## Related topics

$\gamma$-radiation, interaction with material, photoelectric effect, Compton effect, pair formation, detection probability, scintillation detectors.

## Principle and task

The radiation of ${ }^{137} \mathrm{Cs}$ and ${ }^{22} \mathrm{Na}$ is measured with a scintillation detector and the energy spectrum determined with a pulse height analyzer. The fractions of the spectra caused by Compton scattering and those caused by the photoelectric effect are determined on the basis of their areas. The results are used for determining the ratio of the effective cross-sections and examining its energy dependence.

## Equipment

Radioactive source, $\mathrm{Na}-22,74 \mathrm{kBq}$
Source Cs-137, 37 kBq
09047.52

Pulse height analyser
Gamma detector
Operating unit f. gamma detector High-voltage connecting cable
Oscilloscope, 20 MHz , 2 channels xyt recorder
Support rod -PASS-, square, I 400 mm
Right angle clamp -PASS-
Tripod base -PASS-
Adapter, BNC-socket/4 mm plug pair
Screened cable, BNC, I 750 mm
Connecting cord, 750 mm , red
Connecting cord, 750 mm , blue
09096.01
13725.93
09101.00
09101.93
09101.10
$11454.93 \quad 1$
$11416.97 \quad 1$
02026.551
02040.551
$02002.55 \quad 1$
$07542.26 \quad 1$
07542.113
$07362.01 \quad 2$
$07362.04 \quad 2$

## Problems

1. Measurement of the $\gamma$-spectra of ${ }^{22} \mathrm{Na}$ and ${ }^{137} \mathrm{Cs}$, using a scintillation detector.
2. Determination of the ratio of the specific effective crosssections due to the Compton effect and the photoelectric effect in photons having energy values of 511, 662 and 1275 keV .

## Set-up and procedure

The layout of the detection apparatus is shown in Fig. 1. The detailed directions in the operation instructions for the $\gamma$ detector, the pulse height analyzer, the oscilloscope and the recorder must be complied with. The detection apparatus must be calibrated before each measurement. The following settings are selected:

Detector voltage:

Pulse height analyzer:
Mode of operation:
Window:
Zoom:
Base:
Dwell time:
Amplification:

Recorder:
10.00 scale divisions in the case of ${ }^{137} \mathrm{Cs}$
8.50 scale divisions in the case
of ${ }^{22} \mathrm{Na}$
automatic
200 mV
Off
10.00 scale divisions
1.6 s

Coarse and fine adjustment to be made with the oscilloscope in such a way that the highest pulses at the analogue output reach a height of 9 V
Input amplification to conform with the paper format; zero point and coordinates to be marked exactly.

Fig. 1: Experimental set-up for $\gamma$-spectroscopy.


After these settings have been completed, no furhter changes may be made to the detection apparatus or to the counter geometry.

Photon detectionis based on interactions between the electromagnetic radiation and the electrons of the scintillator material. A distinction has to be made between the photelectric effect, the Compton effect and pair formation. The specific cross-sections for each of these effects are critical in determining the shape of a spectrum and the integral detection probability. They are to a large extent energy-dependent.

Pair formation requires an energy of at least 1022 keV , which is double the resting energy of the electron. This effect can therefore be disregarded in the energy range in question.

In the case of the photoelectric effect, the electromagnetic field of a $\gamma$-quantum interacts with an electron in accordance with the laws of quantum mechanics, transmitting its complete energy to this electron which continues to move preferentially in the direction of the $\gamma$-quantum. Photoelectrons are represented by a sharply defined line in the $\gamma$-spectrum. The interaction integral assumes imperceptibly small values at high energy levels.

The Compton effect is a corpuscular property of the $\gamma$-quantum. It results from the laws of conservation of energy and of momentum, using a conventional calculation, that on impact a $\gamma$-quantum imparts a proportion of its energy to the impact electron; this proportion is a function of the direction of scattering and amounts to:

$$
\Delta E=\frac{E 2 \Lambda \sin ^{2} \frac{\Phi}{2}}{2 \Lambda \sin ^{2} \frac{\Phi}{2}+\frac{h \cdot c}{E}}
$$

in which $\Lambda=h / m_{o} c^{2}$ denotes the Compton wavelength.
The effective cross-section for the Compton effect is obtained only from a quantum-electro-dynamic calculation, which yields the Klein-Nishina equation.


Fig. 2: Cross-section for the Compton effect $\sigma_{\mathrm{C}}$ in acordance with the Klein-Nishina equation.

Fig. 3: $\gamma$-spectrum of ${ }^{137} \mathrm{Cs}$.


The Compton-scattered fraction in a $\gamma$-spectrum is represented by a continuous plateau which extends from the energy 0 to a steep fall at the Compton border; the electrons detected here are generated by a central impact and entrain the maximum possible energy. The area portions in the spectrum provide a direct measurement of the occurrence of the particular effect.

Photoelectric lines have an approximately gaussian shape, so that their area $A_{F}$ is obtained from the peak height $\alpha$ and the half width $b$ in accordance with the following formula

$$
A_{F}=\alpha \cdot b \sqrt{\pi} / 2 \cdot \sqrt{\ln 2} .
$$

The Compton background can be approximated by a rectangular area $A_{C}$.

The ${ }^{137} \mathrm{Cs}$ spectrum includes in addition to the photoelectric line a further X-ray line which can be disregarded in the evaluation.

In the ${ }^{22} \mathrm{Na}$ spectrum the Compton background of the 1275 keV line is used as the startline for evaluation of the 511 keV line. A rise is observed in the Compton background of the 511 keV line, which sould not be taken into account in the evalua-


Fig. 4: $\gamma$-spectrum of ${ }^{22} \mathrm{Na}$.

Fig. 5: Compton fraction and photoelectric fraction in the spectra in Fig. 3 and 4.

tion, since it originates from $\gamma$-quanta which have undergone back-scattering in the source holder and in consequence have a lower energy.

The following areas are obtained from the spectra in Figs. 3 and 4:
$A_{\mathrm{F}}=18 \mathrm{~cm}^{2}$ and $A_{\mathrm{C}}=25.6 \mathrm{~cm}^{2}$ where $E=511 \mathrm{keV}$.
The ratio of the areas is therefore $A_{\mathrm{F}} / A_{\mathrm{C}}=0.70$.
$A_{\mathrm{F}}=19.5 \mathrm{~cm}^{2}$ and $A_{\mathrm{C}}=36.7 \mathrm{~cm}^{2}$ where $E=622 \mathrm{keV}$.
The ratio of the areas is therefore $A_{\mathrm{F}} / A_{\mathrm{C}}=0.53$.
$A_{\mathrm{F}}=2,8 \mathrm{~cm}^{2}$ and $A_{\mathrm{C}}=16.3 \mathrm{~cm}^{2}$ in the case of $E=1275 \mathrm{keV}$.
The ratio of the areas is therefore $A_{\mathrm{F}} / A_{\mathrm{C}}=0.17$.

Fig. 6: Specific effective cross-section for the photoelectric effect $\sigma_{F}$.


On the assumption that the area ratio corresponds to the effective cross-section ratio, the specific cross sections for the photoelectric effect $\sigma_{\mathrm{F}}$ are obtained by inserting the numerical values for the cross-section in the case of the Compton effect $\sigma_{\mathrm{C}}$ from Fig. 2.

