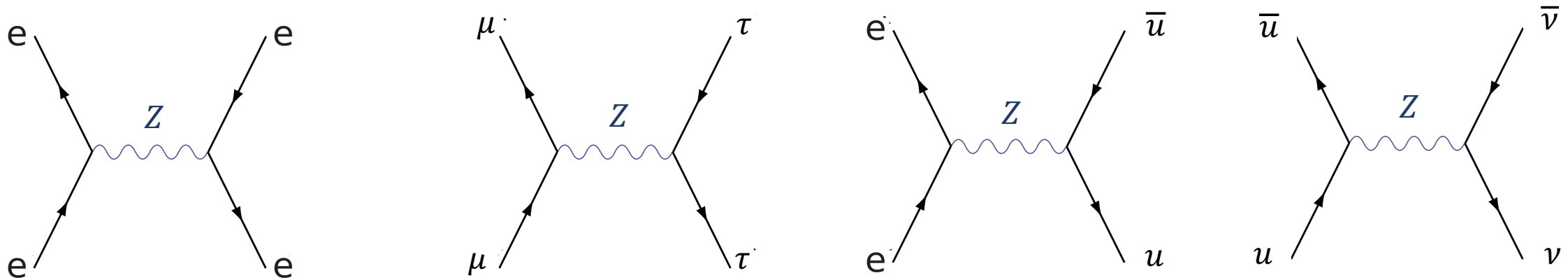


Examples of processes: “neural currents”

- Electromagnetic: couples to electric charge



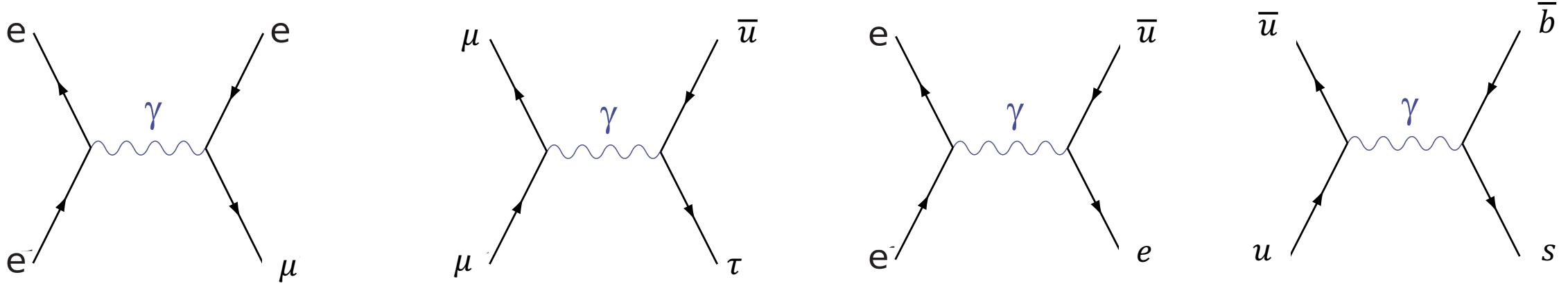
- Weak “Neutral Currents”: couples to weak charge



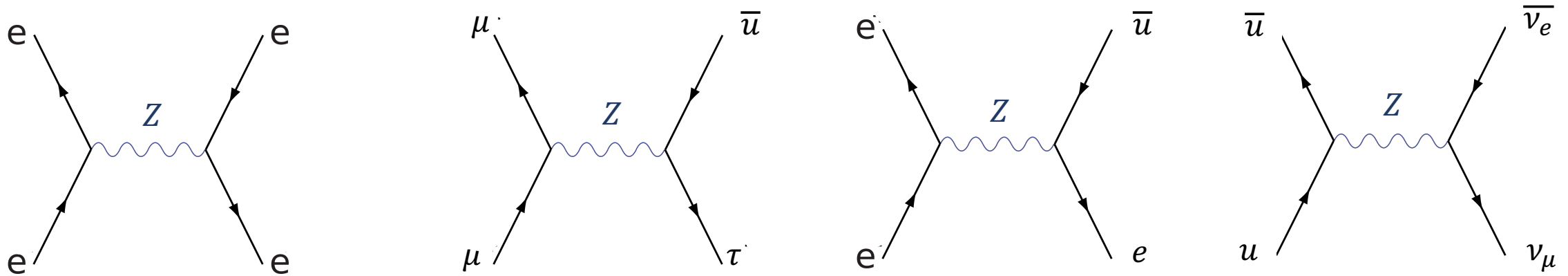
Z boson also couples to neutrinos: neutrinos have weak charge

Examples of processes that *do not exist*

- Electromagnetic:



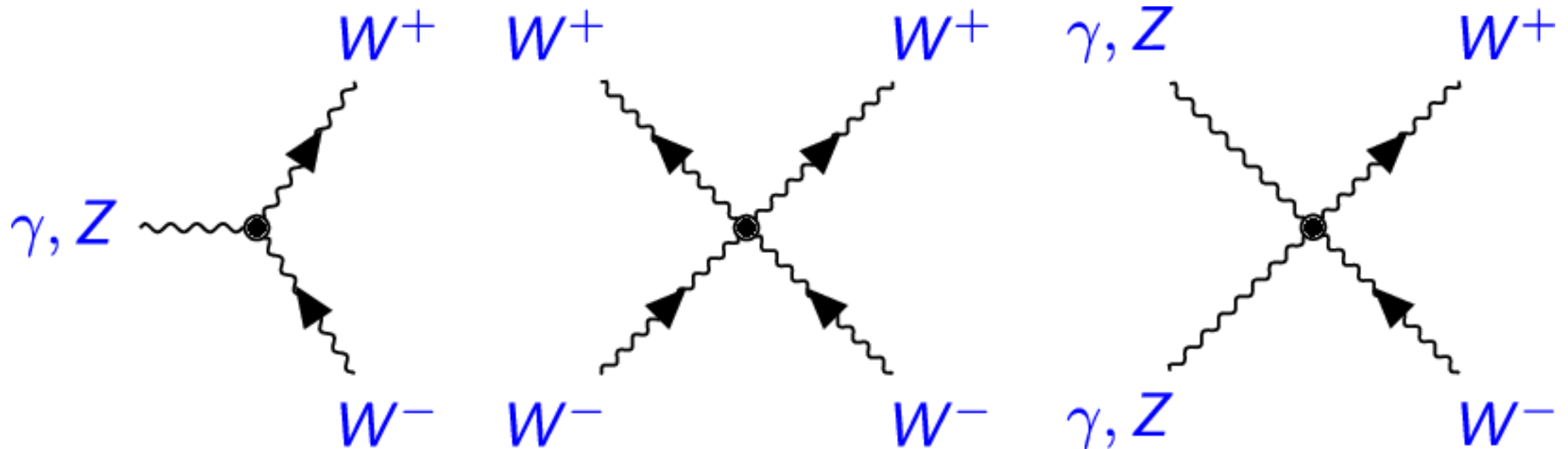
- Weak “Neutral Currents”:



γ and Z only couple to particles *within one generation* and *only to quark or lepton pairs*

Self coupling

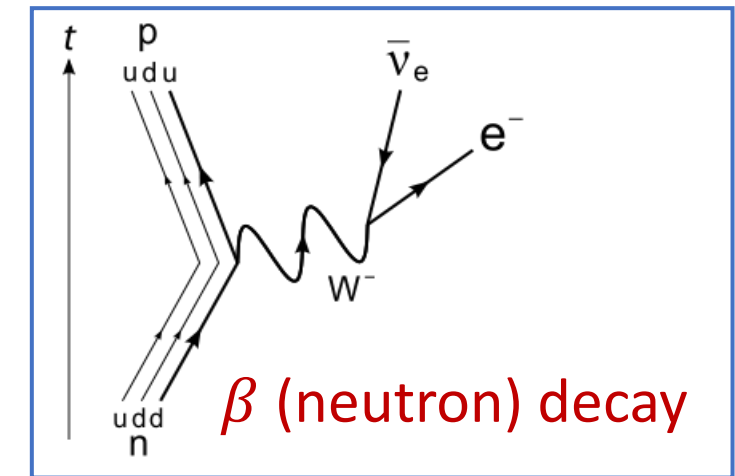
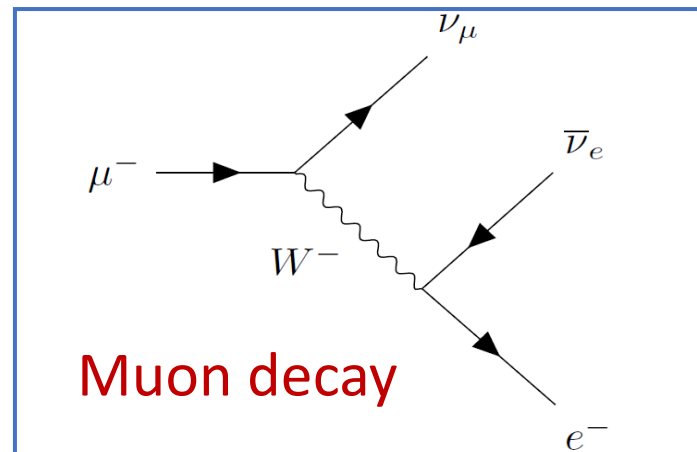
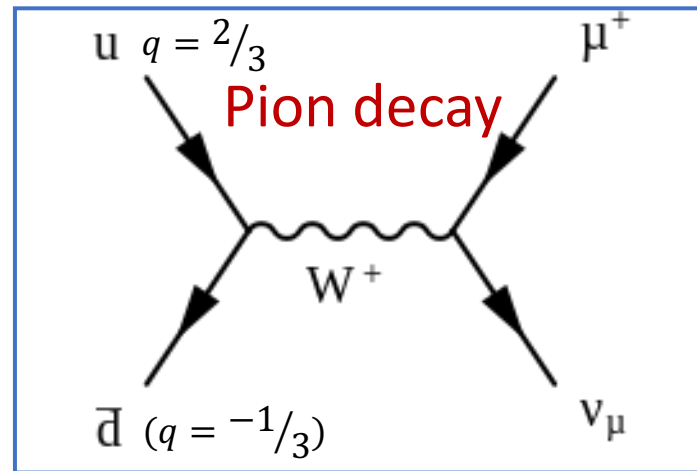
- In electromagnetism a photon has no electric charge
- In the weak interaction the force carriers have charge themselves
 - The weak charge is called “weak isospin”
 - The following diagrams are possible



The W weak force: “charged currents”

- The W -boson carries electric charge
- The W connects the weak isospin (I_3) ‘up’ and ‘down’ type particles of a generation
 - W^+ and W^- do the opposite
- 2nd and 3rd generation muon and tau can decay because neutrino’s are light!

		I	II	III	I_3
Quarks	+	u	c	t	$+\frac{1}{2}$
	-	d	s	b	$-\frac{1}{2}$
Leptons	+	ν_e	ν_μ	ν_τ	$+\frac{1}{2}$
	-	e	μ	τ	$-\frac{1}{2}$
Three Generations of Matter					



- How does a kaon decay?
 - $K^- = s\bar{d}$
- The c -quark is heavier than the s quark:
 - $s \rightarrow c e \bar{\nu}_e$ is not possible

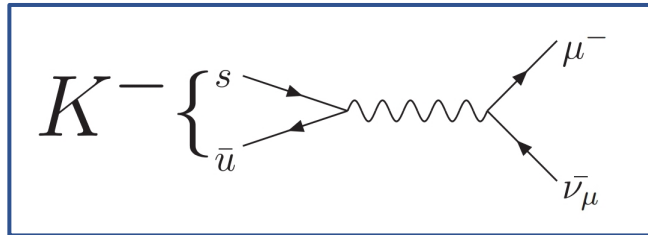
Generations

- Does the W **only** couple between particles of one generation?

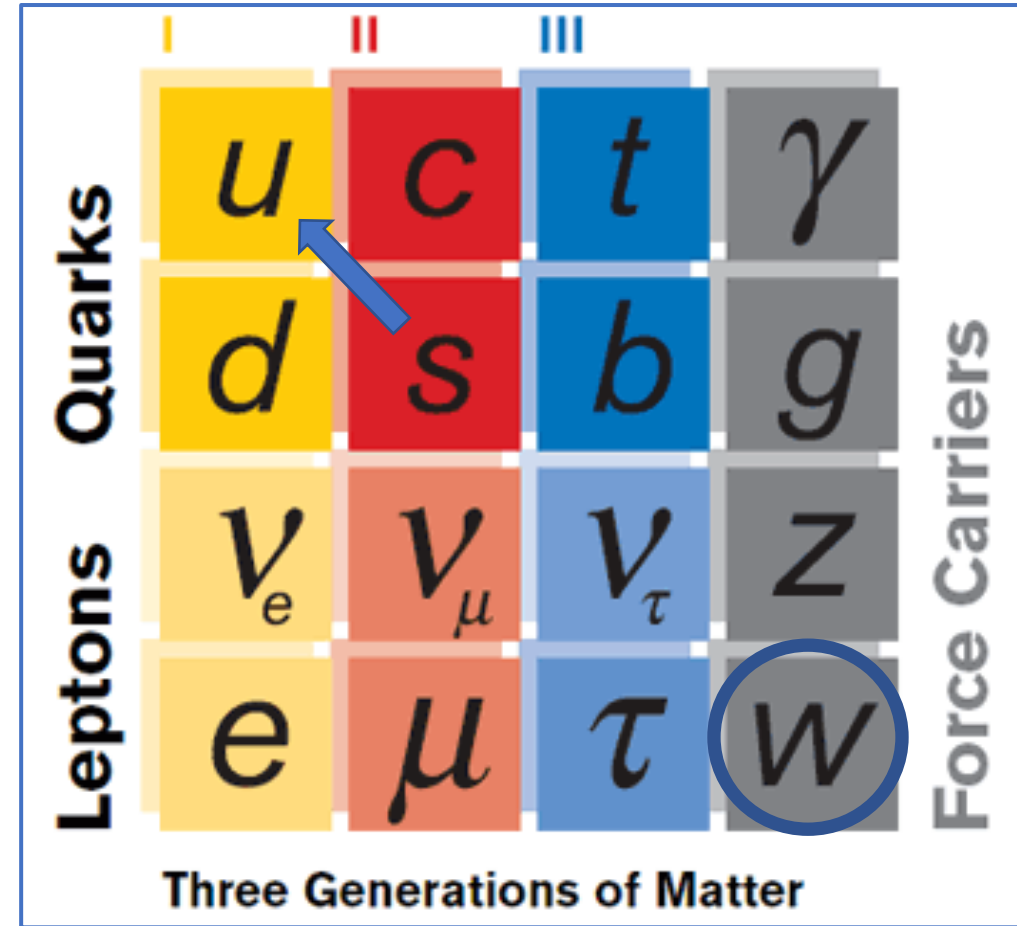
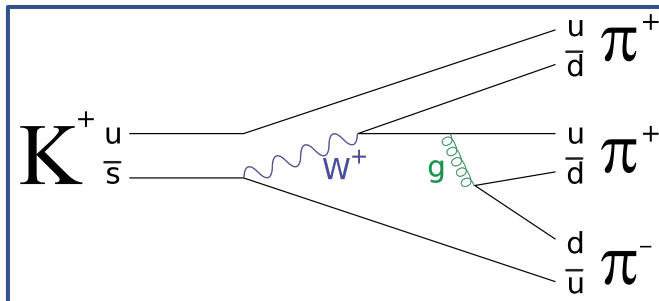
- Leptons: **yes**
- Quarks: **no!**
 - Couplings between generations are possible for the W !

- Kaon decays

- "Leptonic": $K^+ \rightarrow \mu^+ \nu_\mu$

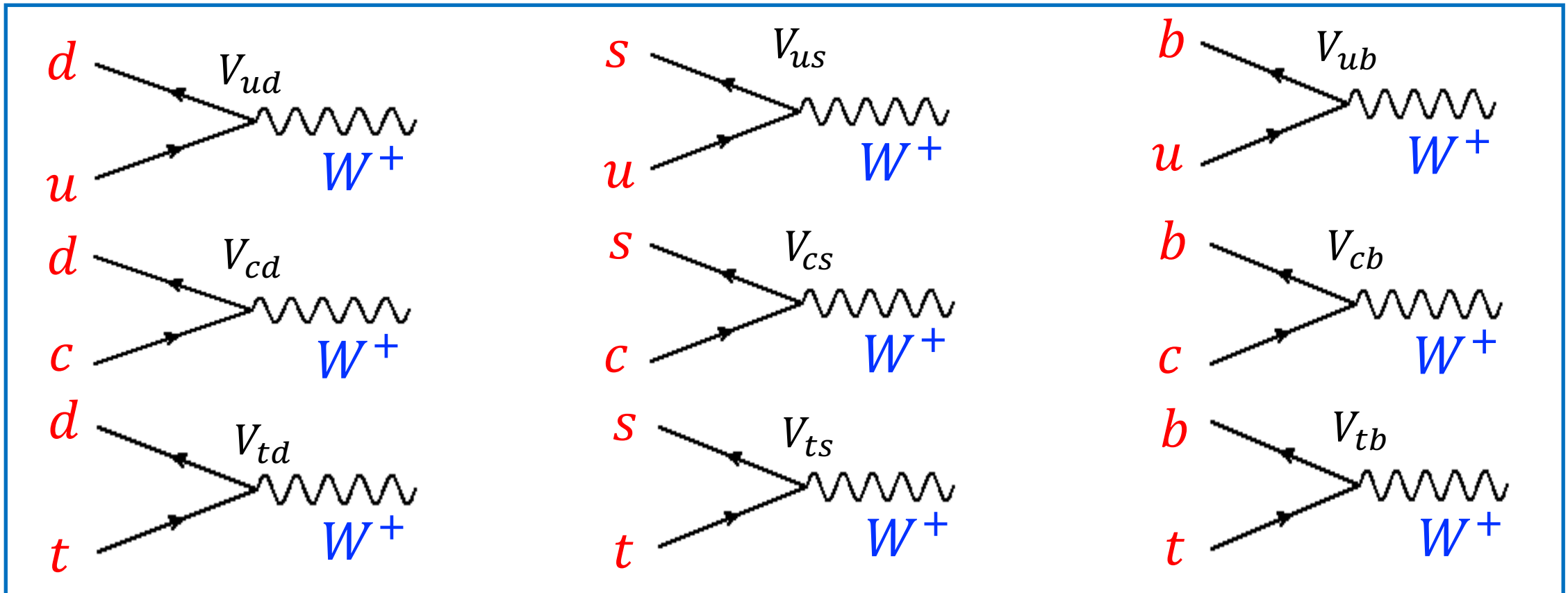


- "Hadronic": $K^+ \rightarrow \pi^+ \pi^+ \pi^-$

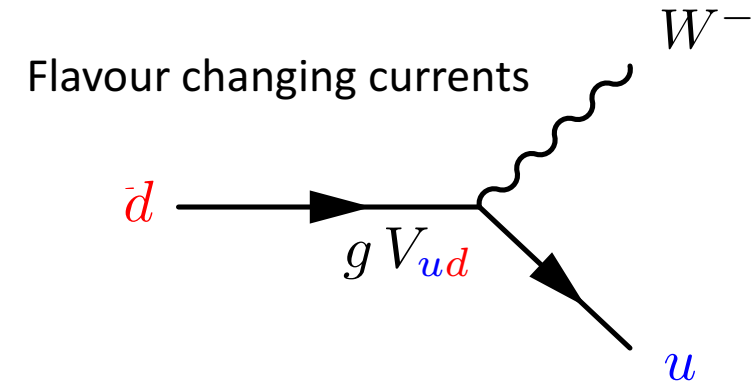
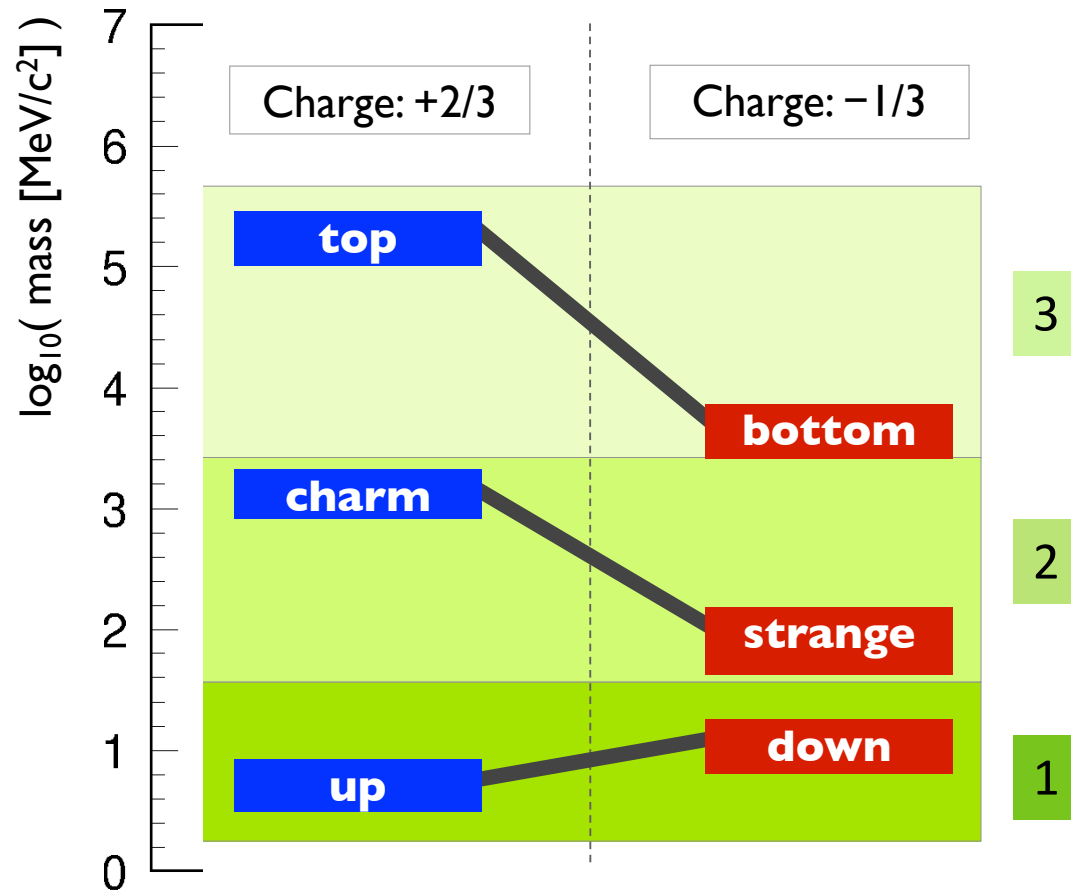


The strange world of W -quark couplings

- All possible combinations of “up”-type to “down”-type quark couplings are possible
 - They are not all equally strong!

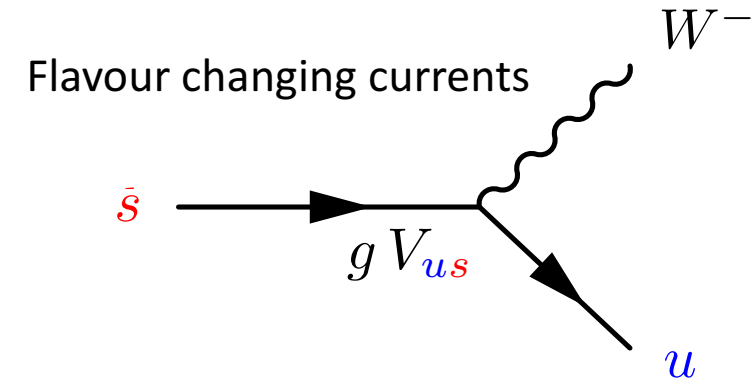
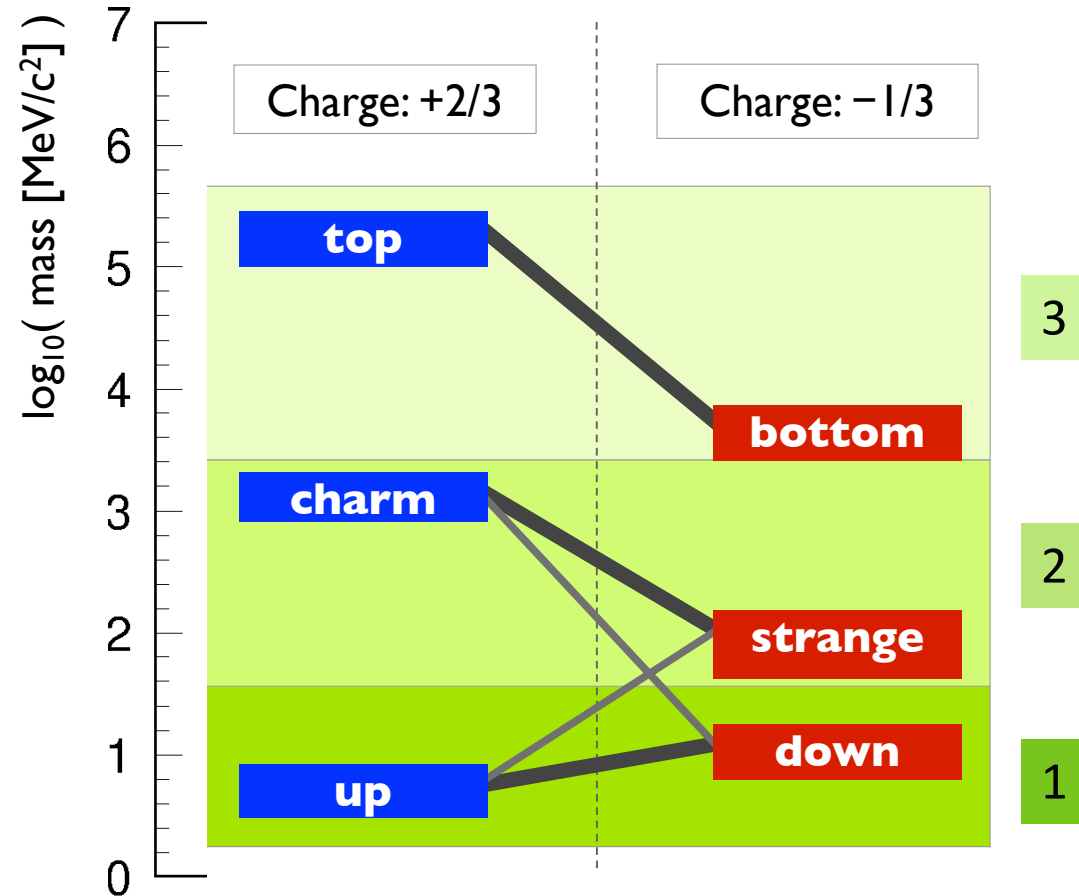


Flavour Changing Quark Interactions



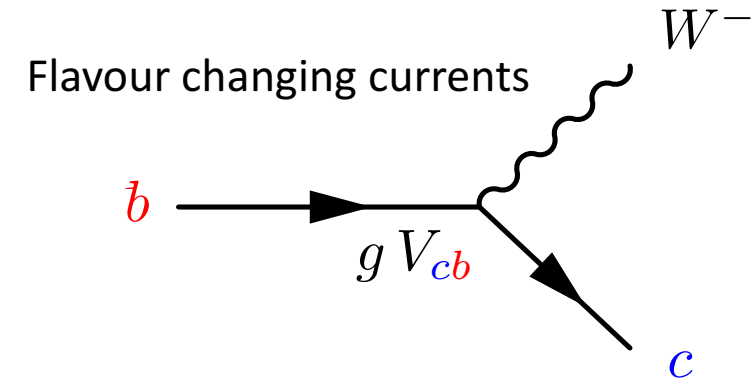
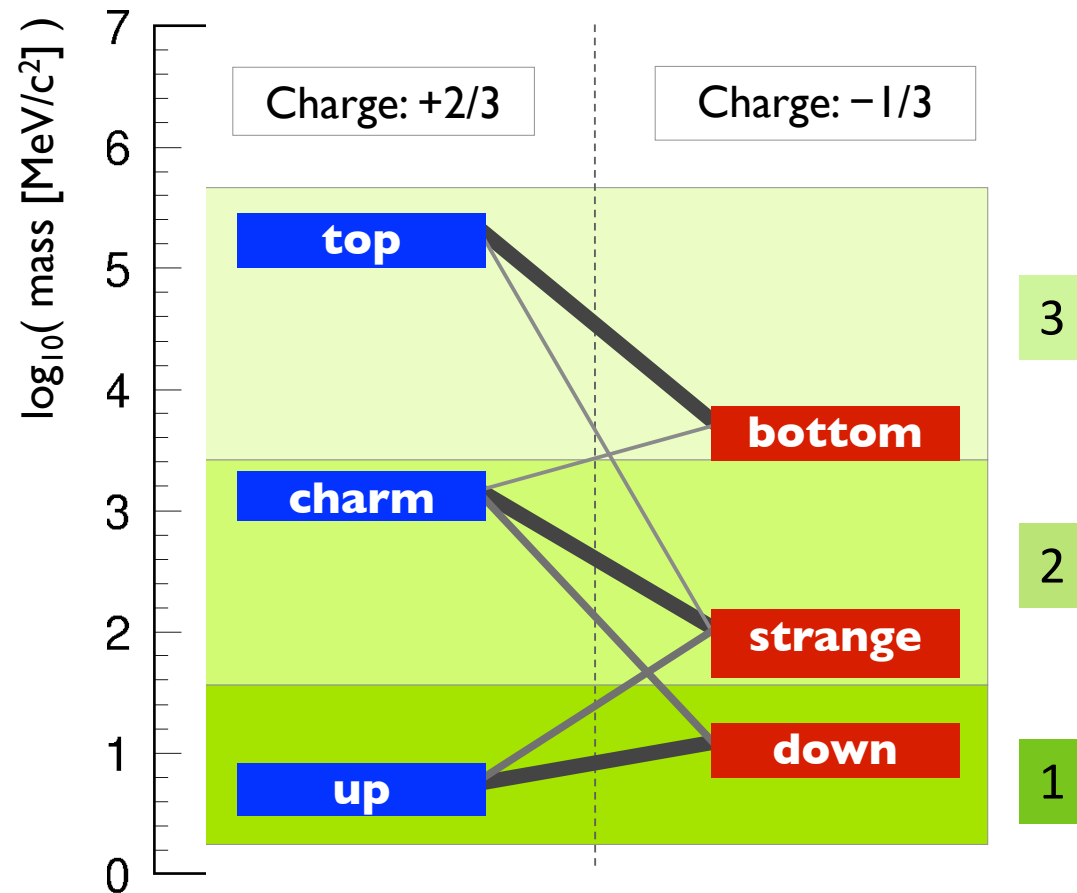
$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & & \\ & V_{cs} & \\ & & V_{tb} \end{pmatrix}$$

Flavour Changing Quark Interactions



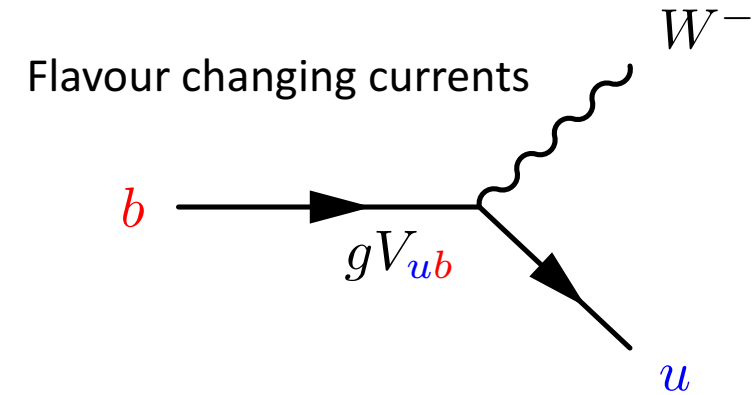
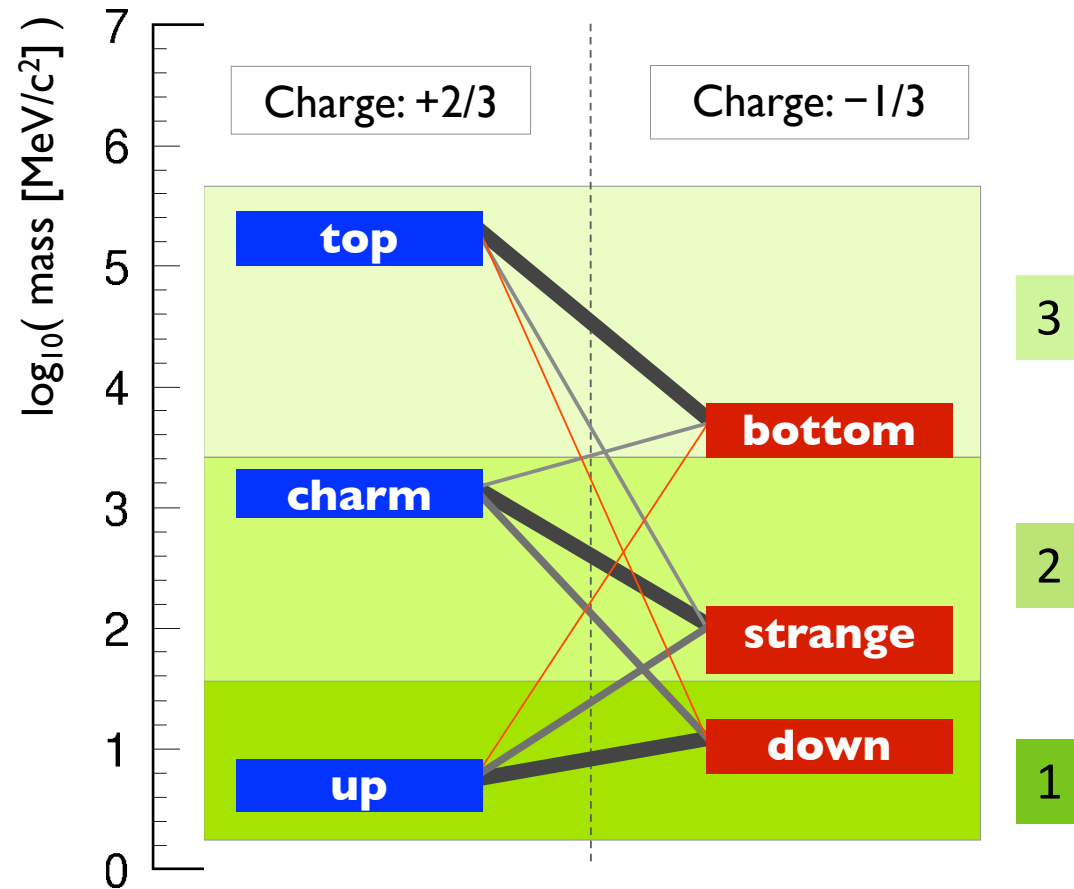
$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & \\ V_{cd} & V_{cs} & \\ & & V_{tb} \end{pmatrix}$$

Flavour Changing Quark Interactions



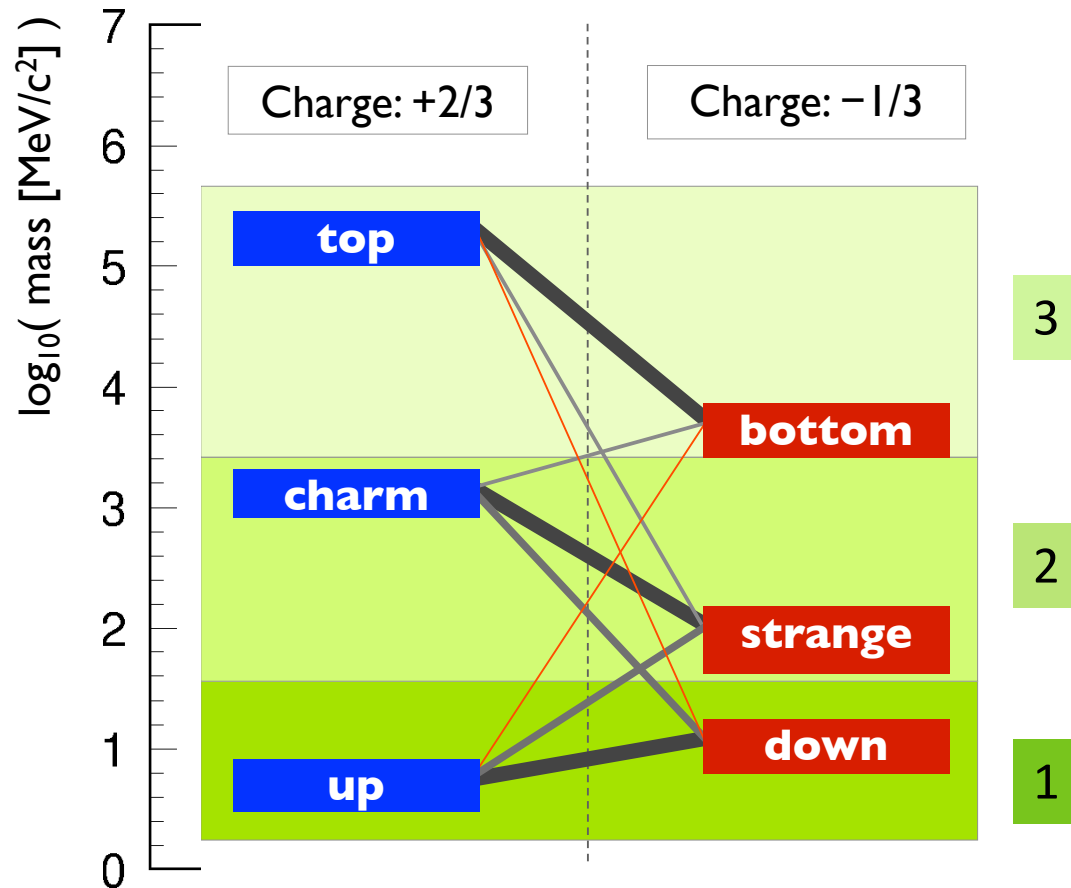
$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Flavour Changing Quark Interactions

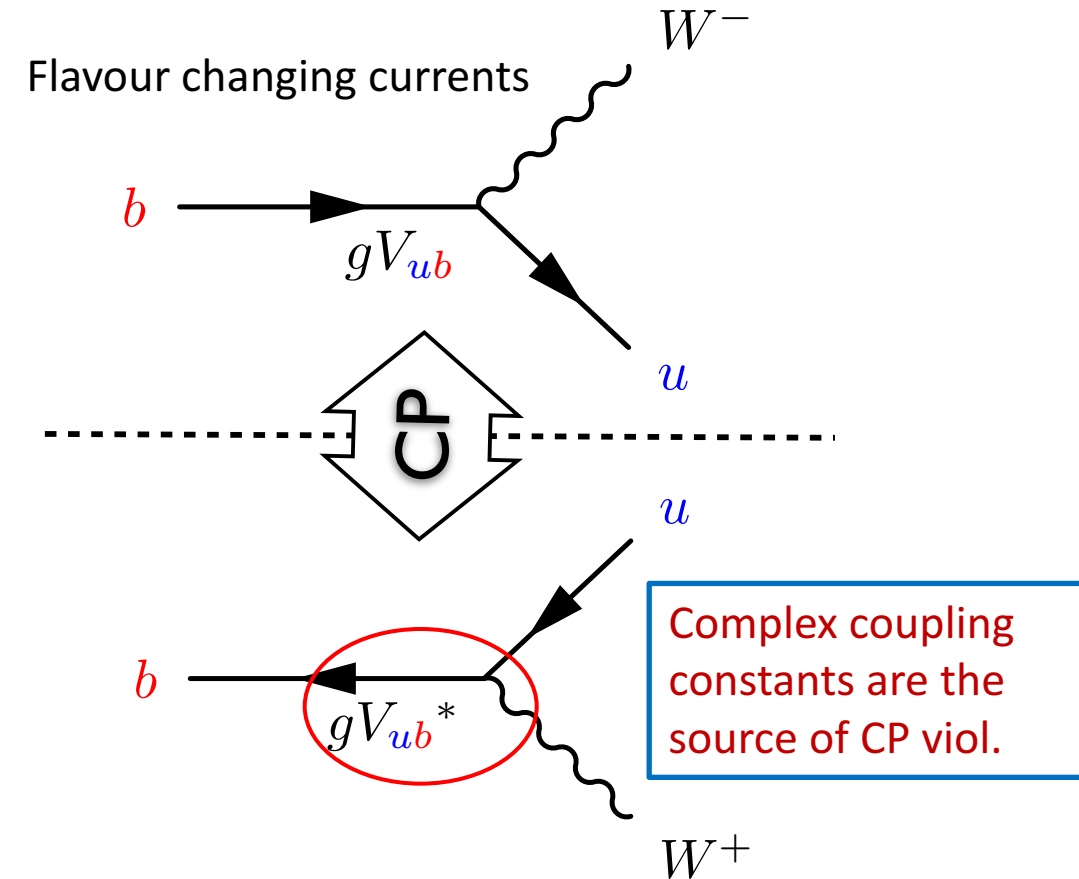


$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Flavour Changing Quark Interactions – CP Violation



- Particles and antiparticles have complex conjugated coupling constants
 - This leads to CP violation
 - Matter dominated universe



$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Why a matrix notation?

- Model:

- Charged weak current does **not** couple to

$$\begin{pmatrix} u \\ d \end{pmatrix}, \begin{pmatrix} c \\ s \end{pmatrix}, \begin{pmatrix} t \\ b \end{pmatrix}$$

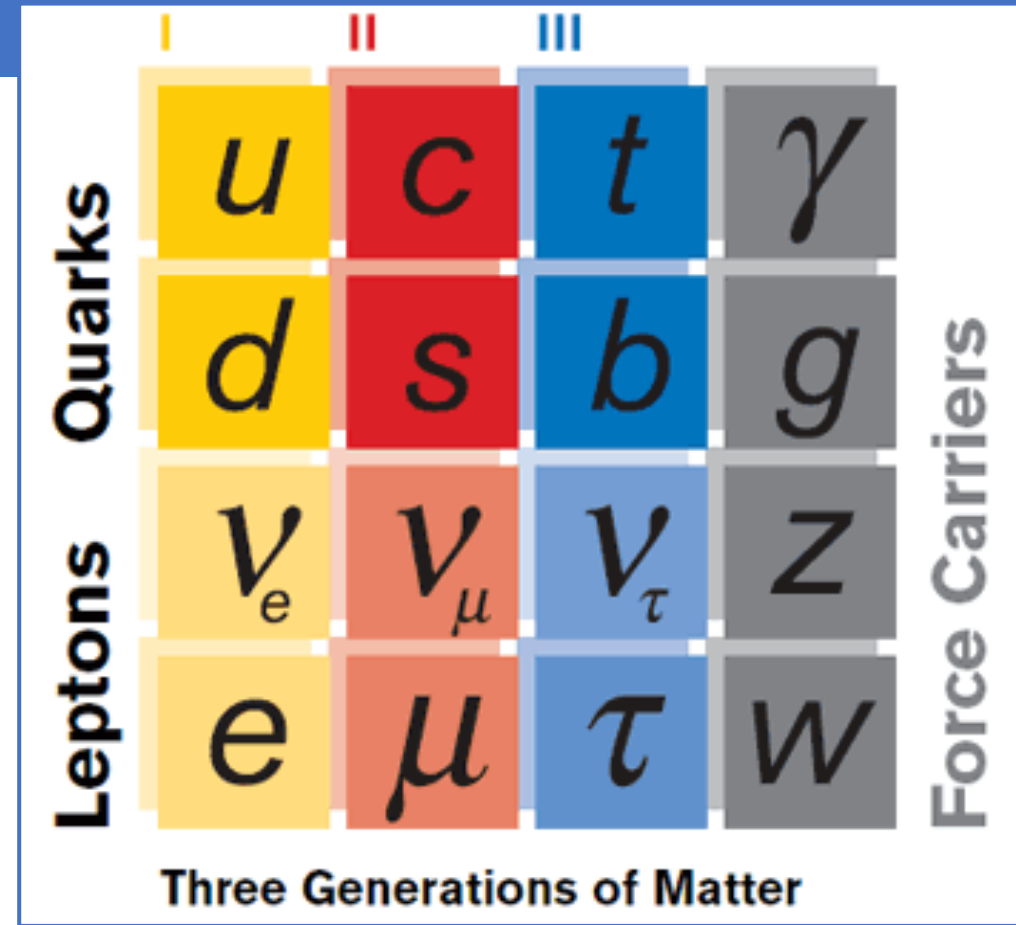
- but instead to

$$\begin{pmatrix} u \\ d' \end{pmatrix}, \begin{pmatrix} c \\ s' \end{pmatrix}, \begin{pmatrix} t \\ b' \end{pmatrix}$$

- Where

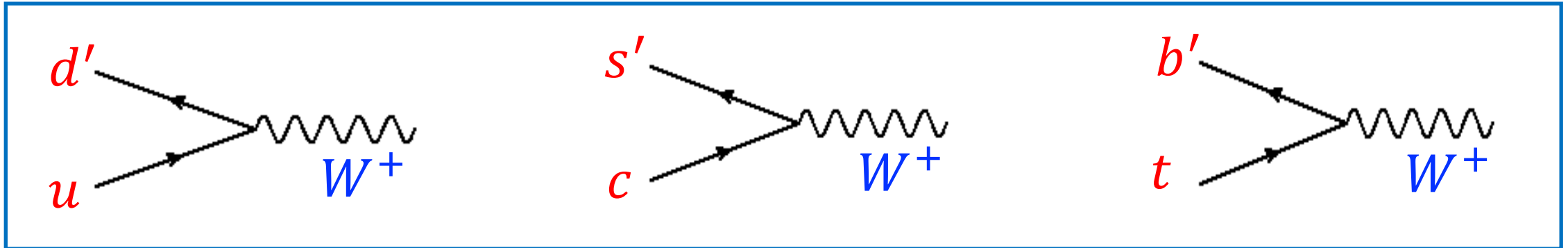
$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \underbrace{\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}}_{\text{CKM matrix}} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

CKM matrix:
Cabibbo,
Kobayashi
Maskawa

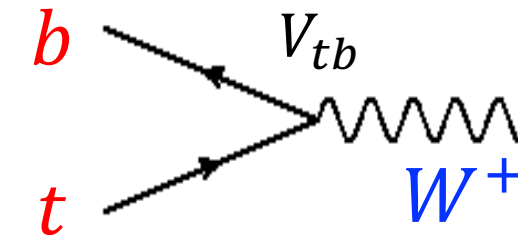
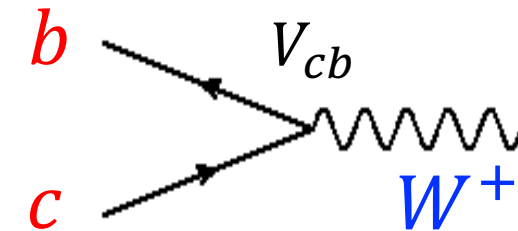
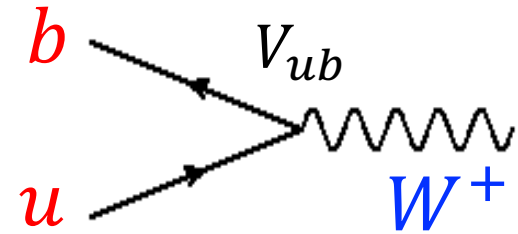
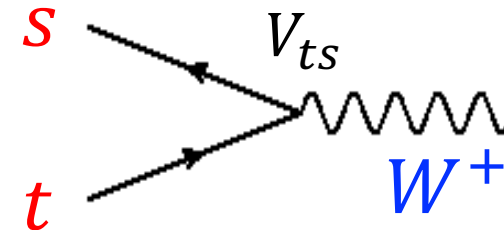
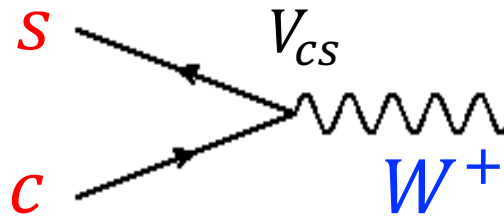
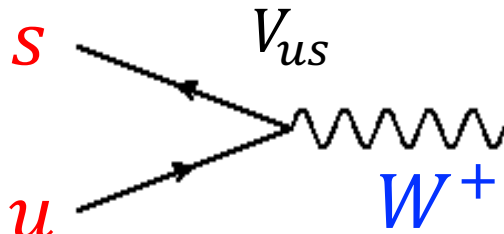
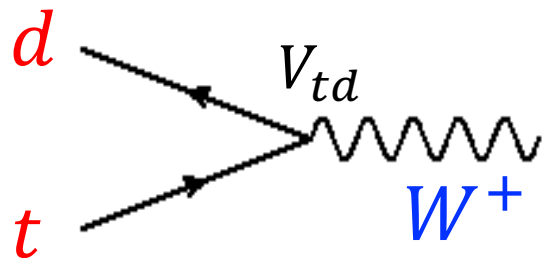
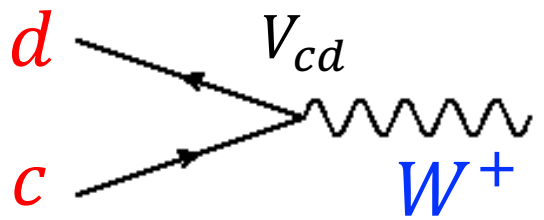
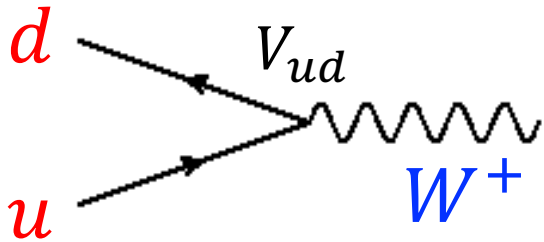
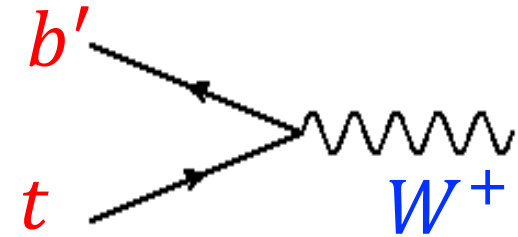
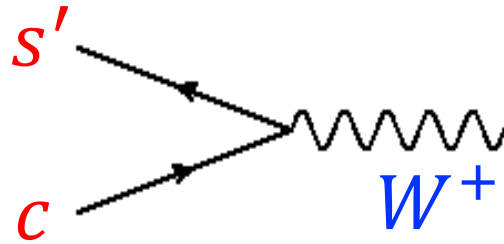
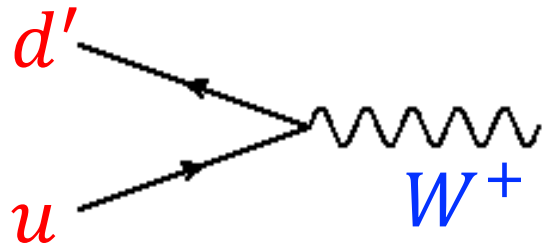


A story of eigenstates

- *Mass* eigenstates $|q\rangle$ are the eigenstate solutions of the free Hamiltonian
- *Weak* or *flavour* eigenstates $|q'\rangle$ are the eigenstate solutions of the weak interaction Hamiltonian.
 - They are unitary linear combination, or “rotation” of mass eigenstates.
- The the weak interaction can be written as:



Flavour eigenstates and Mass eigenstates



The CKM matrix V_{CKM} - 3 vs 2 Generations

$$V_{CKM}: \begin{array}{c} \\ u \\ c \\ t \end{array} \begin{array}{ccc} d & s & b \\ \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \end{array}$$

- Wolfenstein parametrization: $V_{CKM} =$

$$\begin{pmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{1}{2}\lambda^2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

→ 1 CP violating phase

$$V_{CKM}: \begin{array}{c} \\ u \\ c \end{array} \begin{array}{cc} d & s \\ \begin{pmatrix} V_{ud} & V_{us} \\ V_{cd} & V_{cs} \end{pmatrix} \end{array}$$

$$V_{CKM} =$$

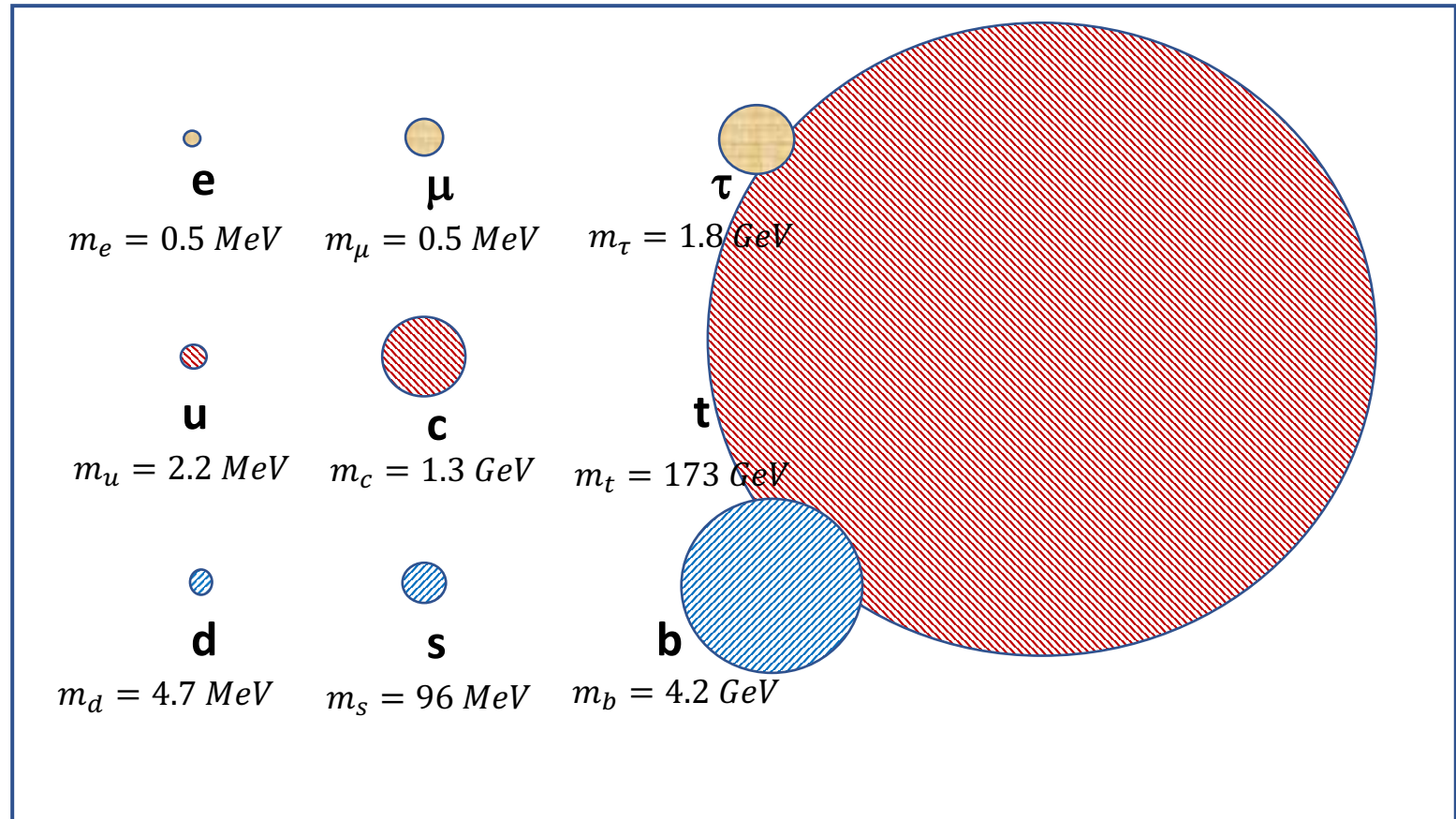
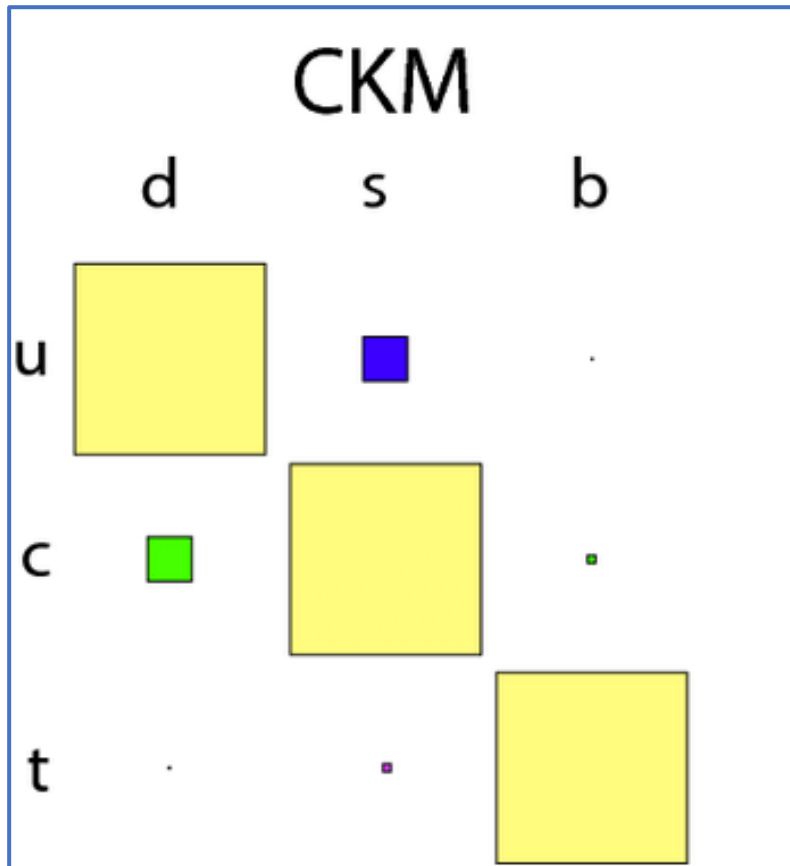
$$\begin{pmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda \\ -\lambda & 1 - \frac{1}{2}\lambda^2 \end{pmatrix}$$

→ No CP violation

- 3 generations is the minimal particle content to generate CP violation (In Standard Model).

The Flavour Puzzle

- Why 3?
- Why are the couplings what they are?
- Is there a relation with the masses of the quarks?



Flavour in the leptons!

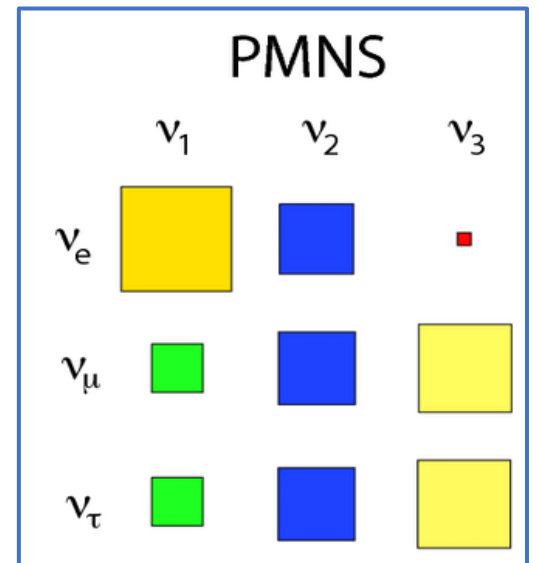
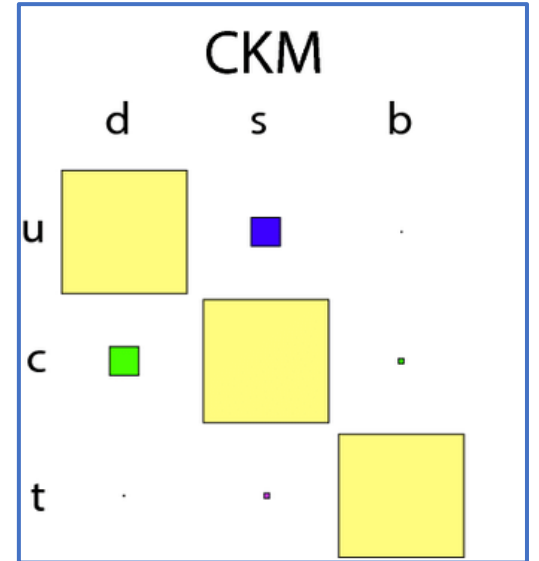
- It turns out neutrino's have mass, too!
 - The mass is very tiny
- The generation mixing also occurs for neutrino's
 - Slightly different nomenclature:

• Quarks (CKM):
$$\begin{pmatrix} u' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

• Leptons (PMNS):
$$\begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} = \begin{pmatrix} U_{1e} & U_{1\mu} & U_{1\tau} \\ U_{2e} & U_{2\mu} & U_{2\tau} \\ U_{3e} & U_{3\mu} & U_{3\tau} \end{pmatrix} \begin{pmatrix} e \\ \mu \\ \tau \end{pmatrix}$$

- U_{PMNS} : Pontecorvo, Maki, Nakagawa, Sakata mixing matrix

- (Difficult) Question: why is lepton mixing not seen in decays?
 - It is only seen in neutrino oscillations

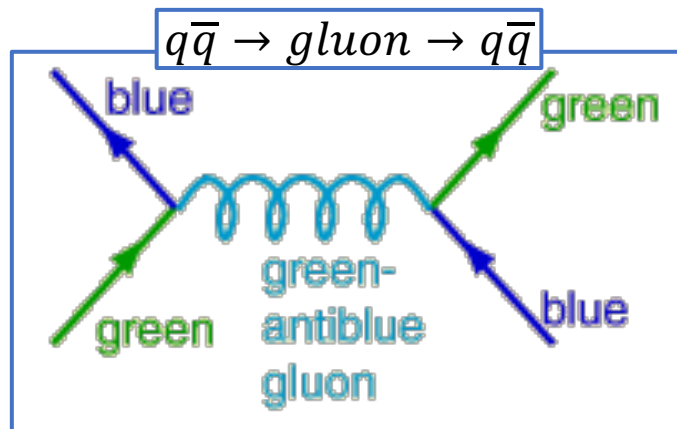


The Strong Interaction

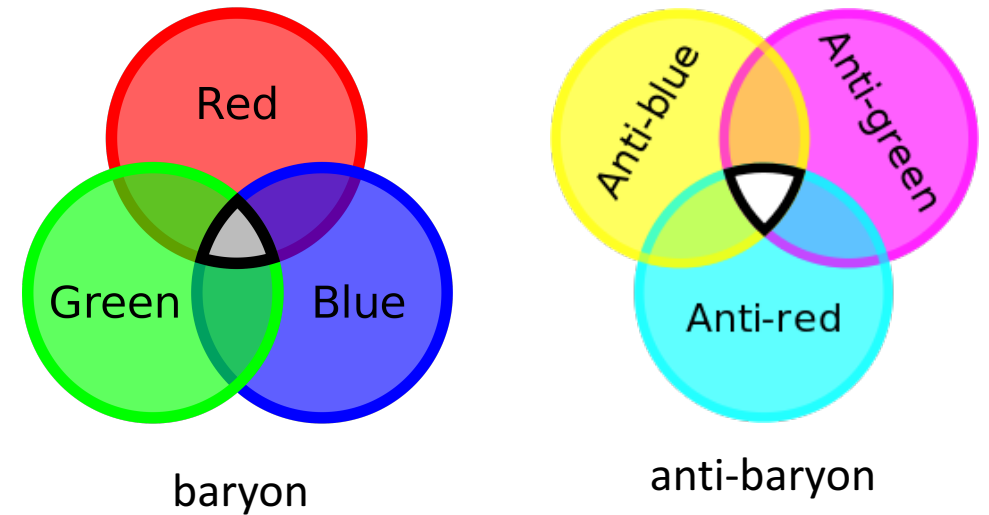
Quantum Chromodynamics (QCD)

Color singlets

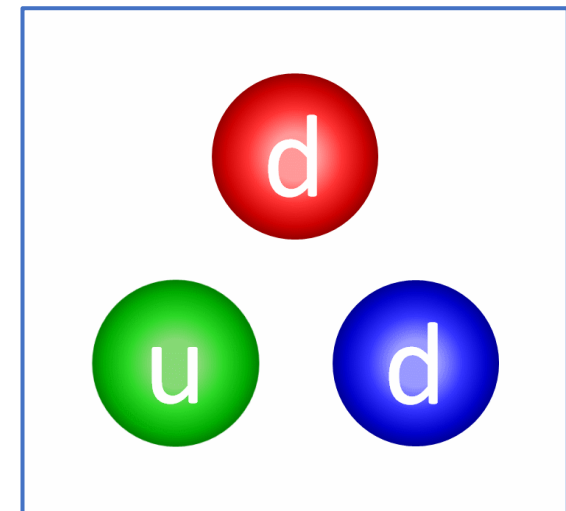
- Quarks are “locked-up” in hadrons
 - Technical term: confinement
- Quarks carry color charge: r, g, b
- All physical objects are color neutral or color *singlets*: *confinement!*
 - Baryons: rgb or $\bar{r}\bar{g}\bar{b}$
 - Mesons: $r\bar{r}$ or $b\bar{b}$ or $g\bar{g}$
- The color force, transmitted by gluons, is very strong
 - Requires a lot of energy to separate a color charge from a color neutral object
 - Trying to do so will produce another color neutral object



Gluon color: $g\bar{b}$
There are 8 gluons
+ 1 colorless singlet



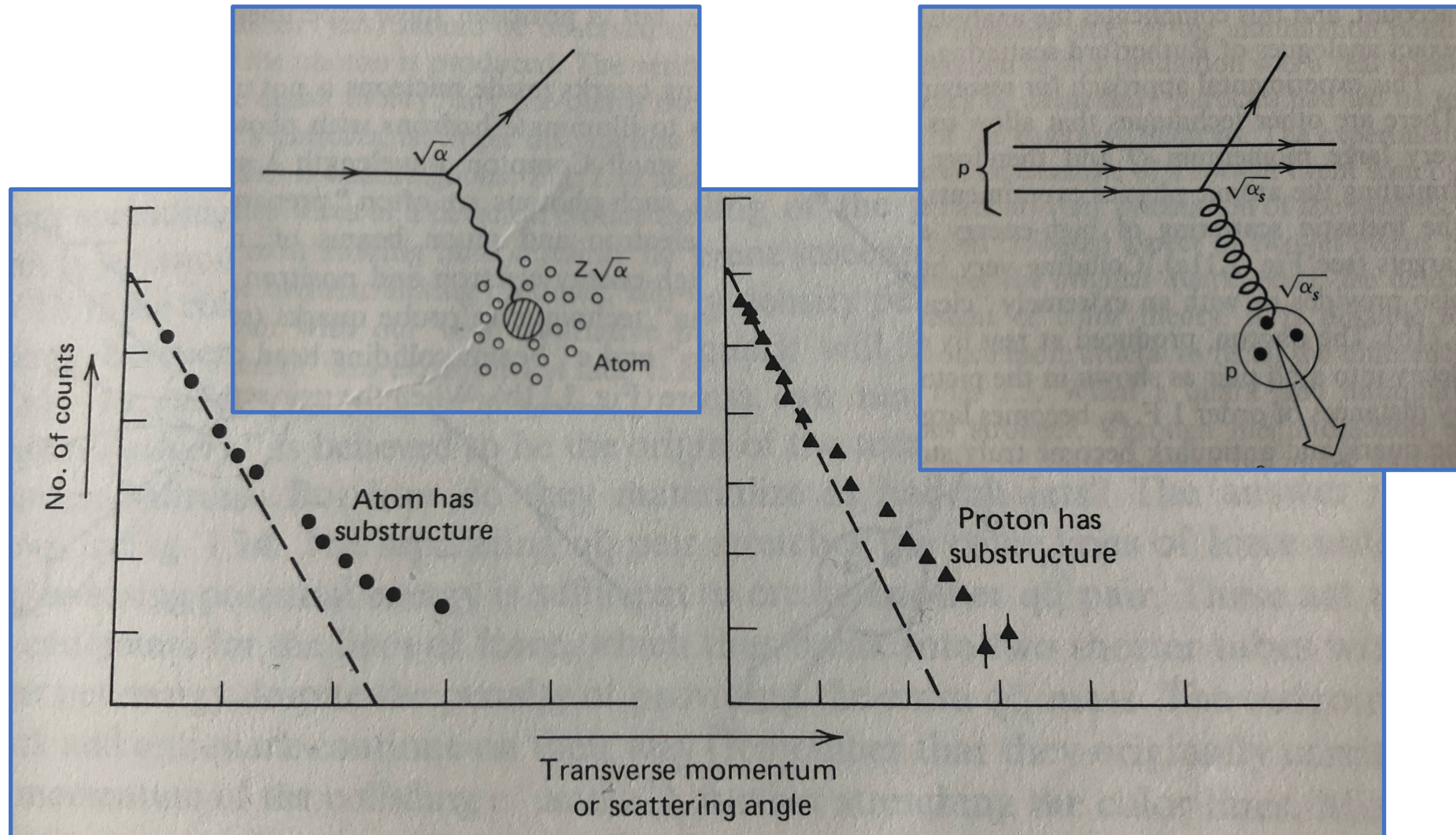
- Animation of a color neutral neutron:



- Are quarks and gluons real or a bookkeeping device?

Proton Substructure: discovery of “partons” or “quarks”(1968)

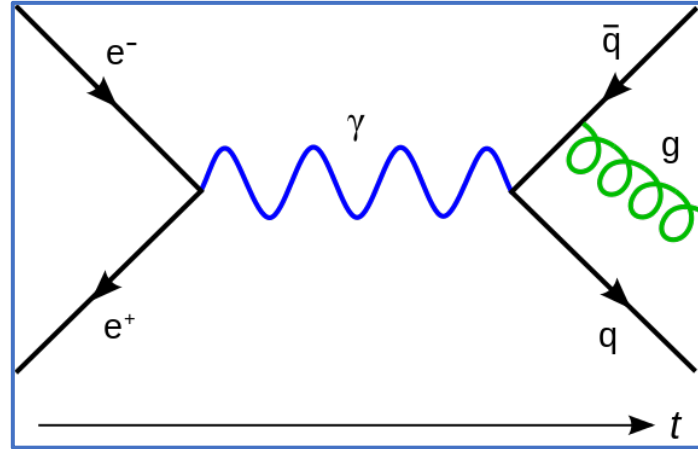
- Similarly to Rutherford scattering (“substructure of the nucleus”) deep inelastic scattering of electrons on protons show the proton substructure



Discovery of the Gluon (1979)

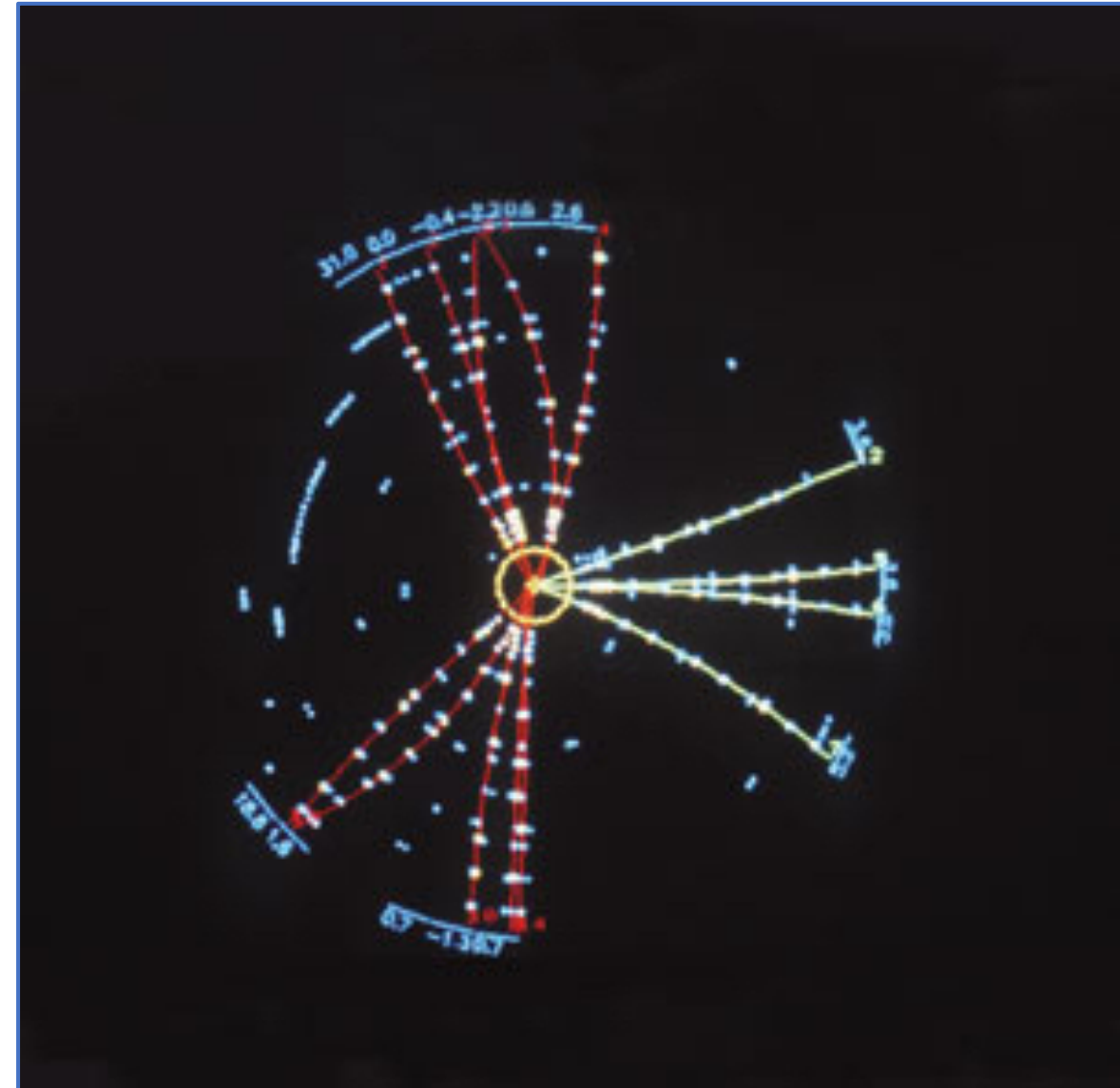
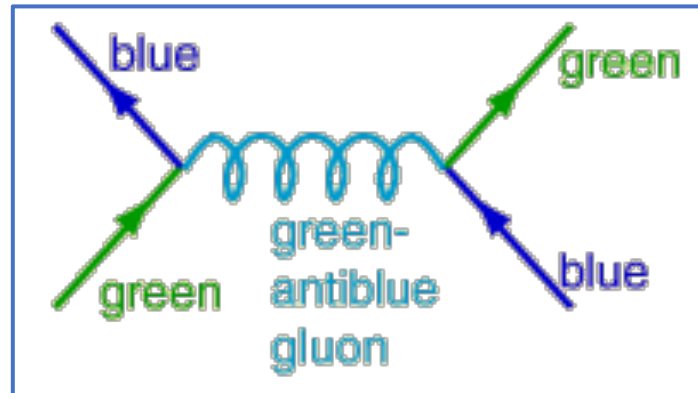
- The gluon was discovered in e^+e^- collisions at DESY in Hamburg

- “3-jet events”

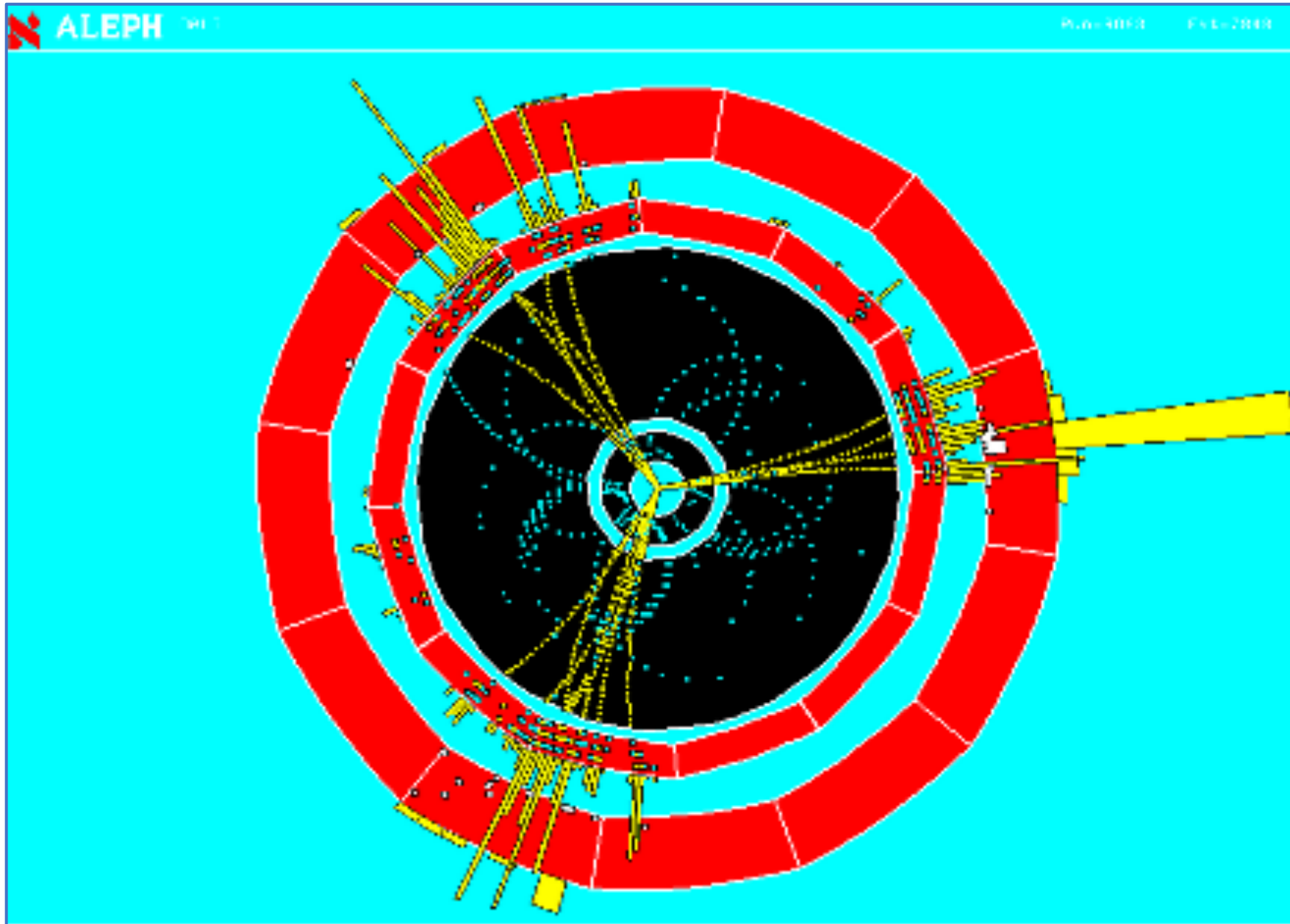


- Feynman diagram for quark-quark interaction

- At LHC

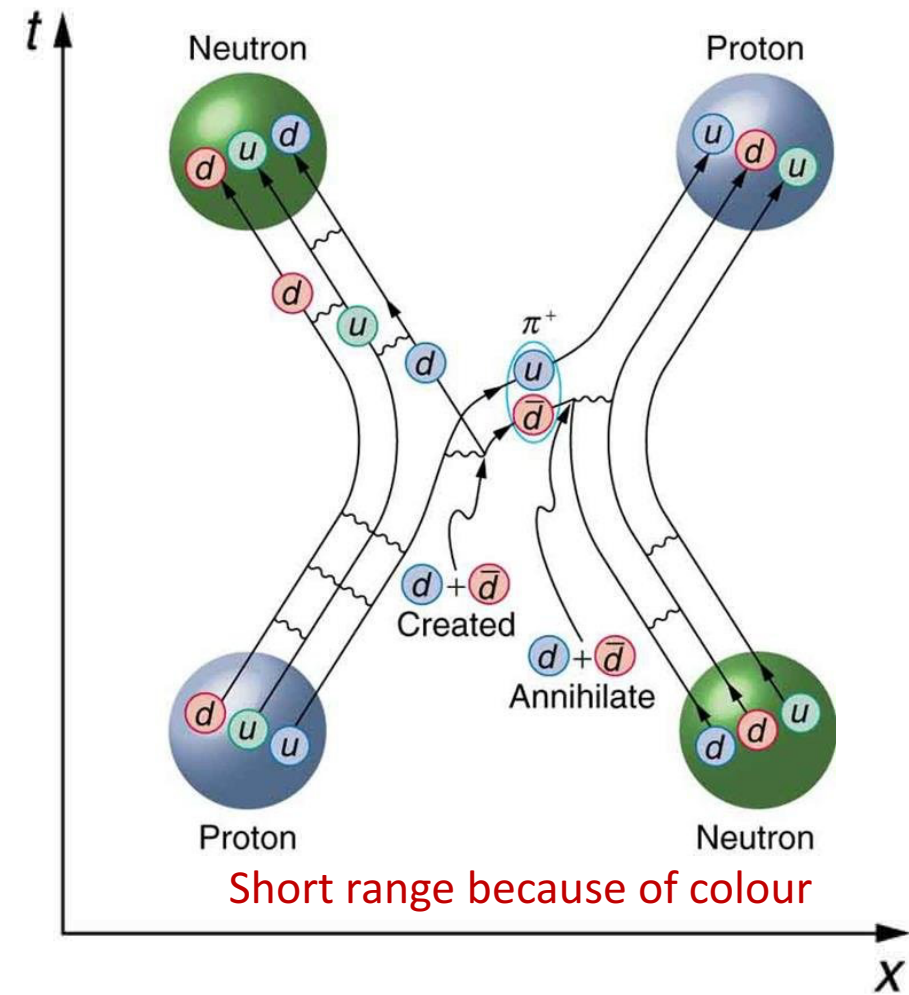
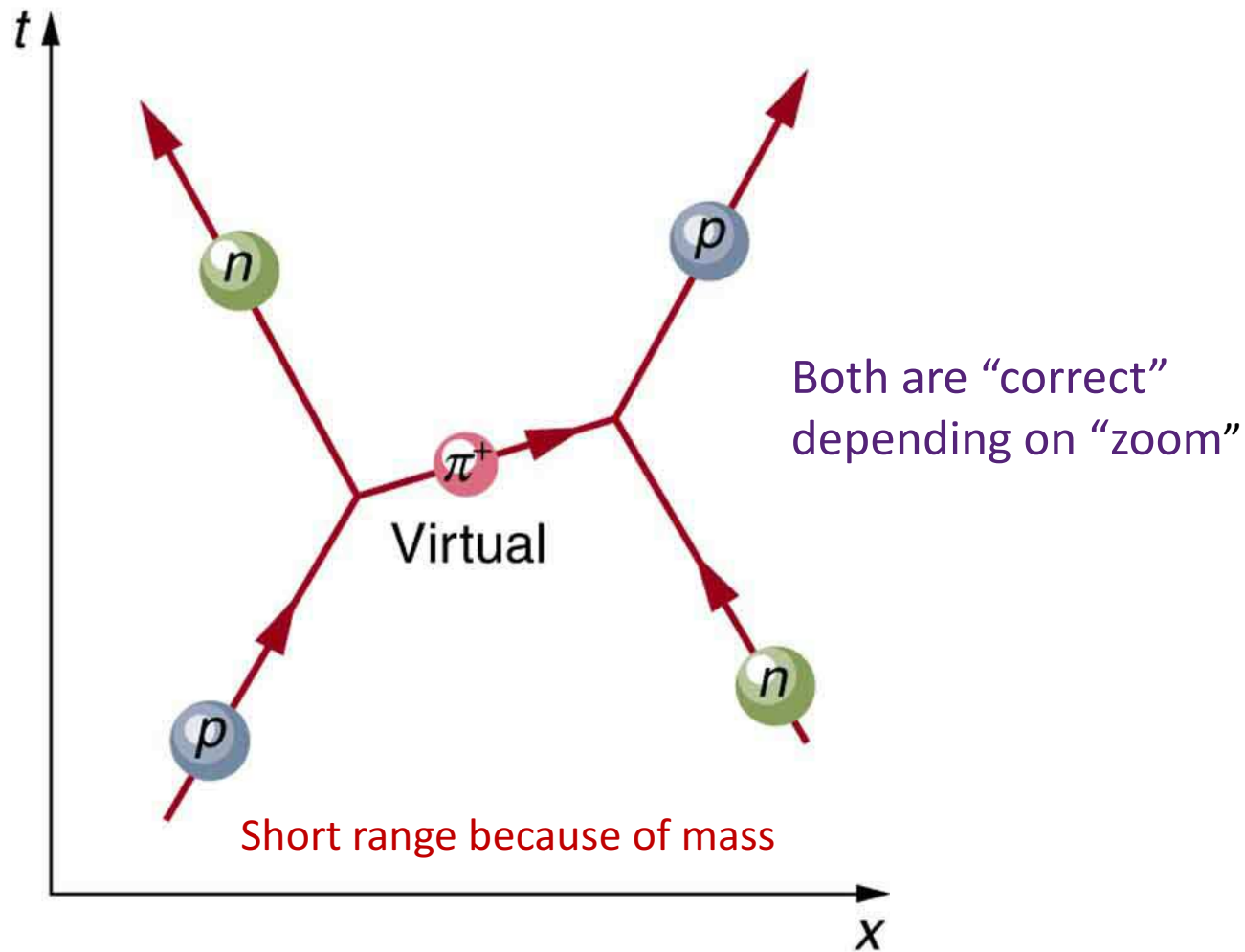
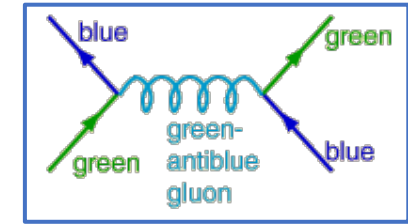
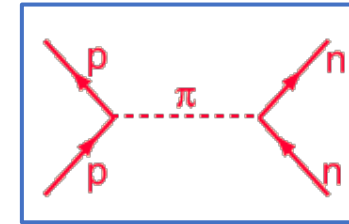


3-Jet event at LEP - Delphi (1989 – 2000)



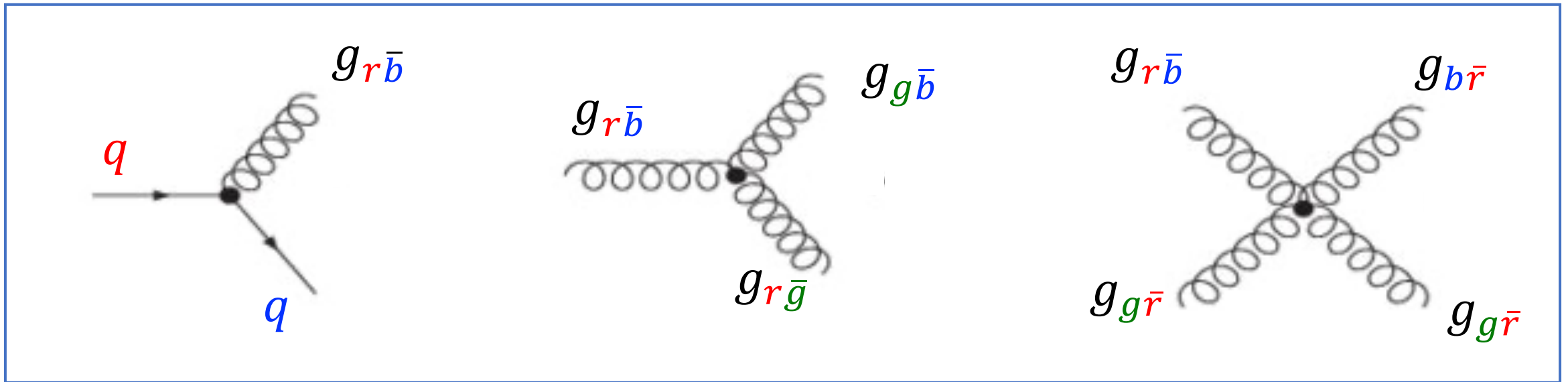
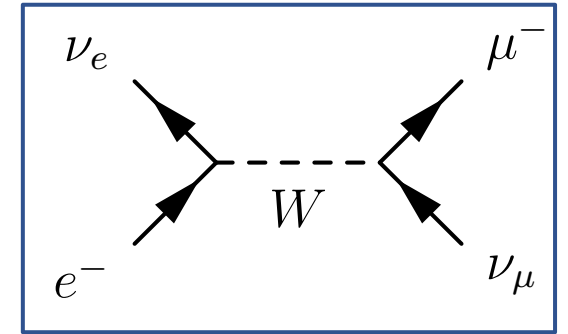
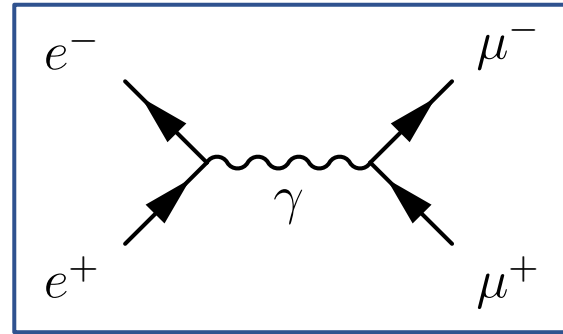
Yukawa's pion exchange vs gluon exchange

- Yukawa: Nuclear force is carried by *massive* gluon
- QCD: Strong force is mediated by *massless* gluon



Fundamental Vertices of the strong interaction

- QED: photon couples to charge
- Weak: W couples to isospin
- Strong gluon couples to color
 - 3 colors and anti-colors
 - Gluon carries charge itself!
 - 8 different gluons (9 – 1 symmetric)

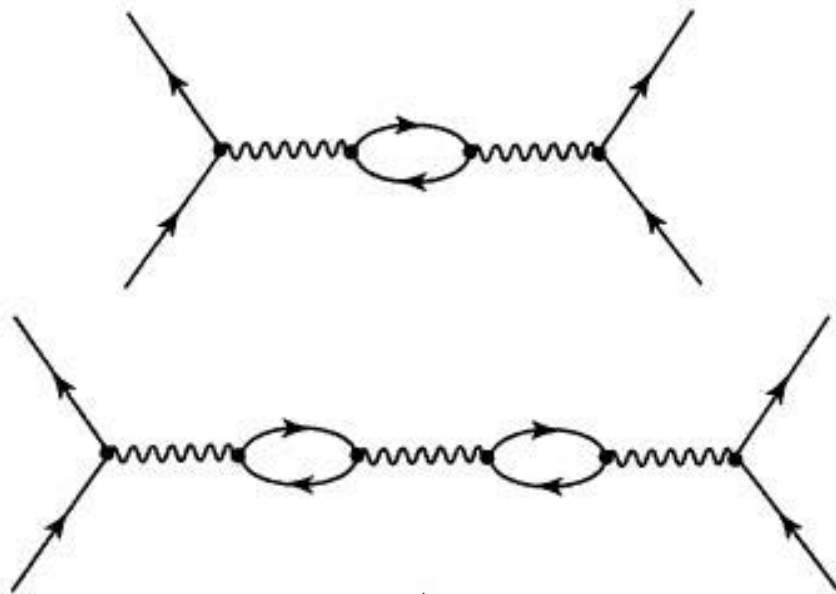


- Leptons carry no color: leptons do not feel strong (nuclear) interaction

Abelian and non-Abelian

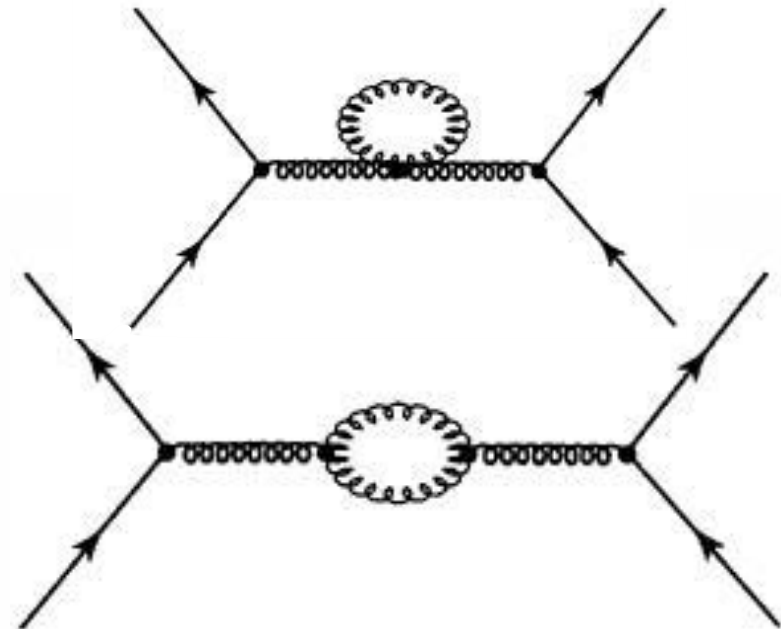
- Photons have no electric charge:
 - Photons *do not* interact with each other (“Abelian” in group theory)
- Gluons carry color charge:
 - Gluons *do* interact with each other (“non-Abelian” in group theory)

QED, QCD



...etc.

QCD



- Vacuum behaves like dielectric: by QED and even more so by QCD

Screening

- Charge screening effect:
 - Around a charge $q(+)$ dipoles align
 - Halo of negative charge
 - Effective charge (at distance) reduced

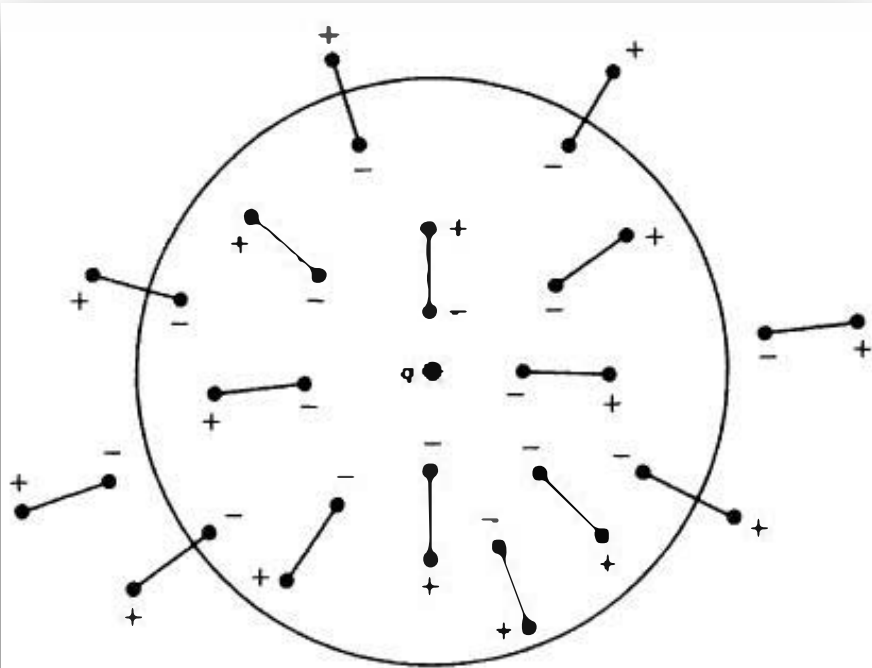
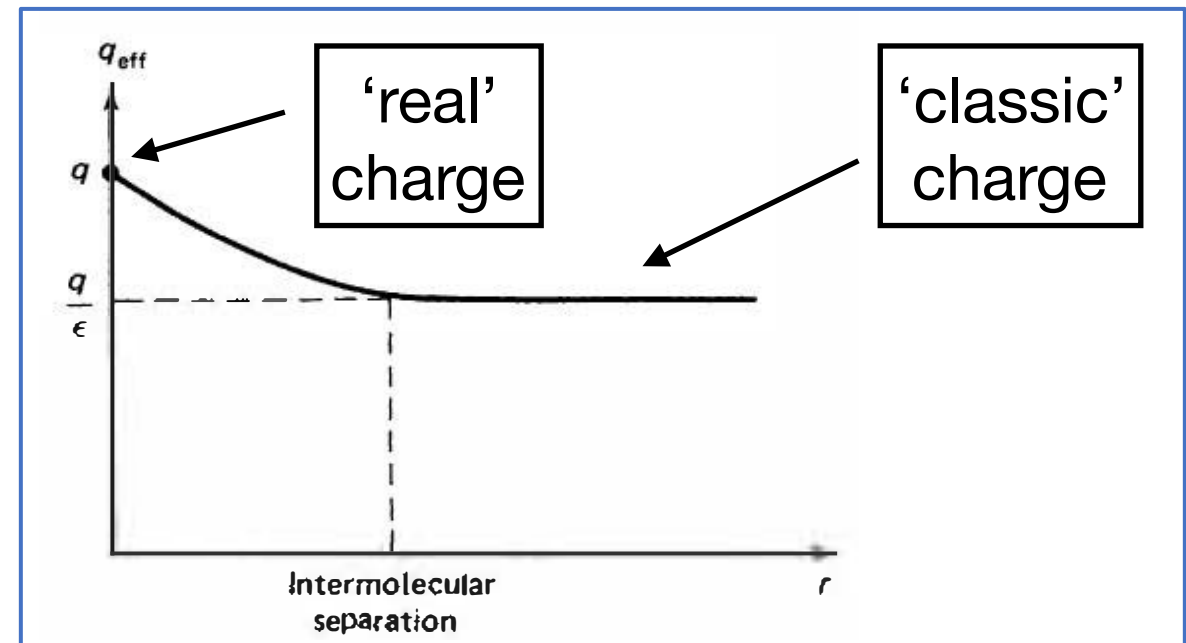
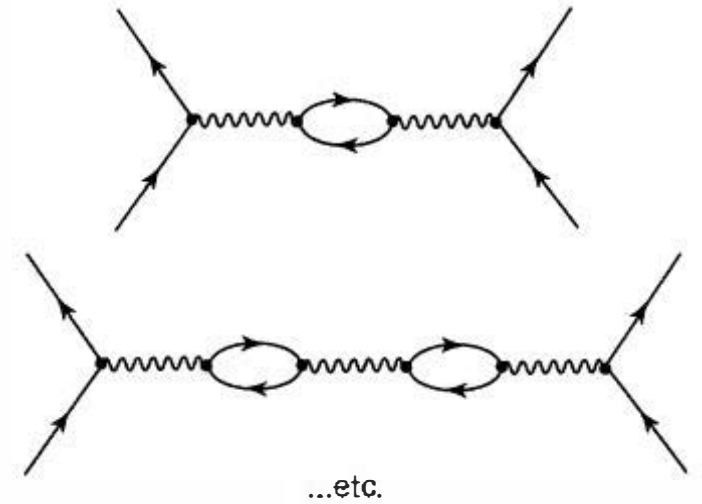


Fig. 2.1 Screening of a charge q by a dielectric medium.

QED, QCD



Anti-screening and asymptotic freedom

- Anti-screening:

- gluon self coupling loops have opposite effect

- Crucial parameter:

$$a \equiv 2f - 11c$$

where: f = number of flavours

c = number of colors

- In case $a > 1$ (QED) coupling *increases* for short distance:

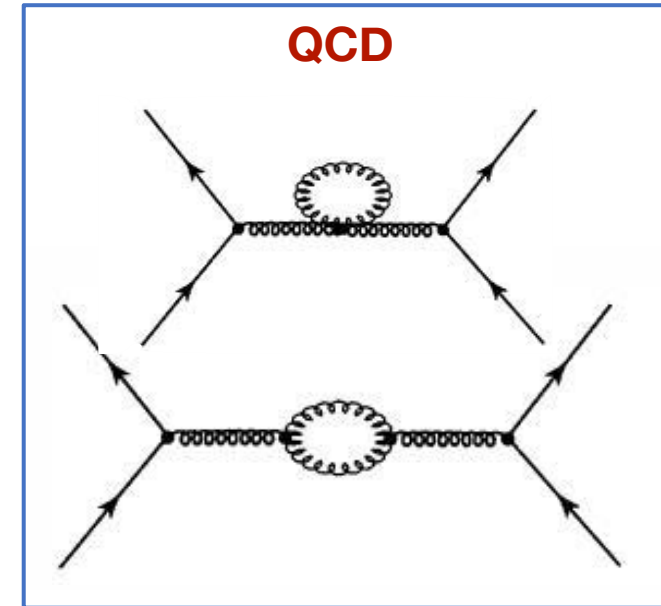
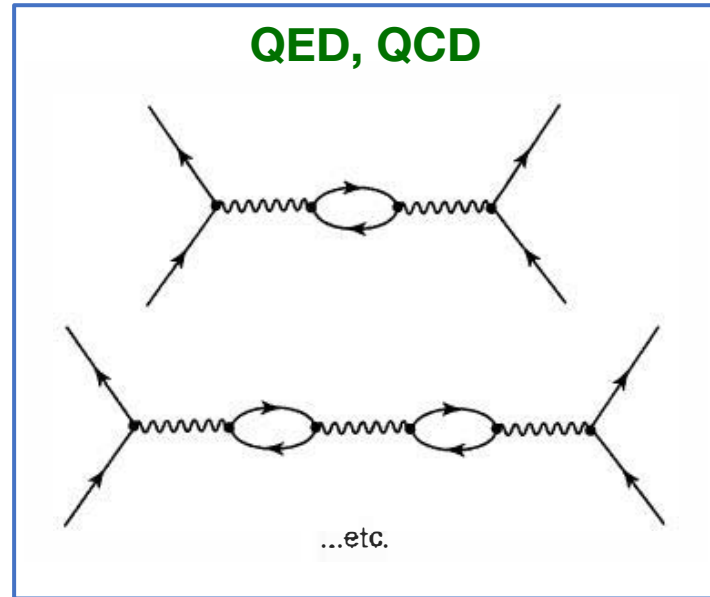
- In case $a < 1$ (QCD) coupling *decreases* for short distance

- Asymptotic freedom: at short distance:

- Coupling $\alpha_s \sim 0$ at very small distances

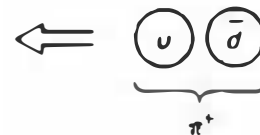
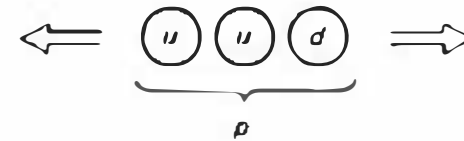
- Quarks and gluons become “free”

➔ Quark gluon plasma



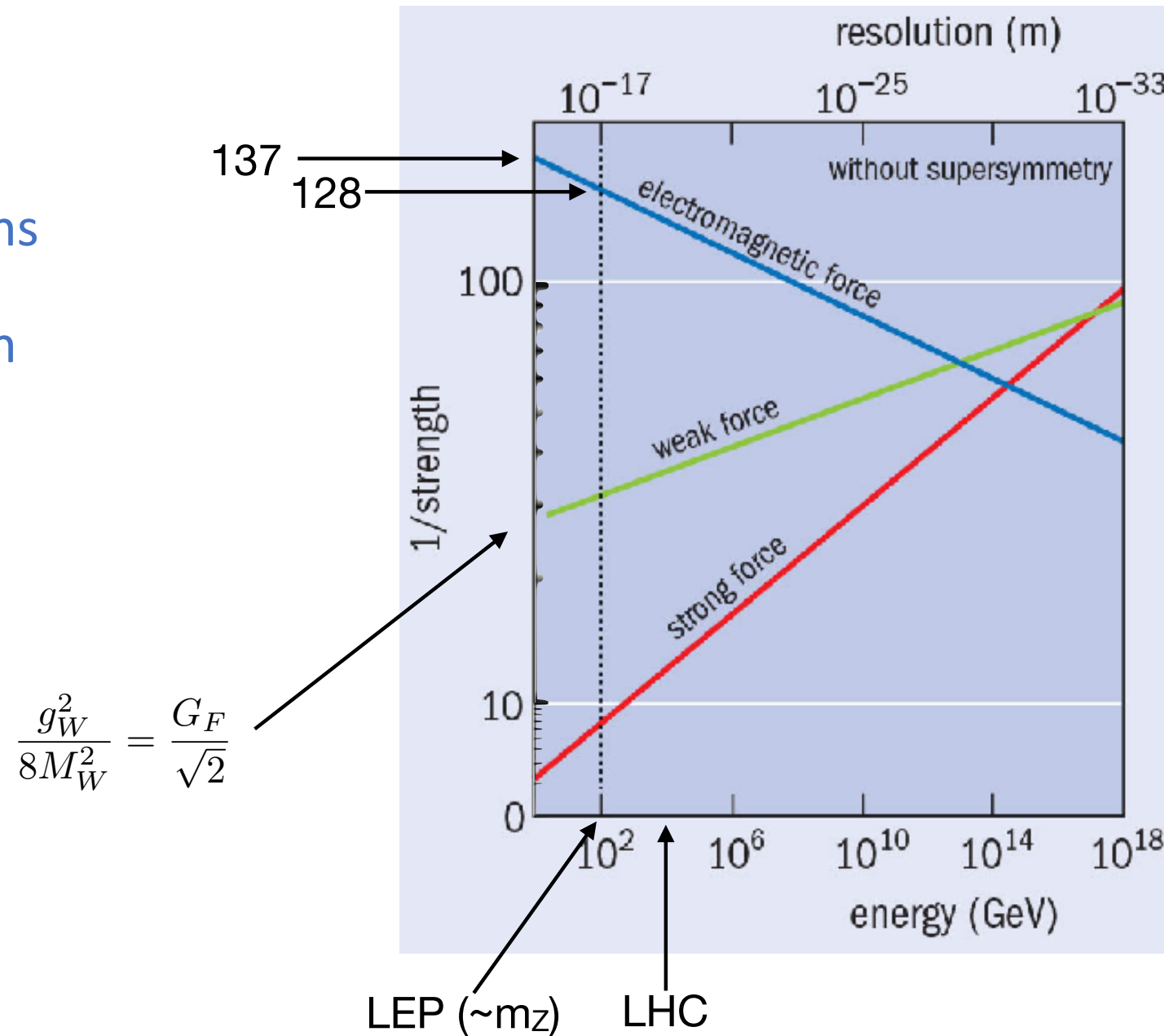
- Confinement: at large distance:

- $\alpha_s > \sim 1$



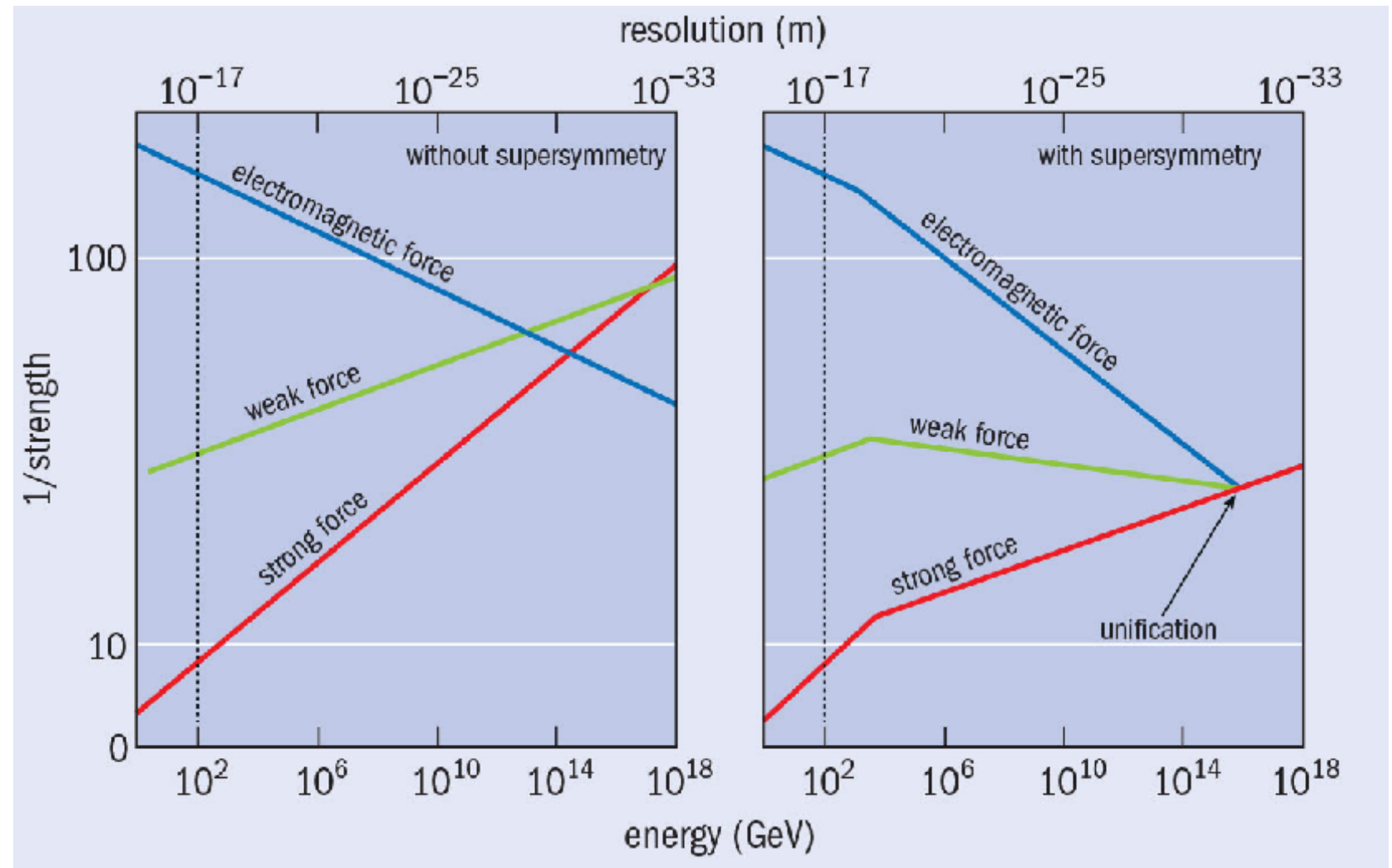
Running coupling constants and unification?

- Running couplings: 'beta'-function
- Higher energy means shorter distance
- Lines do not cross in one point



Running coupling constants and unification?

- Unification with SUSY?
- 10^{15} GeV
 - Grand Unified Theory
- 10^{19} GeV:
 - Planck scale \rightarrow quantum gravity?



- 10^{15} GeV \rightarrow Grand Unified Theory?
- Planck scale ($\sim 10^{19}$ GeV) \rightarrow Gravity! \rightarrow Theory of Everything?

Standard Model of particles and forces

Standard Model of Elementary Particles

three generations of matter (elementary fermions)						three generations of antimatter (elementary antifermions)						interactions / force carriers (elementary bosons)		
		I	II	III			I	II	III					
mass	charge	spin	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0		$\approx 124.97 \text{ GeV}/c^2$			
			$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	0		0		0	
			$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1		0		0	
quarks			u up	c charm	t top	\bar{u} antiup	\bar{c} anticharm	\bar{t} antitop	g gluon	H higgs				
			$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0					
			$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}$	0					
			$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1					
			d down	s strange	b bottom	\bar{d} antidown	\bar{s} antistrange	\bar{b} antibottom	γ photon					
			$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 91.19 \text{ GeV}/c^2$					
			-1	-1	-1	1	1	1	0					
			$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1					
			e electron	μ muon	τ tau	e^+ positron	$\bar{\mu}$ antimuon	$\bar{\tau}$ antitau	Z Z ⁰ boson					
			$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$\approx 80.39 \text{ GeV}/c^2$					
			0	0	0	0	0	0	1					
			$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1					
			ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	$\bar{\nu}_e$ electron antineutrino	$\bar{\nu}_\mu$ muon antineutrino	$\bar{\nu}_\tau$ tau antineutrino	W^+ W ⁺ boson					
									W^- W ⁻ boson					

Summary Forces: Intermediate Vector Bosons

- SM: forces are transmitted by vector mesons, particles with spin-1:

- Electromagnetism:

- Long range : photon γ , $M_\gamma = 0$
- Photon carries no E.M. charge

- Weak interaction:

- Very short range: W and Z bosons, $M_W = 80 \text{ GeV}$, $M_Z = 91 \text{ GeV}$
- Weak isospin charge. W and Z have I-charge.

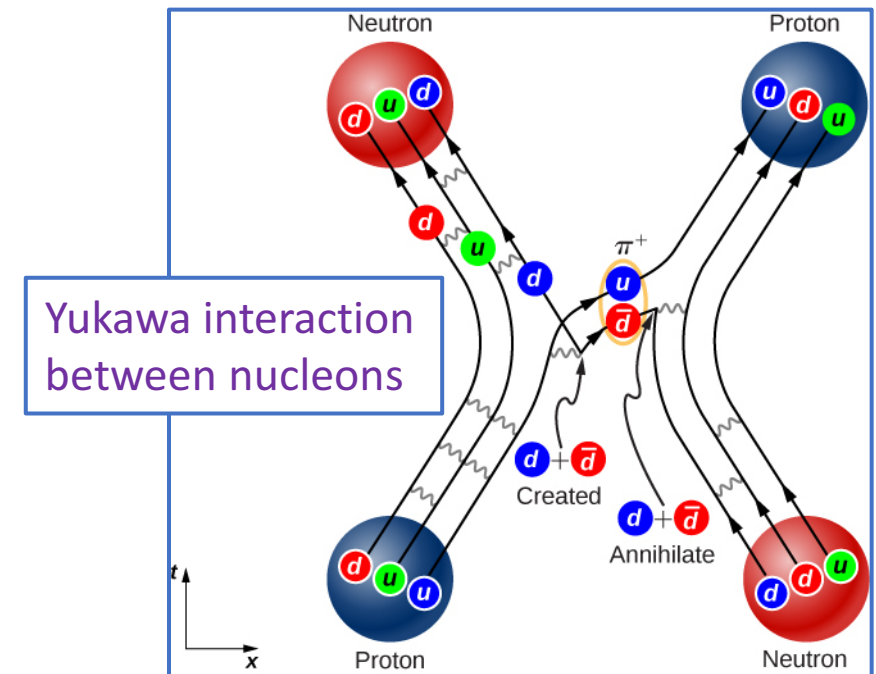
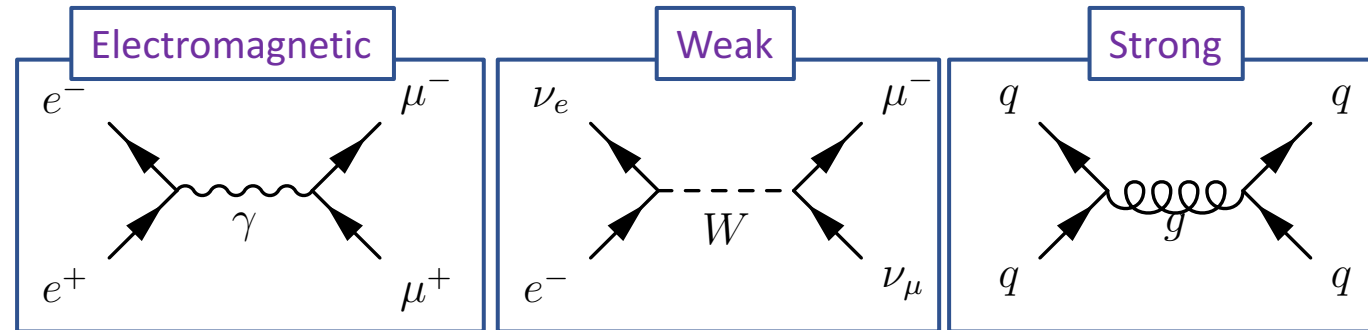
- Strong interaction:

- Very short range: gluon g , $M_g = 0$
 - Gluons carry color \rightarrow confinement \rightarrow very short range
- Pairs of quarks can transmit strong force: Yukawa's mesons
 - Short range

- Gravitation: transmitted by spin-2 particle?

- Graviton

- Long range: $M_g = 0$



Prepare(!): Relativistic and 4-vector notation

See Griffiths chapter 3

- Lorentz transformation along x^1 axis using $\beta = v/c$ and $\gamma = 1/\sqrt{1 - \beta^2}$ is:
$$\begin{aligned}x^{0'} &= \gamma(x^0 - \beta x^1) \\x^{1'} &= \gamma(x^1 - \beta x^0) \\x^{2'} &= x^2 \\x^{3'} &= x^3\end{aligned}$$
- Coordinate four-vector notation: $x^\mu = (x^0, x^1, x^2, x^3)$
 - Time: $x^0 = ct$, Space: $\vec{x} = (x^1, x^2, x^3)$
 - General *Contravariant* vector $A^\mu = (A^0, \vec{A})$ transforms like x^μ
- Lorentz transformations leave the “length” invariant $A \cdot A = |A|^2 = A^{0^2} - |\vec{A}|^2$ which can be regarded as the product of a *contravariant* vector with a *covariant* vector:
$$A_\mu = (A^0, -\vec{A})$$
- Metric tensor: $g_{\mu\nu} = g^{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$
- Scalar product: $A \cdot B = A^\mu B_\mu = g_{\mu\nu} A^\mu B^\nu$, where a sum is always implicit over *contravariant* and *covariant* indices.
- The scalar product is always Lorentz invariant.

Exercise – 7 : 4-Vector notation

- a) Start with the expression for a Lorentz transformation along the x^1 axis. Write down the *inverse* transformation (i.e. express (x^0, x^1) in $(x^{0'}, x^{1'})$)
- b) Use the chain rule to express the derivatives $\partial/\partial x^{0'}$ and $\partial/\partial x^{1'}$ in $\partial/\partial x^0$ and $\partial/\partial x^1$
- c) Use the result to show that $(\partial/\partial x^0, -\partial/\partial x^1)$ transforms in the same way as (x^0, x^1)
- d) In other words the derivative four-vectors transform inversely to the coordinate four-vectors:

$$\partial^\mu = \left(\frac{1}{c} \frac{\partial}{\partial t}, -\vec{\nabla} \right) \text{ and } \partial_\mu = \left(\frac{1}{c} \frac{\partial}{\partial t}, \vec{\nabla} \right)$$

Note the difference w.r.t. the minus sign!

Exercise – 8: Kinematics: Z -boson production

- The Z -boson particle is a carrier of the weak force. It has a mass of 91.1 GeV. It can be produced experimentally by annihilation of an electron and a positron. The mass of an electron, as well as that of a positron, is 0.511 MeV.
 - a) Draw the Feynman interaction diagram for this process.
 - b) Assume that an electron and a positron are accelerated in opposite directions and collide head-on to produce a Z -boson in the lab frame. Calculate the minimal beam energy required for the electron and the positron in order to produce a Z -boson.
 - c) Assume that a beam of positron particles is shot on a target containing electrons. Calculate the beam energy required for the positron beam in order to produce Z -bosons.
 - d) This experiment was carried out in the 1990's. Which method do you think was used? Why?

“Quick” Exercises 9, 10, 11

9. [Griffiths exercise 2.2] “Crossing lightsabers”

- Draw the lowest-order Feynman diagram representing Delbruck scattering: $\gamma + \gamma \rightarrow \gamma + \gamma$
- This has no classical analogue. Explain why.

10. [Griffiths exercise 2.4]

- Determine the invariant mass of the virtual photon in each of the lowest-order Feynman diagrams for Bhabha scattering

11. [Griffiths exercise 2.7]

- Examine the processes in Griffiths exercise 2.7 and state which one is possible or impossible, and why / with which interaction.

Hint: draw the corresponding Feynman diagrams if needed.

Exercise – 12: Penguins

- One of the flagship analyses of the LHCb experiment is the decay $B_s^0 \rightarrow J/\psi \phi$. It is used for the analysis of CP violation.
 - Draw the lowest order (tree) Feynman diagram of the process. ($B_s^0 = (\bar{b}s)$, $J/\psi = (c\bar{c})$, $\phi = (s\bar{s})$). Hint: it consists of an ‘internal’ $W \rightarrow cs$ transition. Why would this diagram be called ‘colour suppressed’?
 - In addition to the tree diagram, there is a famous ‘loop’ contribution, called a *penguin* diagram. Here, the b transforms into an s due to the emission and re-absorption of a W and a *quark*, while this *quark* radiates off a *gluon* that turns into a $c\bar{c}$. Draw this diagram. Try to twist and turn the diagram such that it may look as a penguin?
 - Ask for the funny story behind this name
 - Which flavour can this internal *quark* have? Which option has the largest probability?
 - The precision of LHCb is such that these penguin contributions are becoming a nuisance. Based in the vertices and CKM elements, can you guess how much weaker the penguin diagram is with respect to the tree?