

## SIGHTSEEING

 IN THE LANDSCAPE
## CONTENTS

- Landscape remarks (physics/06041340, Dutch version 1998)
- RCFT orientifolds (with Huiszoon,Fuchs, Schweigert,Walcher)
- 2003-2004 results
(with Dijkstra,Huiszoon)
Q 2005-2006 results
(with Anastasopoulos,Dijkstra,Kiritsis, hep-th/0605226)


## UNIFICATION / UNIQUENESS A BRIEF HISTORY

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- Muon (quark/lepton families)
- Parameters (masses, coupings)

Q Then some theoretical problems arose: Yang-Mills theory: QED is not unique. Many other gauge theories are possible.

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Q But there is another revolution most people preferred to overlook: The string vacuum revolution.

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## A CERN CAFETARIA NAPKIN (~ 1988)

All gauge theories

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Unique vacuum


Complexity
All gauge theories
Life
Intelligence

## A CERN CAFETARIA NAPKIN (~ 1988)



## Naar een waardig slot

Bert Schellekens

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Q 2003: "The Anthropic Landscape of String Theory" (L. Susskind)

## THE ANTHROPIC PRINCIPLE

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Q Is an inevitable consequence of String Theory.

## ANTHROPIC PRINCIPLES

Q The SM gauge theory is not the only solution.
Q Many others do not allow "life".
9 There should be enough to understand why ours exist.

Q Within anthropic regions, we can determine parameters using probabilities.

## HOW MANY "VACUA" ARE NEEDED?

Q Requires understanding of "anthropic" considerations for different gauge theories.

Q Requires some definition of a measure and boundaries.

Wild guess: about $10^{20}$ for SM fine-tunings

The same problems exist in principle for the cosmological constant, but seem less serious there: about $10^{120}$ would be needed. Recent estimates: String Theory has plenty of ground states to understand all fine-tunings.
(Bousso-Polchinski, Douglas Deneff,...

## SUMMARY:

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9 Therefore: A Success for String Theory
Q4-D Quantum gravity implies that the SM is part of a huge landscape: an amazing conclusion! (if correct).

Q Fits nicely with some of the great discoveries in the history of science (heliocentric model, theory of Evolution...)

## Demystification by huge numbers:

Q Planets (Giordano Bruno)

- Mutations (Evolution)

Q Universes (Eternal Inflation)
Q Alternative "Standard Models" (The Landcape)

## Demystification by huge numbers:

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Q Universes (Eternal Inflation)
Q Alternative "Standard Models" (The Landcape)

## A repetion of an old mistake:

There is nothing "special" about us.
This line of thought fits in very well with a series of insights that pointed out our modest place in the cosmos. Our planet is not the center of the solar system, our sun is just one of many stars and not even a very special one, and the same is true for our galaxy. It seems natural to assume that also our universe, including the quarks, leptons and interactions we observe is just one out of many possibilities.
(From physics/06041340)

Q String Theory has never looked better...

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Q ... but it has never looked harder.

## REASONABLE GOALS

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Q Explore unknown regions of the landscape

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Q Explore unknown regions of the landscape
Q Establish the likelyhood of standard model features (gauge group, three families, ....)

Q Convince ourselves that the standard model is a plausible vacuum.

Q Determine if we are the "Chinese" or the "Andorrans" of the landscape.

Q ... and maybe we get lucky


## ORIENTIFOLDS <br> OF <br> GEPNER MODELS

## EARLIER FOOTPRINTS

C. Angelantonj, M. Bianchi, G. Pradisi, A. Sagnotti and Y. S. Stanev, Phys. Lett. B 387 (1996) 743 [arXiv:hep-th/9607229].
R. Blumenhagen and A. Wisskirchen, Phys. Lett. B 438, 52 (1998)
[arXiv:hep-th/9806131].
G. Aldazabal, E. C. Andres, M. Leston and C. Nunez, JHEP 0309, 067 (2003) [arXiv:hep-th/0307183].
I. Brunner, K. Hori, K. Hosomichi and J. Walcher, arXiv:hep-th/0401137.
R. Blumenhagen and T. Weigand, JHEP 0402 (2004) 041 [arXiv:hep-th/0401148].
G. Aldazabal, E. C. Andres and J. E. Juknevich, JHEP 0405, 054 (2004) [arXiv:hep-th/0403262].

## THE LONG ROAD TO THE CHIRAL SSM

(0. Angelantonj, Bianchi, Pradisi, Sagnotti, Stanev (1996)

Chiral spectra from Orbifold-Orientifoldos

* Aldazabal, Franco, Ibanez, Rabadan, Uranga (2000)

Blumenhagen,Görlich,Körs,Lüst (2000)
Ibanez, Marchesano, Rabadan (2001)
Non-supersymmetric SM-Spectra with RR tadpole cancellation

- Cvetic, Shiu, Uranga (2001)

Supersymmetric SM-Spectra with chiral exotics

* Blumenhagen, Görlich, Ott (2002)

Honecker (2003)
Supersymmetric Pati-Salam Spectra with brane recombination
. Dijkstra, Huiszoon, Schellekens (2004)
Supersymmetric Standard Model (Gepner Orientifolds)

* Honecker, Ott (2004)

Supersymmetric Standard Model (Zoorbifoldorientifold)

## CLOSED STRING PARTITION FUNCTION



$$
P(\tau, \bar{\tau})=\sum_{i j} \chi_{i}(\tau) Z_{i j} \chi_{j}(\bar{\tau})
$$

Type IIB

## Orientifold Partition Functions

## Orientifold Partition Functions



## ORIENTIFOLD PARTITION FUNCTIONS



## ORIENTIFOLD PARTITION FUNCTIONS



## Transverse Channel



## GEPNER MODELS

Building Blocks:
Minimal $\mathrm{N}=2 \mathrm{CFT}$

$$
c=\frac{3 k}{k+2}, \quad k=1, \ldots, \infty
$$

168 ways of solving

$$
\sum_{i} c_{k_{i}}=9
$$

Spectrum:

$$
\begin{gathered}
h_{l, m}=\frac{l(l+2)-m^{2}}{4(k+2)}+\frac{s^{2}}{8} \\
(l=0, \ldots k ; \quad q=-k, \ldots k+2 ; \quad s=-1,0,1,2) \\
\quad \text { (plus field identification) }
\end{gathered}
$$

$4(k+2)$ simple currents

## TENSORING

賕 Preserve world－sheet susy
榉 Preserve space－time susy（GSO）
箓 Use surviving simple currents to build MIPFs

蝶 This yields one point in the moduli space of a Calabi－Yau manifold

## Selecting MIPFs And Orientifolds

Each tensor product has a discrete group $\mathcal{G}$ of simple currents：$J \cdot a=b$

Choose：
$\int$ 並 A subgroup $\mathcal{H}$ of $\mathcal{G}$
諩 A rational matrix $X_{\alpha \beta}$ defined on $\mathcal{H}$
$\int$ 絜 An element $K$ of $\mathcal{G}$
曗 A set of signs $\beta_{K}(J)$ defined on $\mathcal{H}$

## A MIPF

$$
\begin{gathered}
\quad(0+2)^{\wedge} 2+(1+3)^{\wedge} 2+(4+6)^{*}(13+15)+(5+7)^{*}(12+14) \\
+(8+10)^{\wedge} 2+(9+11)^{\wedge} 2+(12+14)^{*}(5+7)+(13+15)^{*}(4+6) \\
+(16+18)^{*}(25+27)+(17+19)^{*}(24+26)+(20+22)^{\wedge} 2+(21+23)^{\wedge} 2 \\
+(24+26)^{*}(17+19)+(25+27) *(16+18)+(28+30)^{\wedge} 2+(29+31)^{\wedge} 2 \\
+(32+34)^{\wedge} 2+(33+35)^{\wedge} 2+(36+38)^{*}(45+47)+(37+39)^{*}(44+46) \\
+(40+42)^{\wedge} 2+(41+43)^{\wedge} 2+(44+46)^{*}(37+39)+(45+47)^{*}(36+38) \\
+(48+50) *(57+59)+(49+51)^{*}(56+58)+(52+54)^{\wedge} 2+(53+55)^{\wedge} 2 \\
+(56+58) *(49+51)+(57+59) *(48+50)+(60+62)^{\wedge} 2+(61+63)^{\wedge} 2
\end{gathered}
$$

$$
\begin{aligned}
& +2 \text { * } 2913 \text { ) }{ }^{*}(2915)+2^{*}(2914) *(2912)+2^{*}(2915) *(2913) \\
& +2^{*}(2916)^{\wedge} 2+2^{*}(2917)^{\wedge} 2+2^{*}(2918)^{\wedge} 2+2^{*}(2919)^{\wedge} 2 \\
& +2^{*}(2920)^{\wedge} 2+2^{*}(2921)^{\wedge} 2+2^{*}(2922)^{\wedge} 2+2^{*}(2923)^{\wedge} 2 \\
& +2^{*}(2924) *(2926)+2 *(2925) *(2927)+2 *(2926) *(2924) \\
& +2 \text { * } 2927 \text { )*(2925) }+2^{* *}(2928)^{\wedge} 2+2 *(2929)^{\wedge} 2+2 *(2930)^{\wedge} 2 \\
& +2 *(2931)^{\wedge} 2+2 *(2932) *(2934)+2^{*}(2933) *(2935) \\
& +2 *(2934) *(2932)+2 *(2935) *(2933)+2 *(2936) *(2938) \\
& +2 \text { * } 2937 \text { ) }{ }^{*}(2939)+2^{*}(2938) *(2936)+2 *(2939) *(2937) \\
& +2{ }^{*}(2940)^{\wedge} 2+2 *(2941)^{\wedge} 2+2^{*}(2942)^{\wedge} 2+2 *(2943)^{\wedge} 2
\end{aligned}
$$

## BOUNDARIES AND CROSSCAPS*

## 諩 Boundary coefficients

$$
R_{\left[a, \psi_{a}\right](m, J)}=\sqrt{\frac{|\mathcal{H}|}{\left|\mathcal{C}_{a}\right|\left|\mathcal{S}_{a}\right|}} \psi_{a}^{*}(J) S_{a m}^{J}
$$

粼 Crosscap coefficients

$$
U_{(m, J)}=\frac{1}{\sqrt{|\mathcal{H}|}} \sum_{L \in \mathcal{H}} e^{\pi i\left(h_{K}-h_{K L}\right)} \beta_{K}(L) P_{L K, m} \delta_{J, 0}
$$

*Huiszoon, Fuchs, Schellekens, Schweigert, Walcher (2000)

## COEFFICIENTS

籟 Klein bottle

$$
K^{i}=\sum_{m, J, J^{\prime}} \frac{S^{i}{ }_{m} U_{(m, J)} g_{J, J^{\prime}}^{\Omega, m} U_{\left(m, J^{\prime}\right)}}{S_{0 m}}
$$

䋛 Annulus

$$
A_{\left[a, \psi_{a}\right]\left[b, \psi_{b}\right]}^{i}=\sum_{m, J, J^{\prime}} \frac{S^{i}{ }_{m} R_{\left[a, \psi_{a}\right](m, J)} g_{J, J^{\prime}}^{\Omega, m} R_{\left[b, \psi_{b}\right]\left(m, J^{\prime}\right)}}{S_{0 m}}
$$

䗱 Moebius

$$
M_{\left[a, \psi_{a}\right]}^{i}=\sum_{m, J, J^{\prime}} \frac{P^{i}{ }_{m} R_{\left[a, \psi_{a}\right](m, J)} g_{J, J^{\prime}}^{\Omega, m} U_{\left(m, J^{\prime}\right)}}{S_{0 m}}
$$

$g_{J, J^{\prime}}^{\Omega, m}=\frac{S_{m 0}}{S_{m K}} \beta_{K}(J) \delta_{J^{\prime}, J^{c}}$

## PARTITION FUNCTIONS

## 蟔 Closed

$$
\frac{1}{2}\left[\sum_{i j} \chi_{i}(\tau) Z_{i j} \chi_{i}(\bar{\tau})+\sum_{i} K_{i} \chi_{i}(2 \tau)\right]
$$

諩 Open

$$
\frac{1}{2}\left[\sum_{i, a, n} N_{a} N_{b} A_{a b}^{i} \chi_{i}\left(\frac{\tau}{2}\right)+\sum_{i, a} N_{a} M_{a}^{i} \hat{\chi}_{i}\left(\frac{\tau}{2}+\frac{1}{2}\right)\right]
$$

$N_{a}$ : Chan-Paton multiplicity

## TADPOLES \＆ANOMALIES

絆 Tadpole cancellation condition：

$$
\sum_{b} N_{b} R_{b(m, J)}=4 \eta_{m} U_{(m, J)}
$$

期 Cubic $\operatorname{Tr} F^{3}$ anomalies cancel

暽 Remaining anomalies by Green－Schwarz mechanism

溸 In rare cases，additional conditions for global anomaly cancellation＊

## Abelian Masses

Green-Schwarz mechanism


Axion-Vector boson vertex
-------MWW

Generates mass vector bosons of anomalous symmetries

$$
(e . g . B+L)
$$

But may also generate mass for non-anomalous ones

$$
(Y, B-L)
$$

## SCOPE OF THE SEARCH

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## 靿 168 Gepner models

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Essential to decide what to search for！

## WHAT TO SEARCH FOR

## The Madrid model



Chiral $\operatorname{SU}(3) \times \mathrm{SU}(2) \times \mathrm{U}(1)$ spectrum:

$$
3(u, d)_{L}+3 u_{L}^{c}+3 d_{L}^{c}+3\left(e^{-}, \nu\right)_{L}+3 e_{L}^{+}
$$

Y massless

$$
Y=\frac{1}{6} Q_{a}-\frac{1}{2} Q_{c}-\frac{1}{2} Q d
$$

$\mathrm{N}=1$ Supersymmetry
No tadpoles, global anomalies

## THE HIDDEN SECTOR



## REQUIRED SPECTRUM

## 3 families of $\mathrm{SU}(3) \times \mathrm{SU}(2) \times \mathrm{U}(1)$

+ non-chiral matter


## STATISTICS

| Total number of 4-stack configurations | 45761187347637742772 <br> $\left(45.7 \times 10^{18}\right)$ |
| :--- | :--- |
| Total number scanned | $4.37522 \mathrm{E}+19$ |
| Total number of SM configurations | 45051902 <br> fraction: $1.0 \times 10^{-12}$ |
| Total number of tadpole solutions | 1649642 <br> fraction: $3.8 \times 10^{-14}\left(^{*}\right)$ |
| Total number of distinct solutions | 211634 |

(*) cf. Gmeiner, Blumenhagen,Honecker,Lüst,Weigand: "One in a Billion"

Standard model type: 6

## Number of factors in hidden gauge group: 0

 Gauge group: $\mathrm{U}(3) \mathrm{x} \operatorname{Sp}(2) \mathrm{x} \mathrm{U}(1) \mathrm{x} \mathrm{U}(1)$Number of representations: 19

| 3 | x | (V , V | , 0,0 ) | chirality 3 |
| :---: | :---: | :---: | :---: | :---: |
| 3 | x | (V , 0 | , V , 0 ) | chirality -3 |
| 3 | x | (V , 0 | , V*, 0 ) | chirality -3 |
| 9 | X | ( 0 , V | , 0 , V ) | chirality 3 |
| 5 | x | (0,0 | , V , V ) | chirality -3 |
| 3 | x | ( 0,0 | , V , V*) | chirality -3 |
| 2 | x | (V , 0 | , 0 , V ) |  |
| 10 | x | ( 0 , V | , V , 0 ) |  |
| 2 | x | (Ad, 0 | , 0,0 ) |  |
| 2 | x | ( $\mathrm{A}, 0$ | , 0,0 ) |  |

Higgs: $(2,1 / 2)+2 *, 1 / 2)$


$$
\sin ^{2}\left(\theta_{w}\right)=.5271853
$$

$\frac{\alpha_{3}}{\alpha_{2}}=3.2320501$


## Require only：

谱 $\mathrm{U}(3)$ from a single brane

敖 $\mathrm{U}(2)$ from a single brane

瑔 Quarks and leptons，Y from at most four branes

数 $\mathrm{G}_{\mathrm{CP}} \supset \mathrm{SU}(3) \times \mathrm{SU}(2) \times \mathrm{U}(1)$
＊Chiral $G_{C P}$ fermions reduce to quarks，leptons （plus non－chiral particles）but

龉 No fractionally charged mirror pairs

䗲 Massless Y
（＊）P．Anastasopoulos，T．Dükstra，E．Kiritsis，A．Schellekens

## AlLowed Features

䗲（Anti）－quarks from anti－symmetric tensors

楼 leptons from anti－symmetric tensors

紸 family symmetries

塐 non－standard Y－charge assignments

䩚 Unification（Pati－Salam，（flipped）SU（5），trinification）＊

䅉 Baryon and／or lepton number violation

瞵 ．．．．
＊a，b，c，d may be identical

Chan-Paton gauge group

$$
G_{C P}=U(3)_{a} \times\left\{\begin{array}{c}
U(2)_{b} \\
S p(2)_{b}
\end{array}\right\} \times G_{c} \quad\left(\times G_{d}\right)
$$

Embedding of Y:

$$
Y=\alpha Q_{a}+\beta Q_{b}+\gamma Q_{c}+\delta Q_{d}+W_{c}+W_{d}
$$

Q: Brane charges (for unitary branes)
W: Traceless generators

## CLASSIFICATION

$$
Y=\left(x-\frac{1}{3}\right) Q_{a}+\left(x-\frac{1}{2}\right) Q_{b}+x \underbrace{Q_{C}+(x-1)} Q_{D}
$$

## Distributed over c and d

Allowed values for $x$
1/2 Madrid model, Pati-Salam, Flipped SU(5)
0 (broken) SU(5)
1 Antoniadis, Kiritsis, Tomaras
$-1 / 2,3 / 2$
any Trinification $(x=1 / 3) \quad$ (orientable)

## THE BASIC ORIENTABLE MODEL

$$
\begin{align*}
& U(3) \times U(2) \times U(1) \times U(1) \\
& 3 \times\left(V, V^{*}, 0,0\right) \\
& 3 \times\left(V^{*}, 0, V, 0\right) \\
& \text { (u,d) } \\
& 3 \times\left(V^{*}, 0,0, V\right) \\
& 6 \times\left(0, V, V^{*}, 0\right) \\
& \left(\mathrm{e}^{-}, \nu\right)+\mathrm{H}_{1} \\
& 3 \times\left(0, V, 0, V^{*}\right)  \tag{2}\\
& 3 \times\left(0,0, V, V^{*}\right) \\
& \mathrm{e}^{+}
\end{align*}
$$

"D-branes at singularities"

## RESULTS

龉 Searched all MIPFs with＜ 1750 boundaries （4557 of 5403 MIPFs）

稢 19345 chirally different SM embeddings found
諩 Tadpole conditions solved in 1900 cases
（18＂old＂ones）

## StATISTICS

| Value of x | Total |
| :---: | :---: |
| 0 | 21303612 |
| $1 / 2$ | $124006839^{*}$ |
| 1 | 12912 |
| $-1 / 2,3 / 2$ | 0 |
| any | 1250080 |

*Previous search: 45051902

## UNIFICATION

(1)

## SU(5) MODELS



## SU(5)



Note: gauge group is just $\operatorname{SU}(5)$ !

## SUMMARY

Examples exist of chiral orientifold SSM spectra exist

Q Without mirrors
Q Without adjoints
Q Without (anti)-symmetric tensors
Q Without Observable-Hidden matter
Q Without hidden sector

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Q Without mirrors
Q Without adjoints
Q Without (anti)-symmetric tensors
Q Without Observable-Hidden matter
Q Without hidden sector
....but to get all this simultaneously requires more statistics


IT'S JUST ONE SMALL STEP:
874 HODGE NUMBERS SCANNED
AT LEAST 30000 KNOWN (M. KREUZER)

