

Loops and Mathematics

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- The First Loop
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- More Loops

The First Loop

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REGULARIZATION AND RENORMALIZATION OF GAUGE FIELDS

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Abstract: A new regularization and renormalization procedure is presented. It is particularly well suited for the treatment of gauge theories. The method works for theories that are known to be renormalizable as well as for Yang-Mills type theories. Overlapping divergencies are disentangled. The procedure respects unitarity, causality and allows shifts of integration variables. In non-anomalous cases also Ward identities are satisfied at all scales. It is transparent when anomalies, such as the Bell-Jackiw-Adler anomaly, may occur.

SCALAR ONE-LOOP INTEGRALS

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The completely general one-loop scalar one-, two-, three- and four-point functions are studied. Also an integral occurring in connection with soft bremsstrahlung is considered. Formulas in terms of Spence functions are given. An expansion for Spence functions with complex argument is presented. The whole forms a basis for the calculation of one-loop radiative corrections in the general case, including unstable particles and particles with spin.

⇒ establish a complete system of integrals: 2, 3, 4-point functions

$$\begin{aligned}
 B(k, m_1, m_2) &= \int \frac{d^n q}{(q^2 + m_1^2 - i0)((q+k)^2 + m_2^2 - i0)} \\
 &= -i\pi^2 \left[\frac{2}{n-4} + \ln(-k^2 - i0) \right] + \sum_j \left\{ \ln(1 - x_j) - x_j \ln \left(\frac{x_j - 1}{x_j} \right) - 1 \right\}
 \end{aligned}$$

The 3- and 4-point functions

$$C(p_j, m_i) = \int \frac{d^n q}{(q^2 + m_1^2)((q + p_1)^2 + m_2^2)((q + p_1 + p_2)^2 + m_3^2)}$$

$$D(p_j, m_i) = \int \frac{d^n q}{(q^2 + m_1^2)((q + p_1)^2 + m_2^2)((q + p_1 + p_2)^2 + m_3^2)((q + p_1 + p_2 + p_3)^2 + m_4^2)}$$

the following result obtains:

$$\begin{aligned} \frac{C}{i\pi^2} = & \int_0^1 dy \frac{1}{(c+2ab)y+d+e\alpha+2a+c\alpha} [\ln \{by^2+(c+e)y+a+d+f\} \\ & - \ln \{by_1^2+(c+e)y_1+a+d+f\}] \\ & - \int_0^1 dy \frac{1-\alpha}{(c+2ab)(1-\alpha)y+d+e\alpha} [\ln \{(a+b+c)y^2+(e+d)y+f\} \\ & - \ln \{(a+b+c)y_2^2+(e+d)y_2+f\}] \\ & - \int_0^1 dy \frac{\alpha}{-(c+2ab)\alpha y+d+e\alpha} [\ln \{ay^2+dy+f\} - \ln \{ay_3^2+dy_3+f\}]. \quad (5.6) \end{aligned}$$

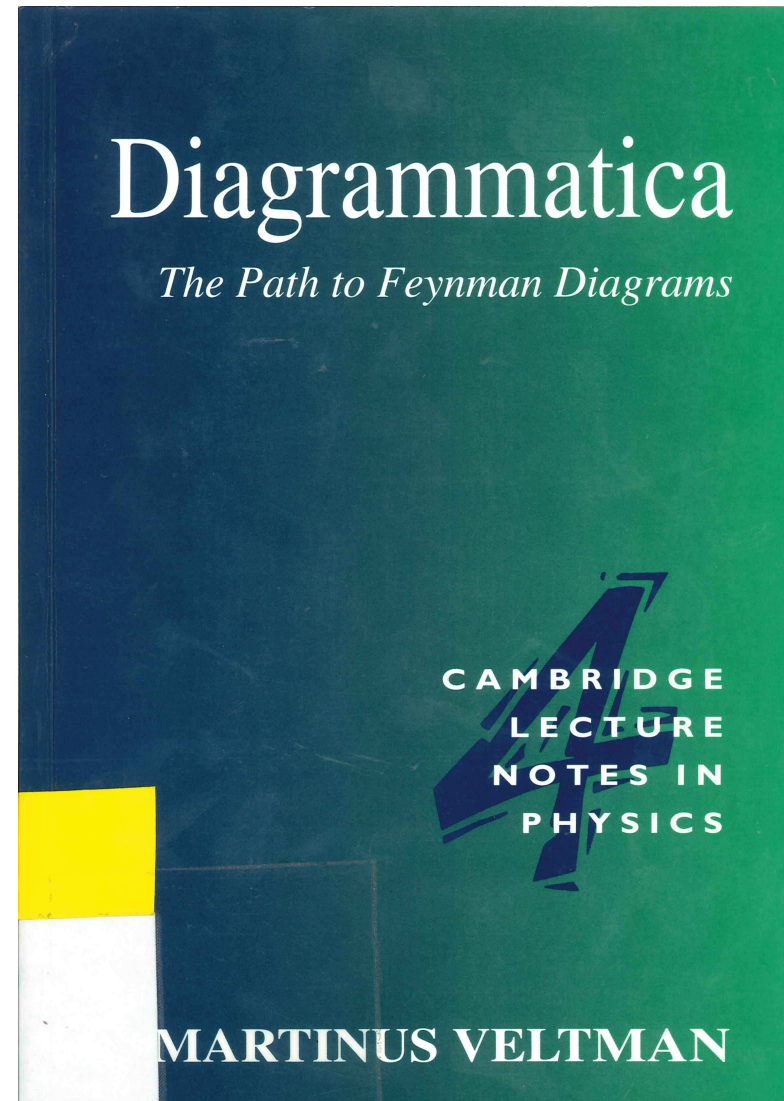
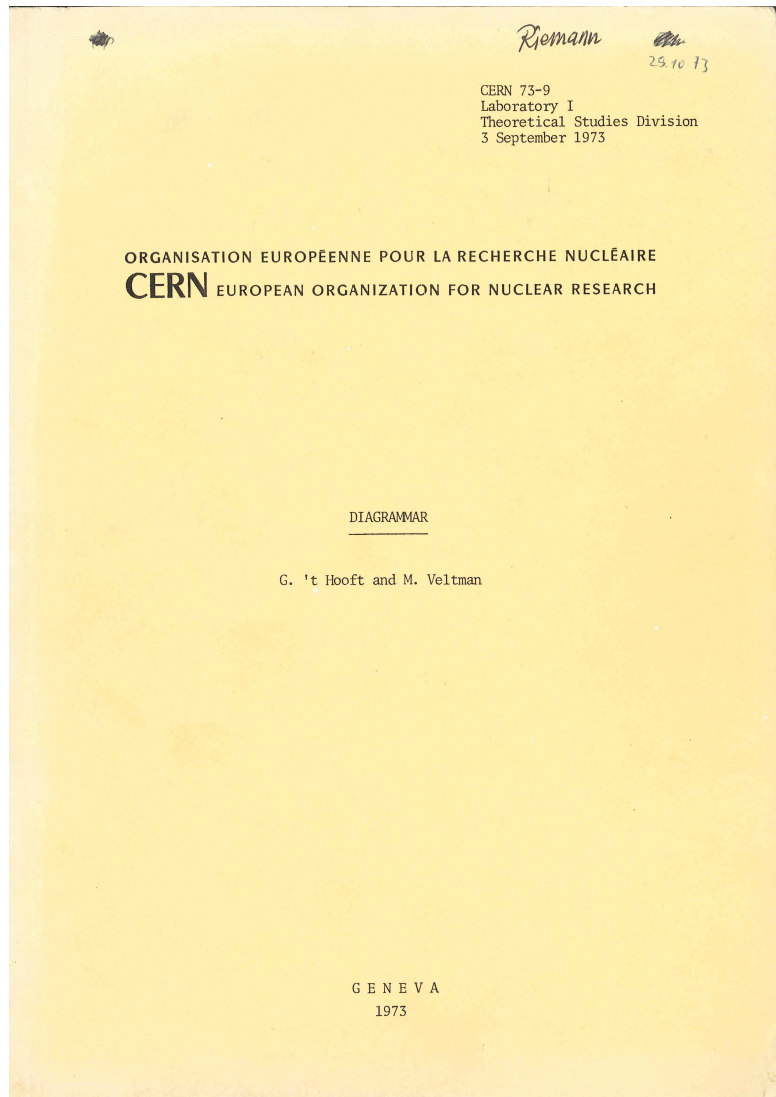
$$Li_2(z) = - \int_0^z \frac{dt}{t} \ln(1-t)$$

$$\begin{aligned} \frac{D}{i\pi^2} = & \frac{A_1 A_2 A_3 A_4}{k} \\ & \times \left[- \int_0^1 dy \frac{1-\alpha}{(c+2ab)(1-\alpha)y+d+e\alpha} \{\ln L_{24}(y) - \ln L_{24}(y_1)\} - R_{24}^1 \right. \\ & - \int_0^1 dy \frac{\alpha}{-(c+2ab)\alpha y+d+e\alpha} \{\ln L_{34}(y) - \ln L_{34}(y_2)\} + R_{34}^2 \\ & + \int_0^1 dy \frac{1}{(c+2ab)y+d+e\alpha+c\alpha+2a} \{\ln L_{23}(y) - \ln L_{23}(y_3)\} + R_{23}^3 \\ & + \int_0^1 dy \frac{1-\alpha}{(c+2ab)(1-\alpha)y+d+(e+k)\alpha} \{\ln L_{14}(y) - \ln L_{14}(y_4)\} + R_{14}^4 \\ & + \int_0^1 dy \frac{\alpha}{-(c+2ab)\alpha y+d+(e+k)\alpha} \{\ln L_{34}(y) - \ln L_{34}(y_5)\} - R_{34}^5 \\ & \left. - \int_0^1 dy \frac{1}{(c+2ab)y+d+(e+k)\alpha+c\alpha+2a} \{\ln L_{13}(y) - \ln L_{13}(y_6)\} - R_{13}^6 \right. \\ & \left. + \theta(-A_1 A_2) S \right], \quad (6.12) \end{aligned}$$

$$L_{ij}(y) = (-l_{ij}A_i A_j + m_i^2 A_i^2 + m_j^2 A_j^2)y^2 + (l_{ij}A_i A_j - 2m_j^2 A_j^2)y + m_j^2 A_j^2 - i\epsilon. \quad (6.13)$$

The One-Loop Problem: Solved Completely

The Basis for Many Calculations :



Integration = Anti-Differentiation

- Integration is easy if one knows the complete target space.
- Try to construct the **target space** for the Feynman integrals, forming certain algebras in general.

Important examples:

- Poincaré-iterated integrals

$$L(a_1, \dots, a_k; y) = \int_0^y \frac{dy_1}{y_1 - a_1} \int_0^{y_1} \frac{dy_2}{y_2 - a_2} \cdots \int_0^{y_{k-1}} \frac{dy_k}{y_k - a_k}$$

Work out all iterations for the alphabet $\mathfrak{A} = \{a_1, \dots, a_k\}$, $a_i \in \mathbb{R}$
e.g. $a_1 = 0, a_2 = 1, a_3 = -1, \dots$ **polylogs, Nielsen functions, HPL's, ...**

- Nested generalized harmonic sums

$$S_{n_1, \dots, n_k}(a_1, \dots, a_k; N) = \sum_{i_1=1}^N \frac{a_1^{i_1}}{i_1^{n_1}} \sum_{i_2=1}^{i_1} \frac{a_2^{i_2}}{i_2^{n_2}} \cdots \sum_{i_k=1}^{i_{k-1}} \frac{a_k^{i_k}}{i_k^{n_k}}$$

The representations are related by a Mellin transform

$$\mathbf{M}[f(x)](N) = \int_0^1 dx x^{N-1} [f(x)]_{(+)}$$

Both types of functions form (quasi) shuffle algebras and further obey structural relations.

$$\text{e.g. : } L(a_3; y)L(a_1, a_2; y) = L(a_3, a_1, a_2; y) + L(a_1, a_3, a_2; y) + L(a_1, a_2, a_3; y)$$

How do these structures emerge ?

Feynman parameter integrals massless or with one mass as regulator map into multi Mellin Barnes integrals

$$\frac{1}{(2\pi i)^l} \prod_l \int_{\gamma_l - i\infty}^{\gamma_l + i\infty} d\sigma_l \frac{\Gamma(a_1 N + b_1(\sigma_l) + r_1 \varepsilon) \dots \Gamma(a_k N + b_k(\sigma_l) + r_k \varepsilon)}{\Gamma(c_1 N + d_1(\sigma_l) + q_1 \varepsilon) \dots \Gamma(c_k N + d_k(\sigma_l) + q_k \varepsilon)}$$

These functions define (generalized) **hypergeometric** functions, Meijer G -functions, and generalizations thereof, which obey (multiple) integral and (multiple) sum representations.

The ε -**expansion** turns

$$\frac{\Gamma(z + r\varepsilon)}{\Gamma(z)} = 1 + \psi(z)\varepsilon r + \frac{1}{2} [\psi'(z) + \psi^2(z)] \varepsilon^2 r^2 \dots,$$

i.e. into product of single harmonic sums $S_k(z)$ and ζ -values.

The multiple sums, e.g. via z , introduce **nesting** which can be solved by modern summation methods **SIGMA**, C. Schneider 2000–

Two Loops

- Very many important calculations all over the world; notably in
The Netherlands: LEP, HERA, pp-colliders ... 

Veltman, Berends, Gastmans, Gaemers, Laenen, van Neerven, Smith, Vermaseren, van der Bij, ... and many other groups world wide, 1980 –

$${}_pF_q(a_k; b_l; x) = \sum_{n=0}^{\infty} \frac{(a_1)_n \dots (a_p)_n}{(1)_n (b_1)_n \dots (b_q)_n} x^n$$

More Loops

... necessarily means less external Legs: **currently two.**

- Massless and massive results up to $O(\alpha_s^5)$

- $g-2$: $O(\alpha^3)$ Laporta, Remiddi, 1996

- Running α_s : $O(\alpha_s^4)$  Larin, van Ritbergen, Vermaseren 1997; Czakon, 2005

- Unpol. anomalous dimensions and Wilson coefficients: $O(\alpha_s^3)$

Moch, Vermaseren, Vogt 2004/05 

- Unpol. NS anomalous dimension 2nd Moment: $O(\alpha_s^4)$ Baikov, Chetyrkin 2006

- Moments of DIS Heavy Flavor Wilson Coefficients for $Q^2 \gg m^2$: $O(\alpha_s^3)$:

Bierenbaum, Blümlein, Klein, 2009

- Contributions to the vacuum polarization : $O(\alpha_s^5)$ Baikov, Chetyrkin, Kühn 2003

- $R(s)$, Adler function, ρ -parameter, Z - and τ -decay ...: $O(\alpha_s^4)$

Baikov, Chetyrkin, Kühn, last 10 years

- Zero scale quantities: up to $O(\alpha_s^{4(5)})$

- Single scale quantities: $O(\alpha_s^3)$

More Loops

Massless(ive) Zero Scale Quantities : $\ln(2)$; ζ_2 ; ζ_3 ; $\text{Li}_4(1/2)$; $(\zeta_5, \text{Li}_5(1/2))$; $(\text{Li}_6(1/2), \sigma_{-5,-1})$; $(\zeta_7, \text{Li}_7(1/2), \sigma_{-5,1,1}, \sigma_{5,-1,-1})$; ... \uparrow 3 loops Vermaseren 1998

Massless Single Scale Quantities : $S_1(N), \dots, S_{-2,2,1}(N), \dots$
 \uparrow 3 loops, weight 5

Massive Single Scale Quantities : in addition ..., $S_{-2,2,1}(\frac{1}{2}, 1, -2; N), \dots$
 \uparrow 3 loops, weight 5

Massive Zero Scale Quantities (more massive lines): (3- and 4 loop)
cyclotomic zeta values, elliptic functions Broadhurst 1998, Ablinger, JB, Schneider, 2011, Laporta 2008

Systematic Integration

Consider graphs with no poles in $1/\varepsilon$. F. Brown, 2008

Zero Scale Graphs can be integrated directly, if all Feynman parameters enter linearly; first problems at 6 loops in scalar field theory.

More Loops

Topics of contemporary research include:

- To which structures does a single mass lead (fixed and variable moments) ?
- Which combinatoric and summation problems arise ?
- What is needed additionally to treat the large amount of diagrams with poles ?
- At which loop level do new structures occur in QED and QCD ?

Symbioses with Mathematics:

- Advanced discrete algorithms and combinatorics \longrightarrow computer algebra
- Complex analysis of new higher transcendental functions
- Theory of irrational numbers, motivic numbers, periods, algebraic geometry

Advanced Technologies to Evaluate Feynman Diagrams

Some Examples:

- 🕒 **Zero-scale Problems** : Euler-Zagier and Multiple Zeta Values
JB, D. Broadhurst, J. Vermaseren, Comput.Phys.Commun. 181 (2010) 582
find all relations : \implies **Tera-Terms** to be processed
alternating: all relations up to $w = 12$ (6-loop level);
non-alternating: all relations up to $w = 22$; determined.
Interesting relations: to $w = 30$;
- 🕒 **Reconstructing recurrent quantities** from Mellin Moments
JB, M. Kauers, S. Klein, C. Schneider, Comput.Phys.Commun. 180 (2009) 2143
Can one find the anomalous dimensions and Wilson coefficients to 3-loops just from their moments ? **Yes** - recurrent quantities in Mellin space.
 ≤ 5114 Moments; difference equation fills 440 books
Complete computation: 5 CPU Months

Tini at Loops and Legs



1996: before the prize
Talk about the physics in momentum space.
Veltman prize \Rightarrow Ettore Remiddi.

Special Talk: 2000



1999



... after the prize



We owe Tini :

- Various key precision predictions
- Many essential technologies in physics
- The example of a straight character
- Reminding us, what is really essential.

Happy Birthday, Tini! - And many Happy Returns.