

# ORIENTIFOLD GUTS

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(WITH E. KIRITSIS AND M. LENNEK)



“GUTS IN STRINGS”  
DESY, 5 FEBRUARI 2009



**ALTERNATIVE TITLE:**



# GUTS IN STRINGS



# GUTS VS. STRINGS



# GUTs: SELLING POINTS

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- Beautiful!
- Unifies 3 of the 4 known interactions
- Explains family structure
- Explains charge quantization
- Predicts  $\sin^2 \theta_w$
- Baryogenesis?



# AROUND 1983:

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Experimental confirmation seemed imminent:

**First Results from a Superconductive Detector for Moving Magnetic Monopoles**

Blas Cabrera

*Physics Department, Stanford University, Stanford, California 94305*

(Received 5 April 1982)

Proton decay experiments were starting...

The expectations were reminiscent of those regarding  
SUSY at the LHC



*H. Georgi,  
Fourth workshop on Grand Unification  
Philadelphia, 1983*



GUTs IN STRINGS?



# GUTS IN STRINGS?

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- Beauty: implausible selection criterion in the landscape.
- Unification of interactions: string theory not only unifies the three gauge interactions, but also gravity, without any need for GUTs.



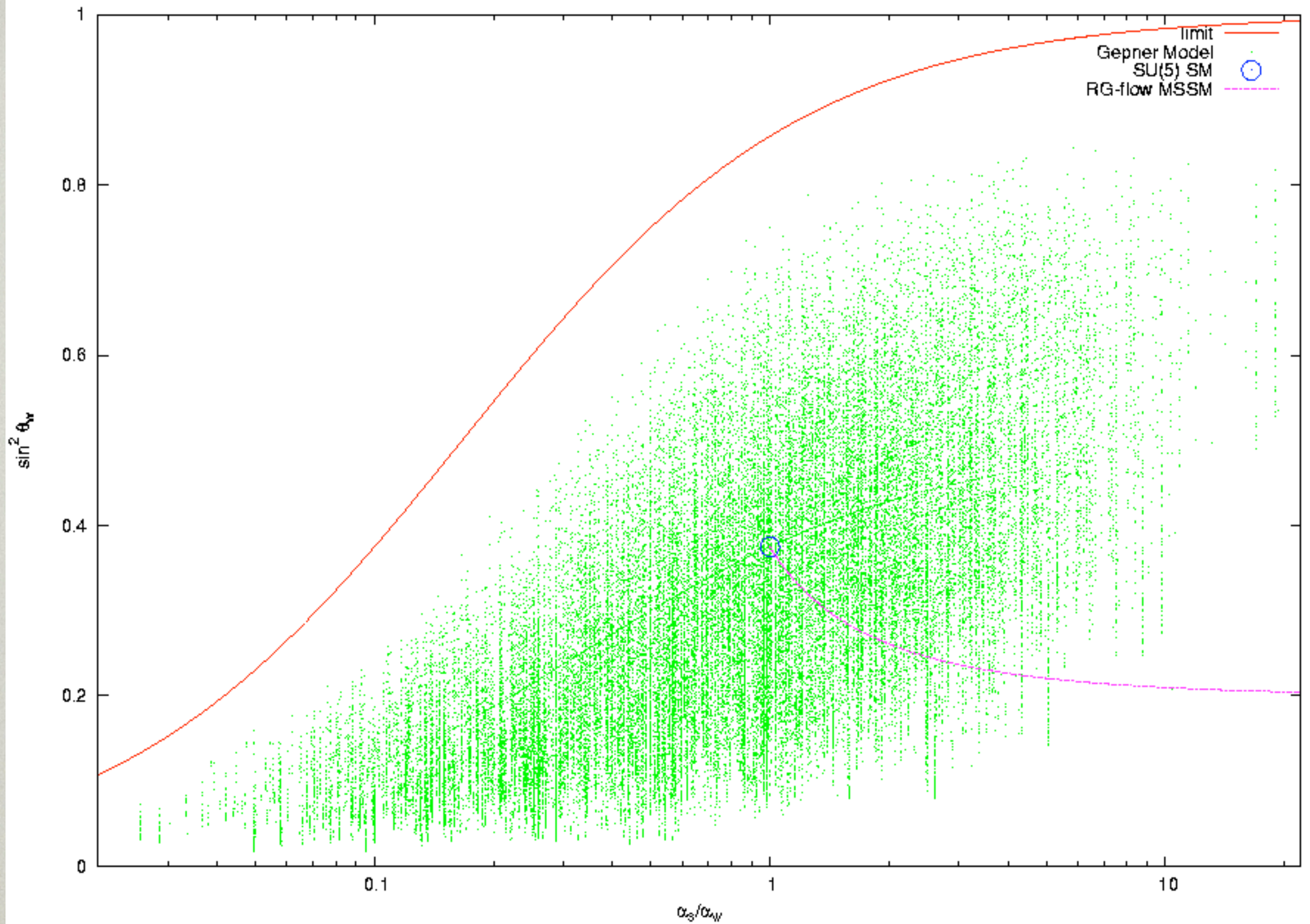
# GUTS IN STRINGS?

COUPLING CONSTANT  
UNIFICATION









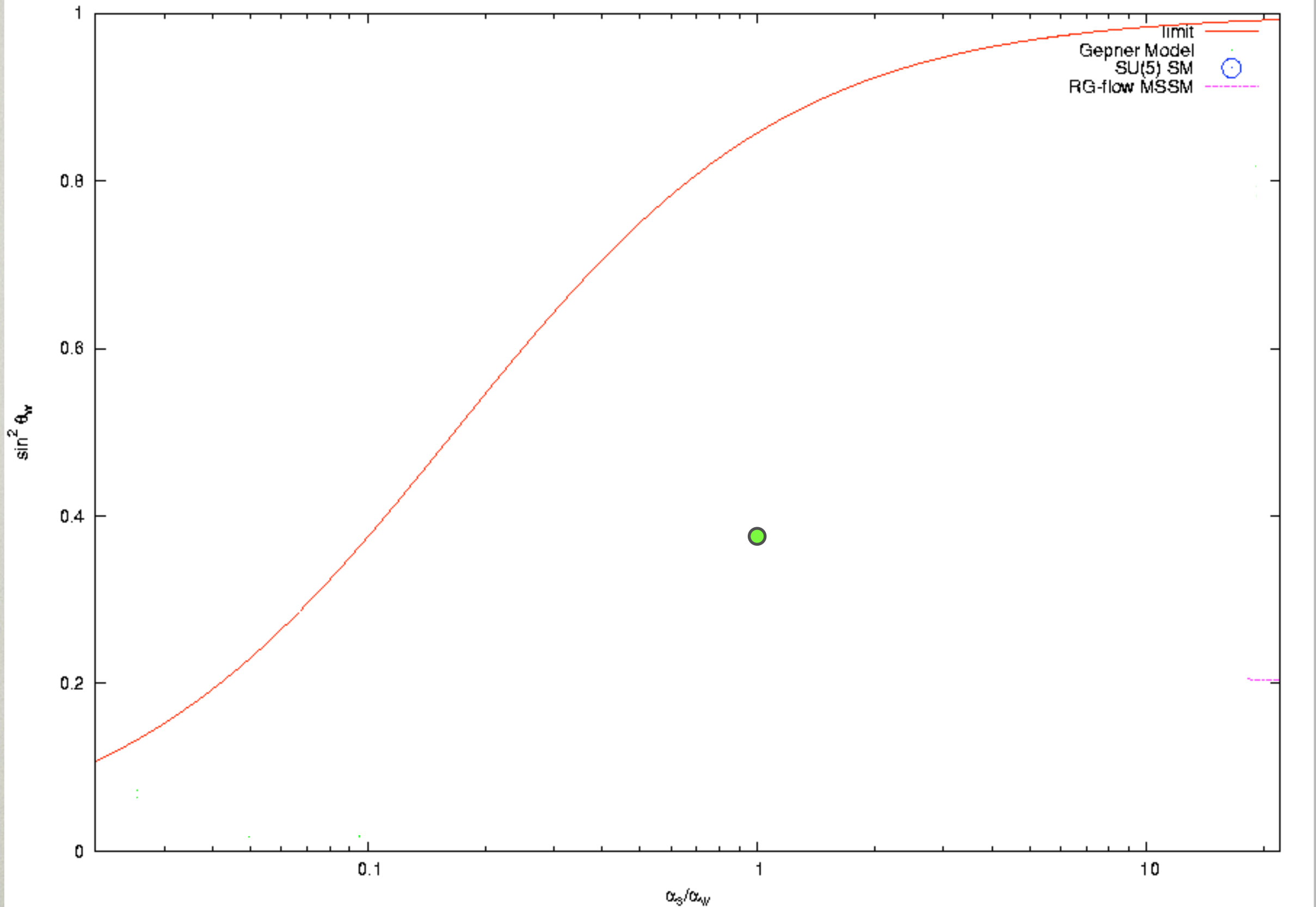
Dijkstra, Huiszoon, Schellekens, *Nucl.Phys.B710:3-57,2005*



We are using orientifolds or heterotic strings to get some idea about generic features of the landscape.

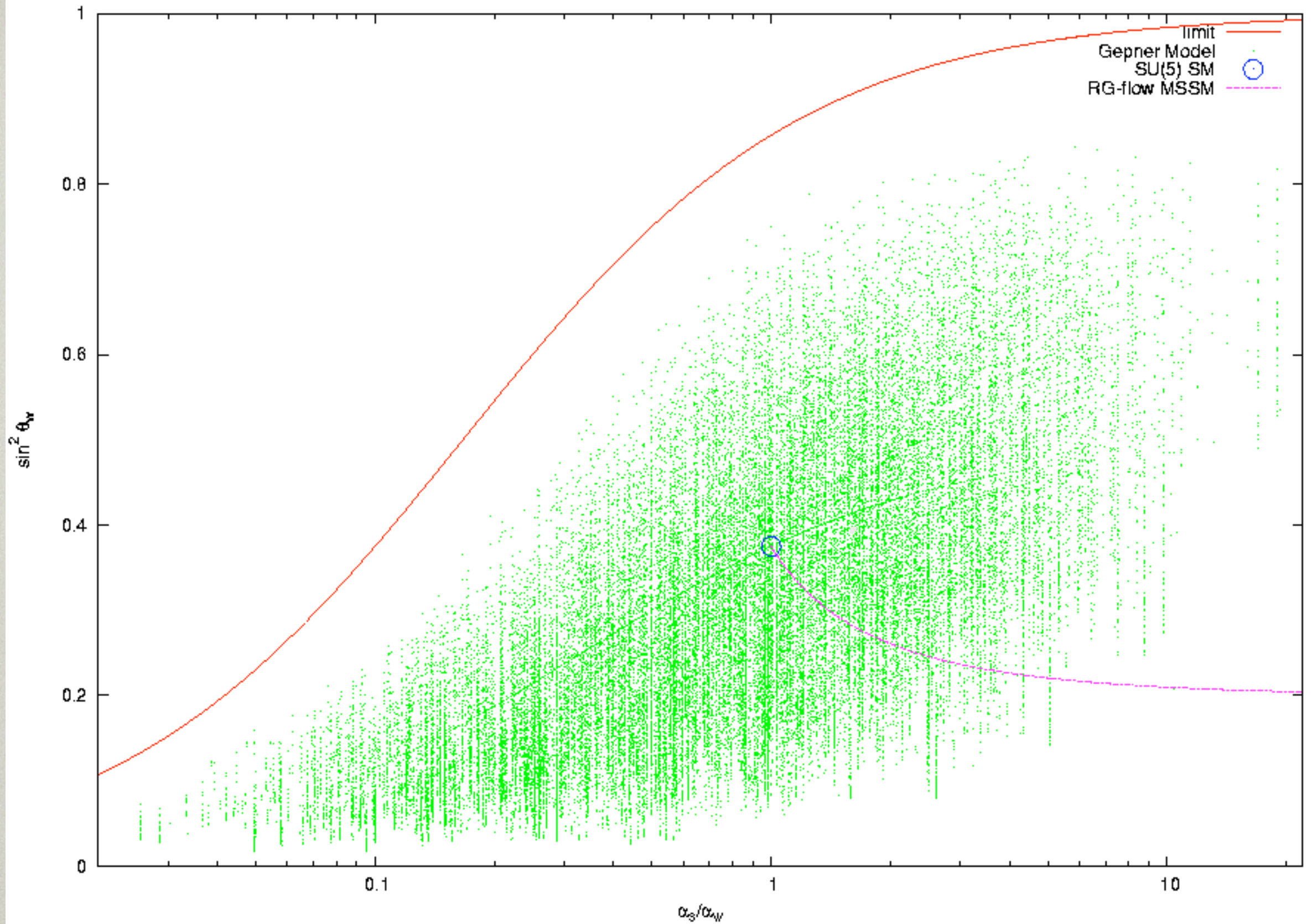
So what looks more generic:





# Heterotic Strings





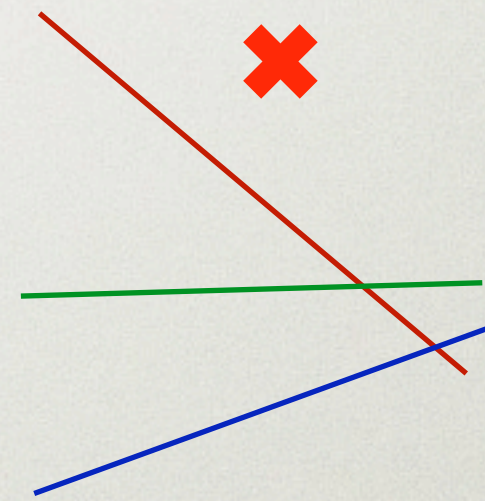
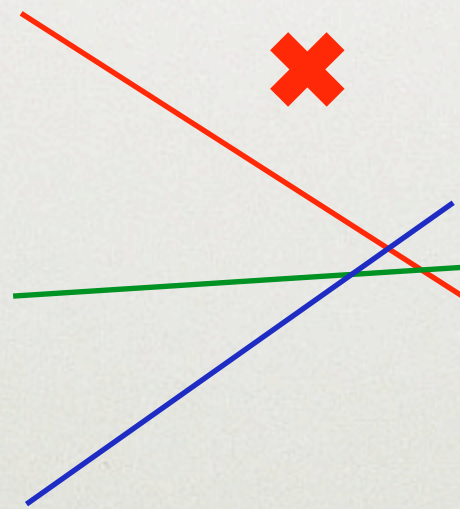
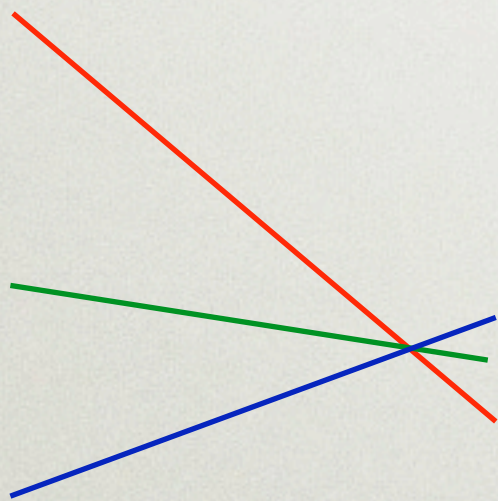
## Orientifolds



Evidence for GUTs (without strings)  
But not predicted by strings  
(but can be “accomodated”)



Better options:





# GUTS IN STRINGS?

## FAMILY STRUCTURE



GUTs emerge “naturally” in  
compactified  $E_8 \times E_8$  Heterotic strings:

- Embedding of the spin-connection in the gauge group  
(CHSW, 1984)
- “Bosonic string map”  
(LLS, 1986)



# BOSONIC STRING MAP\*

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One-to-one map of the characters of the 4-dimensional covariant NSR model to the characters of  $SO(10) \times E_8$  (affine level 1)

This fact can be used to build heterotic string partition functions starting from a diagonal bosonic or type-II partition function:

$$\sum_i \bar{\chi}_i^{\text{NSR}}(\bar{\tau}) \chi_i^{SO(10) \times E_8}(\tau) \times \dots$$

Automatically yields chiral spectra with a number of (16)'s of  $SO(10)$ .

This fact was exploited by Gepner in 1987.

(\* *Lerche, Lüst, Schellekens (1986)*)



# BOSONIC STRING MAP

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Using a simple current extension one can get  $E_6 \supset SO(10)$

By orbifoldings one can get subgroups of  $SO(10)$

There is a large network\* of related string theories for which the structure of one family can be traced back to the characters of the 4-D NSR models:

$$D = 4 \rightarrow SO(10) \text{ - like family structure}$$

The existence of a GUT group  $SO(10)$  at any scale is not really required.

(\*) Includes the “mini-landscape”? (Talks by Ratz, Nilles, Schmidt, Schmidt-Hoberg, Ramos)



$D = 4 \rightarrow SO(10)$  - like family structure

A triumph for (heterotic) string theory!

Conceptually this is far superior to field theory GUTs:

- Dictates the choice of  $SO(10)$  (or  $SU(5)$ ,  $E_6$ ) over most other Lie algebras.
- Dictates the choice of  $(16)$ 's of  $SO(10)$  over any other anomaly free representation.
- Anomaly cancellation is not an ad-hoc constraint as it is in QFT



Even if this does not work, one may appeal to the slightly less powerful statement that the (16) of  $SO(10)$  (and its branchings) is among the few chiral representations allowed at affine level 1.

Even better:

$$SU(3)_1 \times SU(2)_1 \times U(1)_{30}$$

(the standard  $U(1)$  normalization for coupling constant convergence)  
has a simple current extension to  $SU(5)$ .

This means that any such SM realization is an orbifold of an  $SU(5)$ .



# But what about other parts of the landscape?

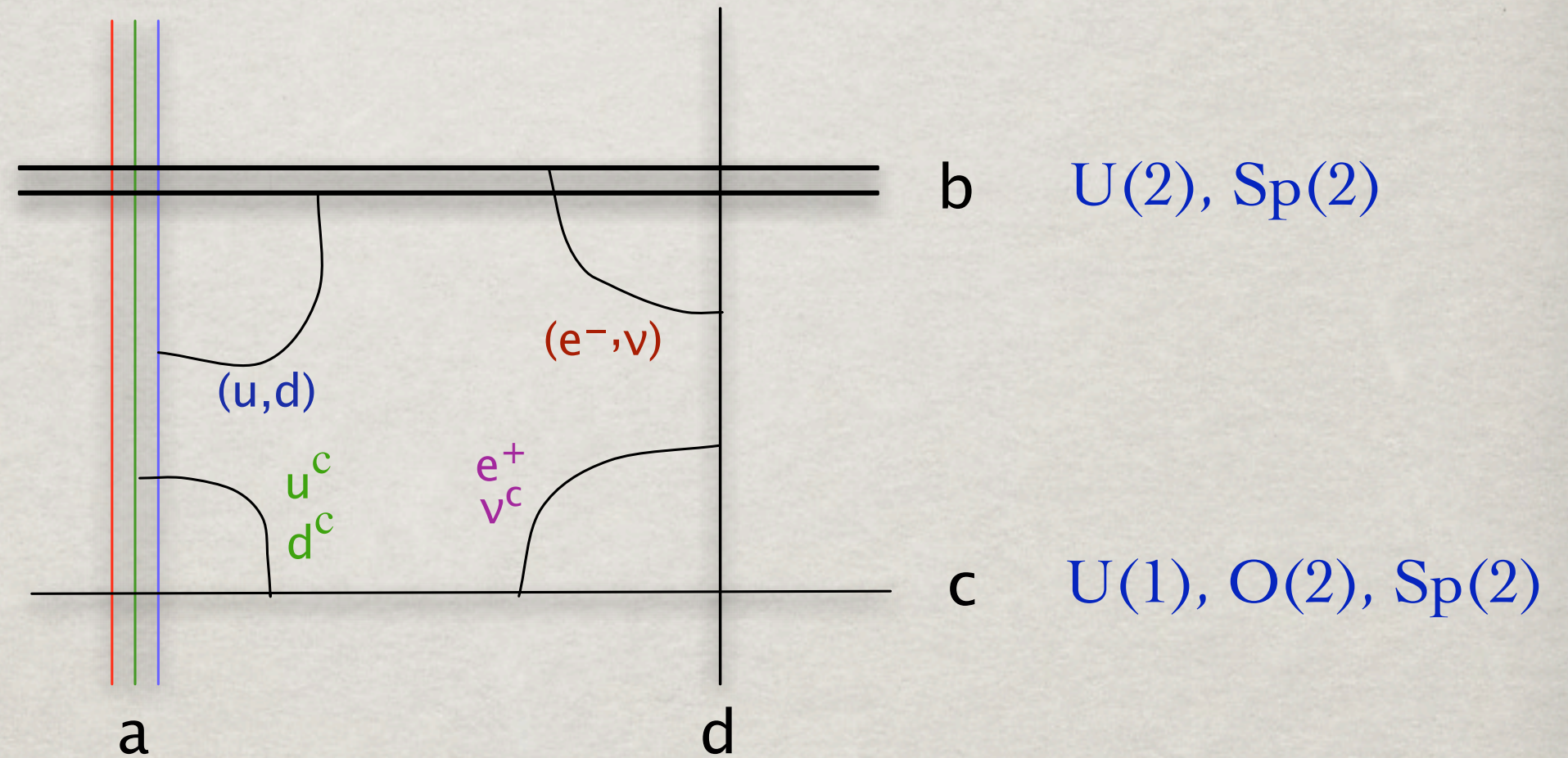
The best studied alternative are orientifold / intersecting brane models.

(F-theory GUTs (*Vafa, Heckman*) assume GUTs *ab initio*, so there is nothing to discuss.)

Orientifolds can produce the right family structure without any apparent relation to  $SO(10)$ :



# THE MADRID MODEL\*



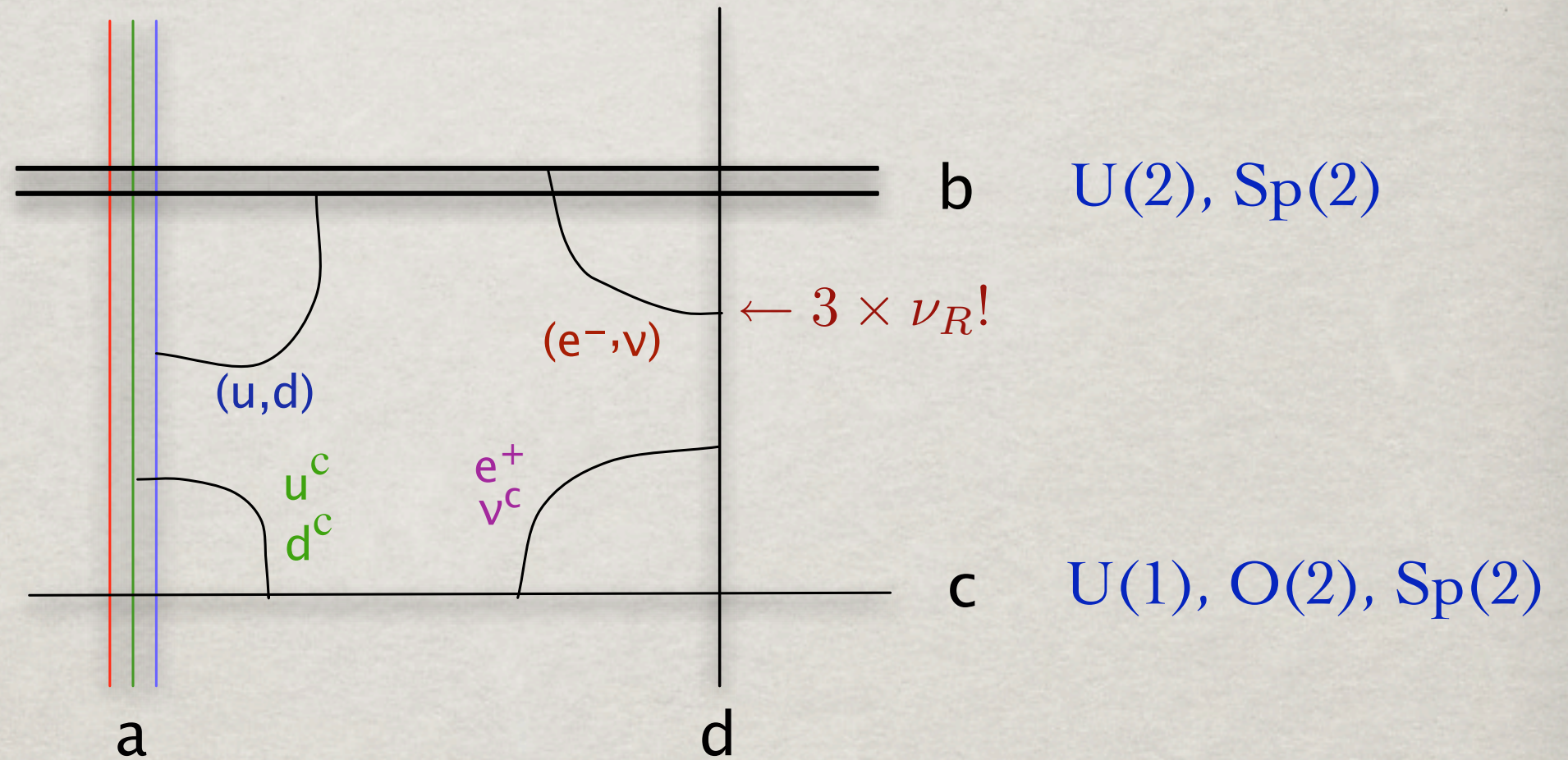
$$Y = \frac{1}{6}Q_a - \frac{1}{2}Q_c - \frac{1}{2}Q_d$$

(\* ) Ibanez, Marchesano, Rabadan (2001)

Lots of further work, reviewed in: Blumenhagen, Cvetic, Langacker, Shiu (2005)



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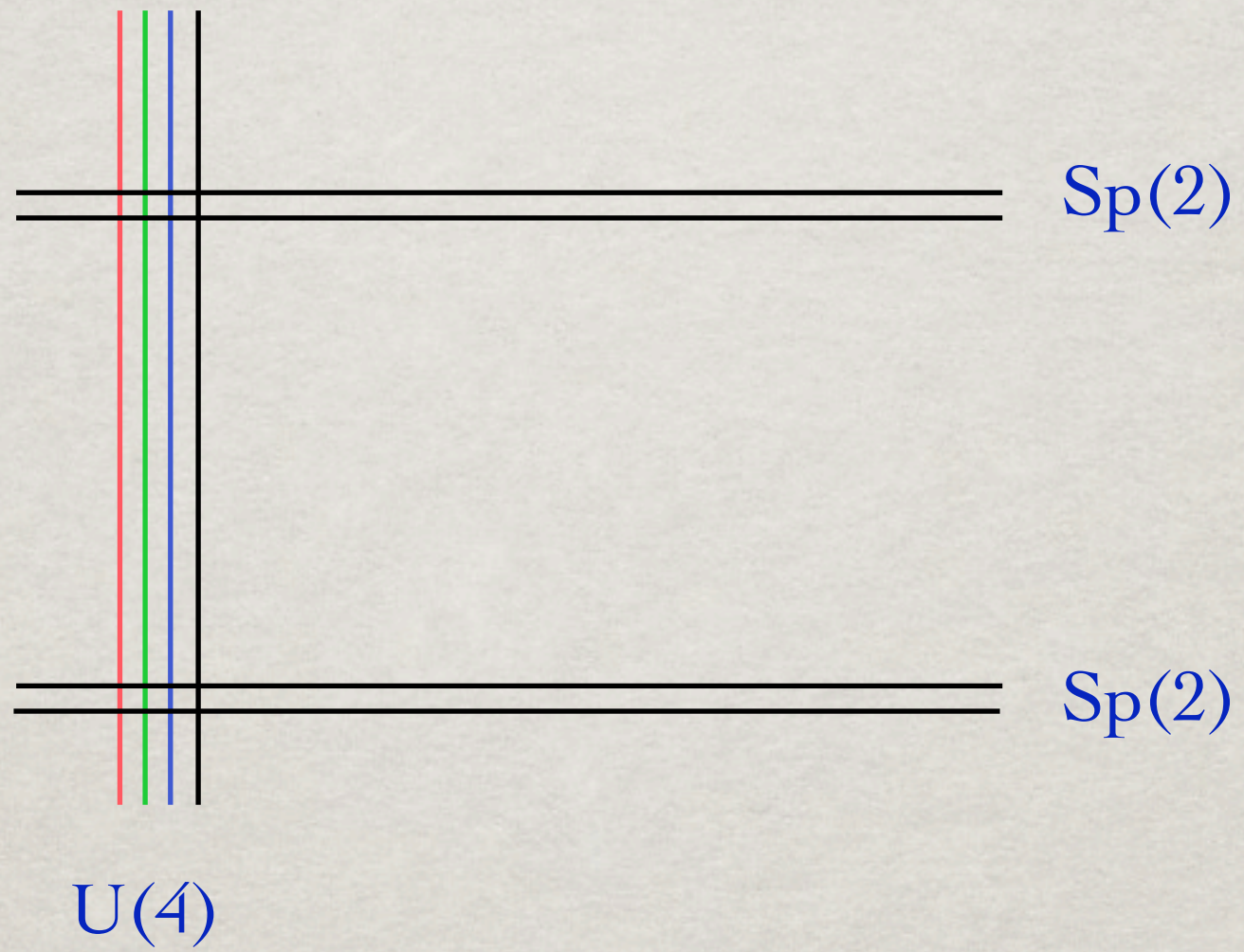
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Pati-Salam model  $\subset$   $SO(10)$





Pati-Salam model  $\subset$   $SO(10)$



OTHER ORIENTIFOLD  
REALIZATIONS OF THE  
STANDARD MODEL

WITH ANASTASOPOULOS, DIJKSTRA, KIRITSIS



# SEARCH CRITERIA

Require only:

- $U(3)$  from a single brane
- $U(2)$  from a single brane
- Quarks and leptons,  $Y$  from at most four branes
- $G_{\text{CP}} \supset SU(3) \times SU(2) \times U(1)$
- Chiral  $G_{\text{CP}}$  fermions reduce to quarks, leptons (plus non-chiral particles)
- Massless  $Y$



# CHAN-PATON GROUP

$$G_{CP} = U(3)_a \times \left\{ \begin{array}{l} U(2)_b \\ Sp(2)_b \end{array} \right\} \times G_c \quad (\times G_d)$$

Embedding of Y:

$$Y = \alpha Q_a + \beta Q_b + \gamma Q_c + \delta Q_d + W_c + W_d$$

Q: Brane charges (for unitary branes)

W: Traceless generators



# CLASSIFICATION

$$Y = \left(x - \frac{1}{3}\right)Q_a + \left(x - \frac{1}{2}\right)Q_b + \underbrace{xQ_c + (x - 1)Q_d}_{\text{Distributed over c and d}}$$

Distributed over  
c and d

Allowed values for  $x$

$1/2$	Madrid model, Pati-Salam, Flipped SU(5)
$0$	(broken) SU(5)
$1$	Antoniadis, Kiritsis, Tomaras model
$-1/2, 3/2$	
any	Trinification ( $x = 1/3$ ) (orientable)



We looked for these configurations in the context of orientifolds of Gepner models.

The “branes” are realized as boundary states of the CFT.

We found a total of more than 19000\* chirally distinct standard model realizations.

(\* ) No chiral exotics, modulo non-chiral exotics



# STATISTICS

(Before tadpole cancellation)

Value of $x$	Total
0	24483441
$1/2$	138837612
1	30580
$-1/2, 3/2$	0
any	1250080



- The dominant classes  $x=0, 1/2$  contain Pati-Salam and SU(5) models  
In this sense one could argue that the family structure is GUT-related, without having a GUT.
- SU(5) GUT models are about .01% of the total, and about .6% of the  $x=0$  models.



# MOST FREQUENT MODELS

nr	Total occ.	MIPFs	Chan-Paton Group	spectrum	x	Solved
1	9801844	648	$U(3) \times Sp(2) \times Sp(6) \times U(1)$	VVVV	1/2	Y!
2	8479808(16227372)	675	$U(3) \times Sp(2) \times Sp(2) \times U(1)$	VVVV	1/2	Y!
3	5775296	821	$U(4) \times Sp(2) \times Sp(6)$	VVV	1/2	Y!
4	4810698	868	$U(4) \times Sp(2) \times Sp(2)$	VVV	1/2	Y!
5	4751603	554	$U(3) \times Sp(2) \times O(6) \times U(1)$	VVVV	1/2	Y!
6	4584392	751	$U(4) \times Sp(2) \times O(6)$	VVV	1/2	Y
7	4509752(9474494)	513	$U(3) \times Sp(2) \times O(2) \times U(1)$	VVVV	1/2	Y!
8	3744864	690	$U(4) \times Sp(2) \times O(2)$	VVV	1/2	Y!
9	3606292	467	$U(3) \times Sp(2) \times Sp(6) \times U(3)$	VVVV	1/2	Y
10	3093933	623	$U(6) \times Sp(2) \times Sp(6)$	VVV	1/2	Y
11	2717632	461	$U(3) \times Sp(2) \times Sp(2) \times U(3)$	VVVV	1/2	Y!
12	2384626	560	$U(6) \times Sp(2) \times O(6)$	VVV	1/2	Y
13	2253928	669	$U(6) \times Sp(2) \times Sp(2)$	VVV	1/2	Y!
14	1803909	519	$U(6) \times Sp(2) \times O(2)$	VVV	1/2	Y!
15	1676493	517	$U(8) \times Sp(2) \times Sp(6)$	VVV	1/2	Y
16	1674416	384	$U(3) \times Sp(2) \times O(6) \times U(3)$	VVVV	1/2	Y
17	1654086	340	$U(3) \times Sp(2) \times U(3) \times U(1)$	VVVV	1/2	Y
18	1654086	340	$U(3) \times Sp(2) \times U(3) \times U(1)$	VVVV	1/2	Y
19	1642669	360	$U(3) \times Sp(2) \times Sp(6) \times U(5)$	VVVV	1/2	Y
20	1486664	346	$U(3) \times Sp(2) \times O(2) \times U(3)$	VVVV	1/2	Y!
21	1323363	476	$U(8) \times Sp(2) \times O(6)$	VVV	1/2	Y
22	1135702	350	$U(3) \times Sp(2) \times Sp(2) \times U(5)$	VVVV	1/2	Y!
23	1050764	532	$U(8) \times Sp(2) \times Sp(2)$	VVV	1/2	Y
24	956980	421	$U(8) \times Sp(2) \times O(2)$	VVV	1/2	Y
25	950003	449	$U(10) \times Sp(2) \times Sp(6)$	VVV	1/2	Y
26	910132	51	$U(3) \times U(2) \times Sp(2) \times O(1)$	AAVV	0	Y
...						



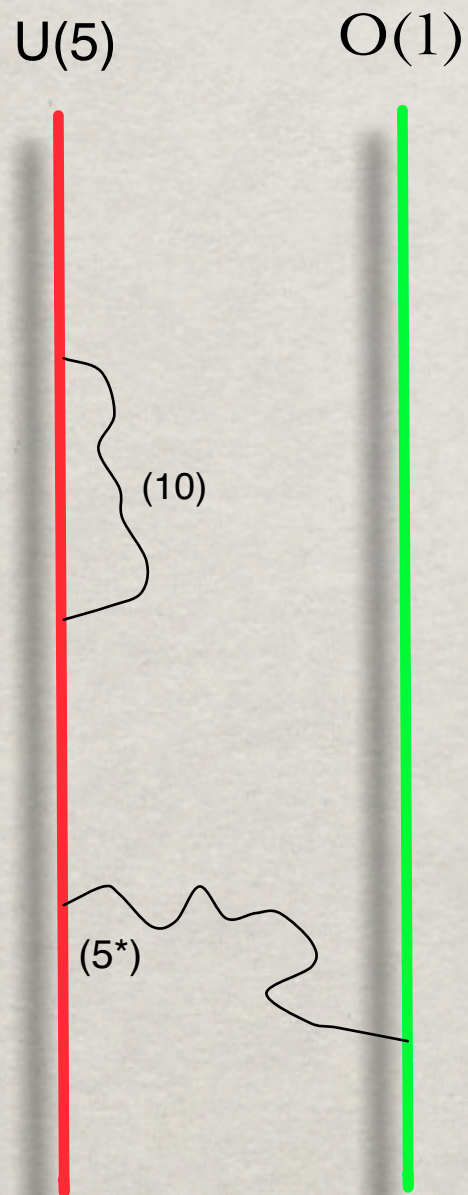
# CURIOSITIES

nr	Total occ.	MIPFs	Chan-Paton Group	Spectrum	x	Solved
617	16845	296	$U(5) \times O(1)$	AV	0	Y
671	14744(*)	29	$U(3) \times U(2) \times U(1) \times U(1)$	VVVV	1/2	
761	12067	26	$U(3) \times U(2) \times U(1)$	AAS	1/2	Y!
762	12067	26	$U(3) \times U(2) \times U(1)$	AAS	0	Y!
1024	7466	7	$U(3) \times U(2) \times U(2) \times U(1)$	VAAV	1	
1125	6432	87	$U(3) \times U(3) \times U(3)$	VVV	*	Y
1201	5764(*)	20	$U(3) \times U(2) \times U(1) \times U(1)$	VVVV	1/2	
1356	5856(*)	10	$U(3) \times U(2) \times U(1) \times U(1)$	VVVV	1/2	Y
1725	2864	14	$U(3) \times U(2) \times U(1) \times U(1)$	VVVV	1/2	Y
1886	2381	115	$U(6) \times Sp(2)$	AV	1/2	Y!
1887	2381	115	$U(6) \times Sp(2)$	AV	0	Y!
1888	2381	115	$U(6) \times Sp(2)$	AV	1/2	Y!
2624	1248	3	$U(3) \times U(2) \times U(2) \times U(3)$	VAAV	1	
2753	1136	74	$U(5) \times U(1)$	AS	0	Y
2880	1049	34	$U(5) \times U(1)$	AS	1/2	Y!
2881	1049	34	$U(5) \times U(1)$	AS	0	Y!
6580	146	18	$U(5) \times U(1)$	AS	0	
14861	12	2	$U(5) \times U(1)$	AS	0	



# MODEL NR. 617

*Gauge group is just SU(5)!*



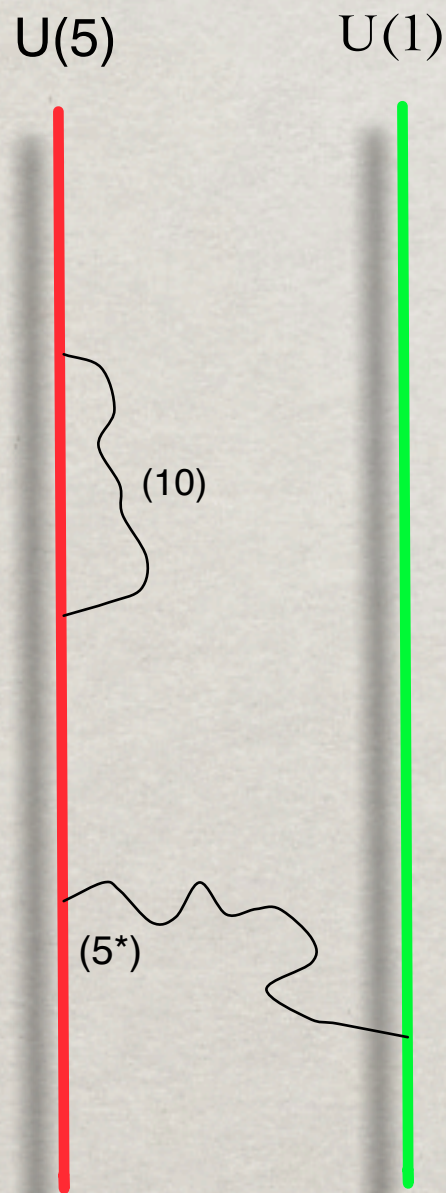
	U5	O1	O1	
3 x	(A ,0	,0	,0	) chirality 3
11 x	(V ,V	,0	,0	) chirality -3
8 x	(S ,0	,0	,0	)
3 x	(Ad ,0	,0	,0	)
1 x	(0 ,A	,0	,0	)
3 x	(0 ,V	,V	,V	)
8 x	(V ,0	,V	,V	)
2 x	(0 ,S	,0	,0	)
4 x	(0 ,0	,S	,S	)
4 x	(0 ,0	,A	,A	)

Hidden sector



# MODEL NR. 2880

Gauge group is  $SU(5) \times U(1)$



**U5 U1**

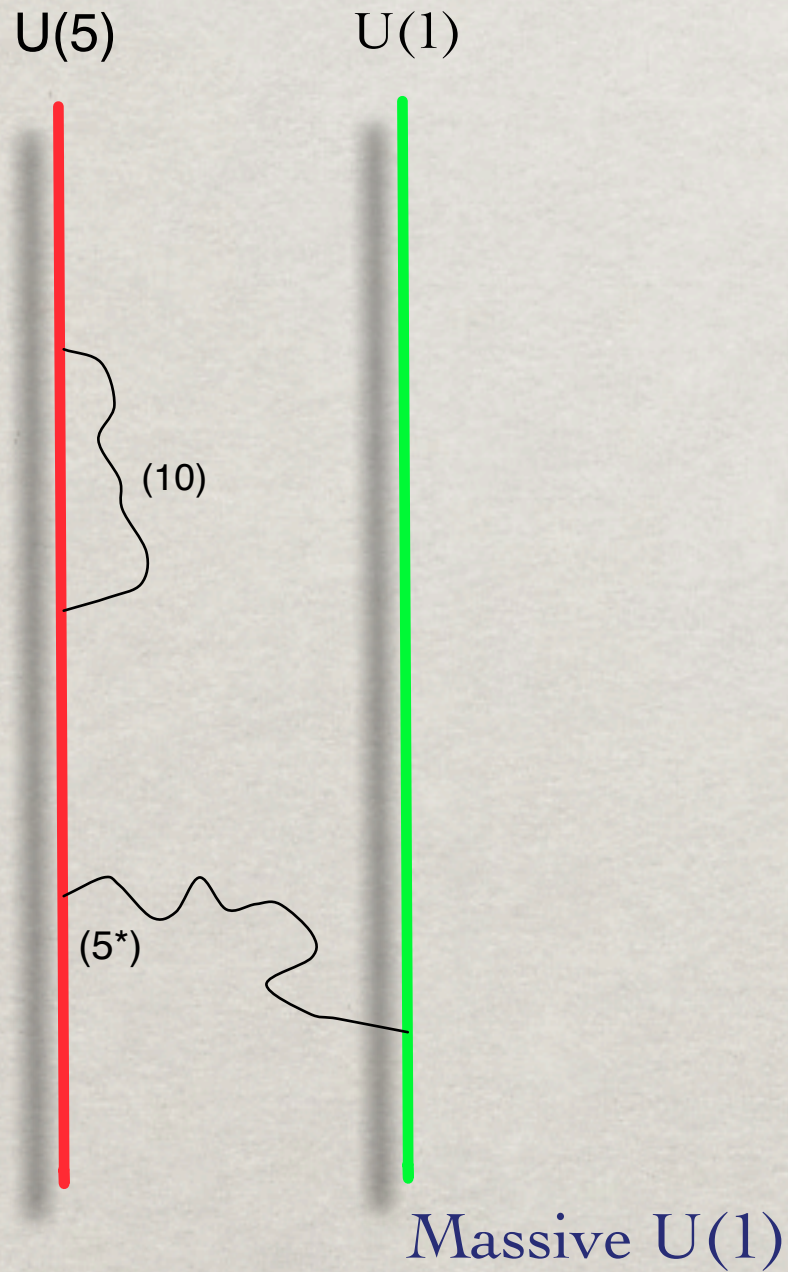
- 11 x ( 0 ,S ) chirality 3
- 3 x ( A ,0 ) chirality 3
- 5 x ( V ,V ) chirality -3
- 8 x ( S ,0 ) chirality 0
- 9 x ( Ad,0 ) chirality 0
- 5 x ( 0 ,Ad) chirality 0
- 4 x ( 0 ,A ) chirality 0
- 12 x ( V ,V\*) chirality 0

Massless U(1)  
Allows flipped SU(5)

No hidden sector!



# MODEL NR. 2753



U5 U1 O2 U2 O2 U5 S4 U1 U1

7 x ( 0 , S , 0 , 0 , 0 , 0 , 0 , 0 , 0 )	chirality 3
3 x ( A , 0 , 0 , 0 , 0 , 0 , 0 , 0 , 0 )	chirality 3
3 x ( V , V , 0 , 0 , 0 , 0 , 0 , 0 , 0 )	chirality -3
2 x ( S , 0 , 0 , 0 , 0 , 0 , 0 , 0 , 0 )	chirality 0
2 x ( 0 , A , 0 , 0 , 0 , 0 , 0 , 0 , 0 )	chirality 0
3 x ( Ad , 0 , 0 , 0 , 0 , 0 , 0 , 0 , 0 )	chirality 0
2 x ( 0 , Ac , 0 , 0 , 0 , 0 , 0 , 0 , 0 )	chirality 0
4 x ( 0 , V , 0 , 0 , V , 0 , 0 , 0 , 0 )	chirality 0
2 x ( V , 0 , 0 , 0 , 0 , 0 , 0 , 0 , V* )	chirality 0
4 x ( V , 0 , 0 , 0 , 0 , 0 , 0 , 0 , V )	chirality 0
2 x ( V , 0 , 0 , 0 , 0 , V , 0 , 0 , 0 )	chirality 0
2 x ( V , 0 , 0 , 0 , 0 , 0 , V , 0 , 0 )	chirality 0
2 x ( V , 0 , 0 , 0 , 0 , 0 , 0 , V* , 0 )	chirality 0
2 x ( V , V* , 0 , 0 , 0 , 0 , 0 , 0 , 0 )	chirality 0
2 x ( 0 , V , 0 , 0 , 0 , 0 , 0 , 0 , V* )	chirality 0
3 x ( 0 , 0 , 0 , 0 , 0 , 0 , 0 , 0 , A )	chirality -1
1 x ( 0 , 0 , 0 , 0 , 0 , 0 , 0 , S , 0 )	chirality 1
1 x ( 0 , 0 , 0 , V , V , 0 , 0 , 0 , 0 )	chirality 1
3 x ( 0 , 0 , 0 , 0 , 0 , 0 , V , 0 , V )	chirality 1
2 x ( 0 , 0 , 0 , 0 , 0 , 0 , 0 , 0 , S )	chirality -2
1 x ( 0 , 0 , V , 0 , 0 , 0 , V , 0 , 0 )	chirality 1
1 x ( 0 , 0 , 0 , V , 0 , V , 0 , 0 , 0 )	chirality -1
1 x ( 0 , 0 , 0 , 0 , 0 , A , 0 , 0 , 0 )	chirality -1
5 x ( 0 , 0 , 0 , 0 , 0 , 0 , 0 , A , 0 )	chirality -1
1 x ( 0 , 0 , 0 , 0 , V , V , 0 , 0 , 0 )	chirality -1
3 x ( 0 , 0 , 0 , 0 , 0 , 0 , 0 , V , V )	chirality 1
2 x ( 0 , 0 , 0 , 0 , 0 , 0 , 0 , V , V* )	chirality 2
2 x ( 0 , 0 , 0 , 0 , V , 0 , 0 , 0 , V )	chirality 2
1 x ( 0 , 0 , V , 0 , 0 , 0 , 0 , V , 0 )	chirality 1
1 x ( 0 , 0 , 0 , V , 0 , 0 , V , 0 , 0 )	chirality 1
1 x ( 0 , 0 , 0 , V , 0 , 0 , 0 , V , 0 )	chirality -1
1 x ( 0 , 0 , 0 , 0 , 0 , V , V , 0 , 0 )	chirality 1
1 x ( 0 , 0 , 0 , 0 , 0 , V , 0 , V , 0 )	chirality -1
3 x ( 0 , 0 , 0 , 0 , V , 0 , 0 , V , 0 )	chirality -3
6 x ( 0 , 0 , 0 , 0 , 0 , 0 , V , V , 0 )	chirality 0
1 x ( 0 , 0 , 0 , 0 , S , 0 , 0 , 0 , 0 )	chirality 0
1 x ( 0 , 0 , 0 , 0 , 0 , 0 , A , 0 , 0 )	chirality 0
2 x ( 0 , 0 , 0 , 0 , 0 , 0 , 0 , Ad , 0 )	chirality 0
2 x ( 0 , 0 , 0 , 0 , 0 , 0 , 0 , 0 , Ac )	chirality 0
1 x ( 0 , 0 , 0 , 0 , 0 , 0 , S , 0 , 0 )	chirality 0

Hidden sector



# GUTS IN STRINGS?

## CHARGE QUANTIZATION



Standard model:

$$\frac{t}{3} + \frac{s}{2} + Y = 0 \pmod{1}$$

Explained by SU(5) GUTs!

Heterotic with standard U(1) normalization:  
spectrum always contains particles violating this relation\*  
(may be heavy or confined by other interactions).

Orientifolds with  $x=1/2$ : half-integer charges in the  
OH-sector (if there is a hidden sector)

Orientifolds with  $x=0$ : SM charge quantization satisfied perturbatively.



True unification (higher level affine,  $x=0$  orientifolds,  
SU(5) F-theory ...)



Heterotic strings, most  $x=1/2$  orientifolds

(\* ) X. Wen and E. Witten, Nucl.Phys.B261:651,1985  
A.N Schellekens, Phys.Lett.B237:363,1990



# VACUUM SELECTION

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Could true GUTs with  $M_{\text{GUT}} < M_{\text{planck}}$  be preferred over generic string vacua?

- Aesthetics: makes no sense.
- Statistics / Vacuum counting: seems unlikely.
- Cosmologically?
- Anthropically?
  - Low energy gauge couplings don't care about unification.
  - Baryogenesis might be a candidate, but does not seem to work



# GUTS IN STRINGS?

---

- Beautiful! 
- Unifies 3 of the 4 known interactions 
- Explains family structure 
- Explains charge quantization 
- Predicts  $\sin^2 \theta_w$  



NOTHING  
IS  
BETTER THAN  
GUTS!

(INSPIRED BY W. SIEGEL)











from:

# THEORY OF MORE THAN EVERYTHING\*

V. Gates, Empty Kangaroo, M. Roachcock, and W.C. Gall\*\*

There is, in fact, a simple proof that superpea theory is superior to superstring theory in describing physics:

- (1) Nothing is better than superstrings.
- (2) Superpeas are better than nothing.
- (3) Therefore, superpeas are better than superstrings.