Name (please print legibly!)

Gravitational Waves: Assignment 5 Astrophysical Sources

- 1. Practice with raising, lowering, and contracting indices: the metric tensor plays the role of allowing one to raise/lower indices on tensors, i.e. switching between contravariant and covariant forms.
 - (a) $g_{\mu\nu}A^{\nu} =$
 - (b) If ∂_{ν} is the contravariant gradient operator, $\eta^{\mu\nu}\partial_{\nu} =$
 - (c) $g^{ij}g_{jk} =$
 - (d) $g^{\alpha\gamma}g^{\beta\delta}A_{\gamma\delta} =$
 - (e) If \hat{n} is a unit vector, then $n^i n_i =$
 - (f) In spatial coordinates, $\delta_i^i =$
- 2. Let's show that $h_{ij}^{\text{TT}} = \Lambda_{ij,kl} \bar{h}^{kl}$ really is transverse and traceless. Recall that the lambda tensor is

$$\Lambda_{ij,kl}(\hat{n}) = P_{ik}P_{jl} - \frac{1}{2}P_{ij}P_{kl} \tag{1}$$

where $P_{ij}(\hat{n}) = \delta_{ij} - n_i n_j$ is a symmetric, transverse projection operator and \hat{n} is the direction of propagation of the GW.

- (a) Using results from Problem 1, show that $P_i^i = 2$.
- (b) Write out $\Lambda_{ij,kl}$ in terms of the Kroneker delta δ_{ij} and \hat{n} and apply the operator to \bar{h}^{kl} .
- (c) The condition for h_{ij}^{TT} to be transverse is that its component in the \hat{n} direction vanishes, i.e., $n^i h_{ij}^{\text{TT}} = 0$. Show that this is true.
- (d) The condition for h_{ij}^{TT} to be traceless is $\delta^{ij}h_{ij}^{\text{TT}} = 0$. Show that this is true.
- 3. When the direction of propagation \hat{n} is equal to \hat{z} , P_{ij} is the diagonal matrix

$$P = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix}.$$
 (2)

This means that P is a projector on the (x, y) plane. Then the quadrupole approximation reduces to

$$\begin{bmatrix} h_{ij}^{\mathrm{TT}} \end{bmatrix}_{\mathrm{quad}} = \frac{1}{r} \frac{2G}{c^4} \Lambda_{ij,kl}(\hat{n}) \ddot{M}^{kl}(t - r/c) \rightarrow \left(P \ddot{M} P \right)_{ij} - \frac{1}{2} P_{ij} \mathrm{Tr}(P \ddot{M}).$$

$$(3)$$

Show that this gives the following two polarization amplitudes for a gravitational wave propagating in the \hat{z} direction:

$$h_{+} = \frac{1}{r} \frac{G}{c^{4}} \left(\ddot{M}_{11} - \ddot{M}_{22} \right),$$

$$h_{\times} = \frac{2}{r} \frac{G}{c^{4}} \ddot{M}_{12}.$$
(4)

4. (a) Consider a binary on a fixed, circular Keplerian orbit. In the center-of-mass point of view for the system, suppose we have a single effective object of mass $\mu = (m_1 m_2)/(m_1 + m_2)$ described by coordinates:

$$x_{0}(t) = R \cos(\omega_{s}t + \pi/2),$$

$$y_{0}(t) = R \sin(\omega_{s}t + \pi/2),$$

$$z_{0}(t) = 0.$$
(5)

Determine the second derivatives of the second mass moments and plug these into the $h_+(t;\theta,\phi)$ and $h_{\times}(t;\theta,\phi)$ expressions for generic gravitational wave propagation that we saw in the lecture slides.

- (b) Consider the cases where the binary orbit is edge-on $(\theta = \iota = \pi/2)$ and face-on $(\theta = \iota = 0)$. What happens to the contributions from h_+ and h_{\times} in these cases?
- 5. Now consider a binary that loses energy due to the emission of gravitational waves. We saw in the lecture slides that the polarization amplitudes evolve as:

$$h_{+}(t) = \frac{1}{r} \left(\frac{GM_{c}}{c^{2}}\right)^{5/4} \left(\frac{5}{c\tau}\right)^{1/4} \frac{1 + \cos^{2}\iota}{2} \cos\left[\Phi(\tau)\right],$$

$$h_{\times}(t) = \frac{1}{r} \left(\frac{GM_{c}}{c^{2}}\right)^{5/4} \left(\frac{5}{c\tau}\right)^{1/4} \cos\iota\sin\left[\Phi(\tau)\right]$$
(6)

where $\Phi(\tau) = -2\left(\frac{5GM_c}{c^3}\right)^{-5/8} \tau^{5/8} + \Phi_0$, $\Phi_0 = \Phi(\tau = 0)$, and $\tau \equiv t_{\text{coal}} - t$. We will generate the plots of a few chirp waveforms using a starting frequency of $f_{\text{gw}}(\tau = t_{\text{coal}}) = 40$ Hz. You will need to invert the following expression to determine the value of $\tau = t_{\text{coal}}$ when $f_{\text{gw}} = 40$ Hz:

$$f_{\rm gw}(\tau) = \frac{1}{\pi} \left(\frac{5}{256\tau}\right)^{3/8} \left(\frac{GM_c}{c^3}\right)^{-5/8}.$$
 (7)

Assuming $\Phi_0 = 0$, $\iota = 0$ and r = 1 Mpc, create plots of h_+ and h_{\times} versus t for the following masses:

- (a) A binary neutron star: $m_1 = 1.4 M_{\odot}$ and $m_2 = 1.4 M_{\odot}$
- (b) A neutron-star-black-hole system: $m_1 = 1.4 M_{\odot}$ and $m_2 = 10 M_{\odot}$
- (c) A binary black hole: $m_1 = 10 M_{\odot}$ and $m_2 = 10 M_{\odot}$